

Achieving Ultrahigh Piezoelectricity in Organic–Inorganic Vacancy-Ordered Halide Double Perovskites for Mechanical Energy Harvesting

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1 Achieving Ultrahigh Piezoelectricity in Organic-

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3 Perovskites for Mechanical Energy Harvesting

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3 Perovskites for Mechanical Energy Harvesting

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1 **ABSTRACT**

2 Piezoelectric charge coefficient (d_{33}) and piezoelectric voltage coefficient (g₃₃) are the two most 3 critical parameters that define output performance of piezoelectric nanogenerators (PNGs). Herein, 4 we propose a vacancy-ordered double perovskite of TMCM₂SnCl₆ (where TMCM is 5 trimethylchloromethyl ammonium) with a large d₃₃ of 137 pC/N and g₃₃ of 980 \times 10⁻³ V •m/N. The 6 piezoelectric coefficients are considered from the halogen-bonding mediated synergistic 7 movements of atomic displacement in inorganic $[SnCl₆]^{2–} octahedrons, as well as the molecular$ 8 rotation of organic $TMCM⁺$, which is revealed by a combined density functional theory (DFT) and experimental study. The TMCM₂SnCl₆ possesses a high saturated polarization (P_s) of 8.7 μC/cm², 10 a high Curie temperature (T_c) of 365 K, and a low coercive field (E_c) of 0.6 kV/cm. The output 11 voltage (V_{oc}) and current (I_{sc}) of the PNGs are 81 V and 2 μ A at an applied mechanical excitation 12 of (4.9 N, 40 Hz). We hope this work will provide guidance in energy harvesting by innovatively 13 designing highly piezoelectric perovskites for the PNGs. 3 critical parameters that define output performance of piezeslectric nanogenerators (PNGs). Herein,

4 ve propose a vacancy-ordered double perrossite of 'UMCM-SnCl_L (where TMCM is

5 trianelly/chloronellyd ammonium) wi

14 **TOC GRAPHICS**

1 Piezoelectric nanogenerators (PNGs) have been emerging as a promising power source for self-2 powered electronics owing to their direct power conversion from mechanical to electrical energy.¹⁻ $5\degree$ To maximize the output power of PNGs, both the d₃₃ and g₃₃ of the piezoelectric host are 4 important, which determines the output current $(I_{sc}=(d_{33}\times\Delta F)/\Delta t$, where ΔF is the applied force 5 and Δt is the time) and voltage (V_{oc}=g₃₃× Δ P×L, where Δ P is the applied pressure and L is the 6 original film thickness), respectively.⁶⁻⁹ In the past decade, a wide range of piezoelectric materials, 7 targeting high g₃₃ or d₃₃, have been synthesized for the efficient PNGs. For example, the organic 8 polyvinylidene fluoride (PVDF) possesses a high g₃₃ (\sim 286 ×10⁻³ V •m/N), leading to a high output 9 piezoelectric voltage. Unfortunately, the resultant current is limited due to its relatively-low d₃₃ 10 (-30 pC/N) .¹⁰⁻¹³ Conversely, inorganic perovskite oxide ceramics, including PbZr_xTi_{1-x}O₃ (PZT) 11 and BaTiO₃ (BTO), exhibit a high d₃₃ (>100 pC/N) but their g₃₃ is rather low (~20 ×10⁻³ V•m/N).¹⁴⁻ 12 ¹⁸ Considering the relation between d₃₃ and g₃₃ (g₃₃=d₃₃/($\varepsilon_0 \times \varepsilon_r$)), where ε_r is material relative 13 permittivity, it is challenging to improve d_{33} while maintaining a high g₃₃ because the increase of 14 d₃₃ is usually associated with an even larger increase of ε_{r} ¹⁹ ³ To mustimize the output power of PNGs, both the d₃ and g₃ of the piezoslectric host are
important, which determines the output current (h_i-(d₃/ΔF)/ΔL, where ΔF is the applied force
for and Δt is the time) and

15 Organic-inorganic hybrid perovskites (OIHPs) combine the merits of both organic molecules 16 and crystalline inorganic solids at a molecular level. The OIHPs were firstly introduced to solar 17 cells (SCs) due to their excellent optical properties in 2009 and then adopted in piezoelectric 18 reseach.²⁰⁻²⁵ In 2015, the MAPbI₃ was firstly employed in PNGs which yielded an output voltage 19 and current density of 2.7 V and 140 nA/cm² at a pressure of 0.5 MPa.²⁶ Subsequently, a series of 20 perovskites, including MASnBr3, FAPbBr2I and Cl/Br-MAPbI3, have been developed for efficient 21 PNGs via the piezoelectric composite films.²⁷⁻³¹ However, the aforementioned OIHPs evolved 22 from the shadows of the SCs possess limited piezoelectric coefficients. Therefore, extensive efforts 23 have been devoted to searching for highly piezoelectric OIHPs. Remarkably, one-dimensional

1 ABX3-type OIHPs with high d33, including the TMCM-MnCl3 and A-site mixed 2 (TMFM)_x(TMCM)_{1-x}CdCl₃, were synthesized in 2017 and 2019, respectively.^{32, 33} Twodimensional (ATHP)₂PbBr₄ OIHPs with large d₃₃ of 75 pC/N and giant g₃₃ of 660 ×10⁻³ V•m/N 4 were reported in 2019.³⁴ In 2020, the OIHPs were further developed and successfully employed in 5 energy harvesting applications, including (FMTMA)PbCl₂I, HETMA-CdCl₃, (TMFM)FeBr₄ and 6 [Ph₃MeP]₄[Ni(NCS)₆].³⁵⁻³⁸ Despite their large piezoelectric coefficients, the newly-developed 7 OIHPs may suffer substantial difficulties in mass production and implementation. For example, 8 the TMCM-MnCl3 has intrinsic chemical instability as the bivalent Mn is prone to be oxidized to 9 tetravalent, while the Cd (or Pb) in $(TMFM)_x(TMCM)_{1-x}CdCl_3$ (or $(ATHP)_2PbBr_4$) is not 10 environmental-friendly. Unlike the ABX₃-type OIHPs, the A₂BX₆ OIHPs (such as K₂PtCl₆) are 11 potentially more chemically stable because their B-site is already in the tetravalent state.³⁹⁻⁴¹ 3 dimensional (ATHP),PhBHz OHPs with large d₃₀ of 75 JC/N and giant g₂₃ of 660 \times 10⁻³ V +m/N
were reported in 2019.⁵⁴ In 2020, the OHPs were further developed and saccessfully employed in
energy harvesting appli

12 Herein, we proposed an alternative crystal model, which was vacancy-ordered double OIHPs, 13 for novel piezoelectric material design. The chemical formula was confirmed as $T M C M_2 S n C l_6$ via 14 powder X-ray diffraction (XRD) and X-ray photoelectron spectroscopy (XPS). The large d_{33} of 15 137 pC/N and giant g₃₃ of 980 \times 10⁻³ V • m/N were experimentally measured for the prepared 16 TMCM₂SnCl₆ crystal, which was further confirmed by the DFT study and the outputs of PNGs. 17 The PNGs were fabricated via sandwiching the TMCM₂SnCl₆@PDMS composite film into two 18 copper electrodes, where PDMS represents Poly(dimethylsiloxane). A self-powered wireless 19 communication node (SWCN) was successfully driven by the PNG. The detailed material 20 synthesis and device fabrication processes are shown in the supporting information.

21

2 **Figure 1. Crystal structure of TMCM2SnCl6 double perovskites:** (a) structural model; (b) 3 calculated and experimental XRD patterns; (c) digital photograph of single crystal; (d) halogen 4 bonding interaction. The arrow from Cl to N indicates the dipole direction of $TMCM^+$; (e) FTIR 5 spectra of pure TMCM-Cl molecules and TMCM₂SnCl₆; (f) calculated energy mapping of halogen 6 bond for the simplified model in Figure 1d.

 8 TMCM₂SnCl₆ is a vacancy-ordered double perovskite because its 50% $[SnCl₆]$ ² octahedrons 9 are replaced by periodic vacancies, in which the asymmetric TMCM⁺ takes the A sites (**Figure** 10 **1a**). This is different from the reported vacancy-ordered double perovskites with spherical 11 symmetric cations like trimethylmethyl ammonium (TMA) or monovalent alkali metals (Na, K, C s...) taking the A sites.⁴²⁻⁴⁴ Figure 1b shows the XRD measured from the TMCM₂SnCl₆ single 13 crystals (black curve) in comparison to the simulated XRD (red curve) from the projected structure 14 5 spectra of pure TMCM-Cl molecules and TMCM₂.
6 bond for the simplified model in Figure 1d.
7
7 TMCM₂SnCl₆ is a vacancy-ordered double per
9 are replaced by periodic vacancies, in which the
10 **1a**). This is differ in Figure 1a. The single crystal of TMCM₂SnCl₆ with a length of 1.8 cm is shown in **Figure 1c**.

1 The agreement of XRD peaks confirms the crystal structure and its high phase purity. The 2 TMCM₂SnCl₆ was aged to 6 months in an ambient atmosphere and the XRD was measured in 3 **Figure S1a**. The reflection peaks of TMCM₂SnCl₆ almost remained unchanged and no new 4 reflection peaks appeared. These indicate the TMCM₂SnCl₆ crystals have not degraded into other 5 chemicals and exhibit good stability. The TMCM₂SnCl₆ also exhibits good thermal stability, as 6 indicated by a negligible change in the measured permittivity at 340K for one hour (**Figure S1d)**. 7 The chemical states of (TMCM)2SnCl6 were measured via XPS in **Figure S2**. The peaks at binding 8 energy of 495.8 eV and 487.3 eV are assigned to Sn^{4+} and corresponding to its 3d_{3/2} and 3d_{5/2} 9 orbitals, respectively. The peaks at 200.6 eV and 198.8 eV are assigned to $3p_{1/2}$ and $3p_{3/2}$ orbitals 10 of Cl⁻. This is well consistent to the chemical states in [SnCl₆]^{2−} octahedrons, where Sn is in 11 tetravalent state and Cl in minus state.

12 A halogen bond has been considered as an efficient driving force to achieve the well-defined 13 structures during crystal growth for piezoelectric properties.^{45, 46} In TMCM₂SnCl₆, there is a net 14 attractive interaction between an electrophilic region associated with the Cl atom in TMCM⁺ and 15 the nucleophilic Cl⁻ in the [SnCl₆]^{2−} octahedron (**Figure 1d**). The attractive nature of the halogen 16 bond is determined by the bond distance (R, in the range of 2-6 Å) and the bond angle (θ, in the 17 range of 150-180^o), especially when R is shorter than the sum of Van der Waals radii (\sim 3.5 Å). 18 The halogen bond in the $TMCM₂SnCl₆$ was reflected by Fourier transform infrared spectroscopy (FTIR) in **Figure 1e**. The absorption peak at 793 cm[−]¹ 19 (black line) is attributed to the stretching 20 vibration of Cl-C bond of TMCM⁺. For the TMCM₂SnCl₆, the absorption peak has a red-shift of 14 cm^{-1} to 807 cm⁻¹ (red line). The red shift should be attributed to the increased force constant 22 when the electrophilic Cl in TMCM⁺ is attracted by the nucleophilic Cl in $[SnCl_6]^2$ octahedron 23 according to harmonic motion for diatomic model.^{47, 48} The halogen bond energy was calculated 3 Figure S1a. The reflection peaks of TMCM₂SnCl₆ almost remained unchanged and no new
reflection peaks appeared. These indicate the TMCM₂SnCl₆ exposuls have not degraded into other
chemicals and exhibit good stabi

1 by using DFT as a function of the bond distance (R) and the bond angle (θ) in **Figure 1f and Figure S3**. With the decreasing of R, the calculated system energy of $TMCM^+ \cdots [SnCl_6]^2$ 3 decreases as the increased Coulomb attraction, and followed by a quick increase afterwards due to 4 the dominating nuclear repulsion. The halogen bond energy is evaluated from the absolute value 5 of valley energy. The lowest energy of halogen bond is observed to be 0.021 Hartree at a $Cl...Cl$ 6 bond distance of 3.06 Å for a certain C-Cl…Cl bond angle (θ = 150°). The bond energy increases 7 from 0.021 to 0.024 Hartree (dash line) as θ increases from 150° to 180°. ion, and followed by a quick increase afterwards due to

ogen bond energy is evaluated from the absolute value

ogen bond is observed to be 0.021 Hartree at a CI···CI
 \cdot ··CI bond angle ($\theta = 150^\circ$). The bond energy in

9 **Figure 2. Piezoelectric and ferroelectric properties of TMCM2SnCl6 double perovskite**: 10 vertical PFM amplitude (a) and phase (b) images; (c) current density-voltage (J-V) curve and the 11 integrated polarization-voltage (P-V) hysteresis loop of the TMCM2SnCl₆ film at room 12 temperature (25 \degree C) by using the double-wave method; (d) temperature-dependent J-V curves. 13 Accept to the same of the sam Voltage was swept from -70 to 70 V. The inset is the differential scanning calorimetry (DSC) curves

1 of TMCM₂SnCl₆ crystals with the scanning rate of 10 K/min; (c) frequency-dependent permittivity 2 tested from 0.01 to 1 MHz; (f) calculated density of states of $TMCM_2SnCl_6$.

3 The d₃₃ of TMCM₂SnCl₆ perovskites was measured to be 137 pC/N by using a quasi-static 4 method in **Figure S4**. The ferroelectric domains were confirmed by a clear contrast in amplitude 5 and phase images of vertical piezoelectric force microscopy (V-PFM) in **Figure 2a-b**. There was 6 no crosstalk between the topography (**Figure S5b)** and the V-PFM phase. Meanwhile, the lateral 7 PFM (L-PFM) was recorded in **Figure S5c-d** where the signal strength was found relatively weak, 8 indicating that the spontaneous polarization was almost along the out-of-plane direction.³³ The P-9 V hysteresis loop was measured for $TMCM_2SnCl_6$ films with the thickness of 400 μ m by using 10 the double-wave method (DWM) in **Figure 2c**.^{49, 50} The arrows in the J-V curve indicate the 11 voltage sweep directions (-60 V \rightarrow 0 V \rightarrow 60 V \rightarrow 0 V \rightarrow -60 V). The J-V curve exhibits two 12 apparent current peaks at the ferroelectric switching voltage of ± 23 V. The E_c of TMCM₂SnCl₆, 13 expressed by the ratio of switching voltage to the film thickness, is 0.6 kV/cm. This value is 14 comparable to that of conventional piezo-ceramics such as bulk BTO single crystals (1 kV/cm) .⁵¹ The P_s from the P-V hysteresis loop is up to 8.7 μ C/cm², close to the calculated value of 8.4 \pm 0.3 16 $μC/cm²$. When the temperature is higher than T_c, the ferroelectric material will undergo a phase 17 transition and lose its spontaneous polarization. The phase transition of TMCM2SnCl₆ was 18 considered from orthorhombic to cubic, as indicated by some disappearing diffraction peaks at 19 around 26.3° , 33.5° and 36.5° in temperature-dependent XRD (**Figure S1b**). The T_c of 20 TMCM₂SnCl₆ was studied via monitoring the change of peak intensity at 23 V in the temperature-21 dependent J-V curves. As indicated in **Figure 2d**, the current peaks disappear when the 22 temperature rises to 353-363 K. This is consistent to a pair of thermal anomalies at 365/387 K in 23 the differential scanning calorimetry (DSC) curves (the inset of Figure 2d) and an increase at 350- 3 The day of TMCMASnCl₄ perovskites was measured to be 137 pCN by using a quasi-state
method in Figure **S4**. The ferroelectric domains were confirmed by a clear contrast in amplitude
5 and phase images of vertical pieco

1 380 K in the temperature dependent permittivity (**Figure S1c**). The TMCM2SnCl6 possesses a Tc 2 above room temperature, making it suitable for PNG applications. The ε_r of perovskite oxide 3 ceramics, such as PZT, is observed as thousands, which stems from the hopping polaron within 4 semiconducting grains and interfacial polarization at the insulating grain boundary.⁵² The 5 TMCM2SnCl6 exhibits a relative low εr of 14-12 with the frequency from 0.01 to 1 MHz in **Figure** 6 **2e.** This should be attributed to the insulating properties of TMCM2SnCl6: where the calculated 7 bandgap of above 3.2 eV and a narrow conduction band width are defined by the 5s orbits of Sn 8 and the 3p orbits of Cl, as indicated in **Figure 2f and Figure S6**. According to $g_{33}=d_{33}/(\epsilon_0 \times \epsilon_r)$, 9 where ε_0 is the permittivity of free space and $\varepsilon_r = 16$, the calculated g₃₃ is 980 ×10⁻³ V • m/N. Notably, 10 a higher ε_r was chosen to match the following PNG working frequency at 40 Hz. In comparison to 11 traditional metal oxides and other molecular materials listed in **Figure 3a** (or **Table S1**), the 12 TMCM₂SnCl₆ perovskites are more favorable for the PNG applications due to the both large d_{33} 13 and giant g₃₃. 3 ceramics, such as PZT, is observed as thousands, which stems from the hopping polaron within
4 semiconducting grains and interfactal polarization at the insulating grain boundary.⁵² The
5 TMCM:SaCl_e exhibits a relat

15 • Figure 3. Origin of piezoelectricity in the TMCM₂SnCl₆ double perovskite: (a) the d_{33} and g_{33}

1 of TMCM2SnCl6 compared with the reported inorganic and molecular materials, including BTO, 2 PZT, Triglycine sulfate (TGS), Imidazolium perchlorate (IP), PVDF, and (ATHP)₂PbBr₄; (b) the 3 calculated polarization as a function of c_1 by using Berry phase method; (c) side views of the 4 unstrained structure and (e) corresponding deformation charge density for $c = 9.33864 \text{ Å}$; (d) the 5 strained structure and (f) deformation charge density for $c = 9.29311$ Å. The above deformation 6 charge density is in the $[001]$ planes including the Sn-Cl…Cl-C chain. The blue and red colors 7 represent the electron loss and gain, respectively. by using Berry phase method; (c) side views of the
deformation charge density for $c = 9.33864 \text{ Å}$; (d) the
ge density for $c = 9.29311 \text{ Å}$. The above deformation
ding the Sn-Cl···Cl-C chain. The blue and red colors
tiv

8

9 To investigate the piezoelectric origin of $T M C M_2 S n C l_6$ double perovskites, the polarization 10 (P_{total}) was calculated by the Berry phase method as a function of c_1 , where c_1 is the lattice constant 11 during the applied strain along the c axis (**Figure 3b**). The slope of the linear P_{total} vs. c_1 curve 12 yields the piezoelectric stress coefficient (e₃₃).^{53, 54} By fitting the P_{total} vs. c_{1/c0} curve in Figure 3b, 13 e₃₃ is evaluated as 1.7 C/m^2 . According to the calculated elastic modulus tensor (c₃₃ = 187.7 kBar) 14 in **Table S2,** the d₃₃ is calculated as 90 pC/N by using the expression of d₃₃=e₃₃/c₃₃. The calculated 15 d₃₃ is lower than the experimental value of 137 pC/N as the P_{total} contribution resulted from lateral 16 strain was neglected.⁵⁵ Figure 3c and 3d show the crystal structure (side view) of unstrained 17 TMCM₂SnCl₆ with a lattice constant c₁ of 9.33846 Å and that of strained TMCM₂SnCl₆ with c₁ of 18 9.29311 Å (corresponding to a strain of -0.4%), respectively. Under the applied strain, the angle $(α)$ in the dipole direction of TMCM⁺ with respect to c axis decreases from 12.31^o to 9.907^o. 20 According to the equation of $P_{TMCM} = \mu \times \cos(\alpha)/V_u$, where μ is the dipole moment of TMCM⁺ (4.65 21 Debye) and V_u is the cell volume, the polarization from the TMCM⁺ increases when α decreases. 22 The deformation charge density profiles of the unstrained and strained crystal were shown in 23 Figure 3e and 3f, respectively. The Cl atom in TMCM⁺ is involved in the formation of the covalent e₃₃ is evaluated as 1.7 C/m². According to the calc

14 in **Table S2**, the d₃₃ is calculated as 90 pC/N by us

15 d₃₃ is lower than the experimental value of 137 pC

16 strain was neglected.⁵⁵ **Figure 3c** and

1 bond. In addition, there is a region of lower electron density (σ -hole) where the potential is positive, 2 which generates a cap of depleted electron density on the elongation of the covalent bond. Thus, a 3 halogen-bonding interaction will be formed between the Cl of TMCM⁺ and the electron-rich Cl 4 atoms in $[SnCl₆]²$. Therefore, the rotation of TMCM⁺ can trigger the atomic displacement in 5 $[SnCl₆]²⁻ octahedrons simultaneously, indicated by the increased angle (β) of Cl-Sn-Cl in$ 6 [SnCl₆]^{2–} octahedrons from 175.6[°] to 176.5[°]. The polarization from [SnCl₆]^{2–} is calculated by the 7 equation of $P_{Sn}=P_{total}-P_{TMCM}$. Via fitting the decomposed polarization P_{TMCM} and P_{Sn} in **Figure S7**, 8 the e₃₃ of $[SnCl_6]^2$ ⁻ and TMCM⁺ is 1.5 C/m² and 0.2 C/m², respectively. It is clear that the $[SnCl_6]^2$ ⁻ 9 contributes nearly all of the piezoelectricity. After replacing the asymmetric $TMCM⁺$ by spherical 10 symmetric TMA⁺, the [SnCl₆]^{2−} in TMA₂SnCl₆ is not distorted (**Figure S8**). The polarization from 11 TMA2SnCl6 is 0 and not changed under an applied strain, indicated the loss of piezoelectricity. 12 Unlike the traditional perovskite ceramics where the piezoelectricity is primarily arising from the 13 atomic displacement (**Figure S9**), (i) the organic TMCM⁺ in TMCM₂SnCl₆ can break its crystal 14 symmetry and (ii) the piezoelectricity in TMCM₂SnCl₆ is from the halogen bond-mediated 15 synergistic movement of atomic displacement in inorganic $[SnCl₆]²⁻ octahedrons$ and the 16 molecular rotation of TMCM⁺. 3 halogen-bonding interaction will be formed between the Cl of TMCM and the electron-rich Cl

4 atoms in [SaCk]². Therefore, the triation of TMCM² can trigger the atomic displacement in

5 [SaCk]² octabedrons simult

17 The piezoelectricity in the $T M C M_2 SnCl_6$ was further demonstrated by the outputs of fabricated 18 PNGs. The precursor of piezoelectric composite films was prepared via dispersing 10 wt. % 19 TMCM2SnCl6 nanoparticles (NPs) uniformly in the PDMS by mechanical stirring (**Figure S10**). 20 The size of the NPs was in the range of 200-1000 nm, as shown in the SEM and elemental mapping 21 (Figure S11b-c). The films were grown via drop-casting the above precursor solution on a glass 22 substrate (Figure S11a). After the drop-casting, NPs with a larger size tended to precipitate to the 23 bottom of central area rather than diffuse to the marginal area (**Figure S11a, d-f**). The fraction of

1 TMCM2SnCl6 NPs in the composite film for the PNG integration was 18 wt. % measured in the 2 dotted circle in Figure S11d by the SEM-EDS. The film thickness at the central area was 300 μm. 3 Finally, the TMCM2SnCl6@PDMS composite film was sandwiched between two copper 4 electrodes for PNG fabrication (**Figure 4a** and **Figure S12**) (see the experimental details in the 5 supporting information).

7 **Figure 4. The piezoelectric output of single PNG:** (a) device structure of PNGs via inserting 8 TMCM₂SnCl₆@PDMS composite film between copper electrodes; (b) V_{oc} and (c) I_{sc} of PNGs with 9 the force from 2.8 to 4.9 N (40 Hz); (d) the device reliability tested for 10 mins; (e) $I_{\rm sc}$ and (f) 10 enlarged view of the PNGs under forward and reverse electrical connections. The inset is schematic 11 diagrams of the PNGs under forward and reverse connections; (g) polarity switching test of V_{oc} 12 (black curve) and $I_{\rm sc}$ (red curve) of pure PDMS based devices; (h) $V_{\rm oc}$ and (k) $I_{\rm sc}$ of PNGs based 13 on TMCM2SnCl6@PDMS film (without poling). For the Figures 4d-i, the applied mechanical

1 excitation is $(2.8 \text{ N}, 40 \text{ Hz})$.

2 The outputs of the $TMCM_2SnCl_6@PDMS$ based PNGs was measured by mounting the device 3 on a shaker in a vertical compressing and releasing process (**Figure S13**). For the PNGs with 18 4 wt. % TMCM₂SnCl₆ in the composite film, the output V_{oc} and I_{sc} of 40 V and 0.9 μ A was achieved 5 under the force of 2.8 N at a frequency of 40 Hz. The device output is reproducible as indicated 6 by the statistical analysis of V_{∞} in **Figure S14**. With the applied force increased from 2.8 to 4.9 7 N, the piezoelectric outputs (V_{oc} and I_{sc}) of the PNGs increase linearly in **Figure 4b-c**. The V_{oc} 8 and Isc reach 81 V and 2 μA, respectively, at an applied force of 4.9 N. The frequency-dependent 9 Voc and Isc are shown in **Figure S15** as the force keeps 2.8N. Meanwhile, the device outputs were 10 also recorded at low frequency (3 Hz) via a cylindrical hammer mounted on a linear motor and a 11 metallic stopper. The V_{oc} and current density (J_{sc}) were measured as 11.6 V and 1.17 μ A/cm², 12 respectively, at an applied pressure of 0.032 MPa (**Figure S16**). The concentration of 13 TMCM₂SnCl₆ on the electricity generation capacity of the PNGs was optimized in **Figure S18**. 14 The output voltage and current gradually increase to 72 V and 1.3 μA with the increasing mass 15 fraction of TMCM2SnCl6 NPs up to 29 wt. %, under an applied force of 2.8 N. In comparison with 16 traditional inorganic perovskite ceramics in **Table S3**, the TMCM2SnCl6 were prepared at room 17 temperature and the PNGs exhibit the competitive piezoelectric outputs. The reliability of the 18 PNGs was explored by continually testing the device for 10 minutes and the amplitude of the 19 output voltage shows a negligible degradation (**Figure 4d** and Video S1), confirming its excellent 20 mechanical durability. Finally, the PNGs were used as an integrated power source to meet up the 21 energy demand of practical electrical loads via harvesting energy from local environment, such as 22 13 TMCM₂SnCl₆ on the electricity generation capac
14 The output voltage and current gradually increase
15 fraction of TMCM₂SnCl₆ NPs up to 29 wt. %, und
16 traditional inorganic perovskite ceramics in **Table**
17 t SWCN (see the details in **Figure S17** and Video S2). releasing process (**Figure S13**). For the PNGs with 18
the output V_{∞} and I_{∞} of 40 V and 0.9 µA was achieved
40 Hz. The device output is reproducible as indicated
S14. With the applied force increased from 2.

1 To demonstrate the generated outputs were originating from the piezoelectric phenomenon, the 2 following experiments were conducted: (i) the measured current output of the PNGs switched its 3 polarity during the electrical connection reversal (**Figure 4e-f** and Video S3), indicating that the 4 outputs were produced from the PNGs; (ii) For comparison, the output V_{oc} and I_{sc} of PNGs made 5 from pure PDMS were merely 0.2 V and 0.02 μA and lacking of the polarity switching property in 6 **Figure 4g**. The Voc is enhanced with the loading TMCM2SnCl6 NPs in the composite film as shown 7 in **Figure S18**. These verify the claim that - the high performance of PNGs is determined by the 8 incorporation of $TMCM_2SnCl_6$; (iii) the effect of electrical poling on the PNG outputs was 9 investigated (see the poling procedures in the experimental section). Compared with the outputs 10 of unpoling PNGs (**Figure 4h-k**), the PNGs after poling at 4.5 kV (Figure 4b-c) exhibited drastic 11 increase in V_{oc} (from 5.4 to 40 V) and I_{sc} (from 0.05 to 0.9 μ A) under the same measurement 12 conditions (2.8 N, 40 Hz). A detailed comparison of output voltage of PNGs measured under 13 various applied poling fields is shown in **Figure S19**. 3 polarity during the electrical connection reversal (Figure 4e-f and Video S3), indicating that the outputs were produced from the PNGs; (ii) For comparison, the output V_× and L_5 of PNGs made
5 from pure PDMS were

14 In summary, a novel vacancy-ordered double perovskite of TMCM₂SnCl₆ with a large d₃₃ of 15 137 pC/N and giant g₃₃ of 980 \times 10⁻³ V \cdot m/N was proposed for the energy harvesting applications. 16 The TMCM₂SnCl₆ perovskite possesses a high P_s of 8.7 μ C/cm², a high T_c of 365 K, and a low E_c 17 of 0.6 kV/cm. The ultrahigh piezoelectricity in TMCM₂SnCl₆ is attributed to the halogen bonding-18 mediated synergistic movements of atomic displacement in the inorganic [SnCl₆]^{2−} octahedron and 19 the dipole alignment of organic $TMCM⁺$. The piezoelectric coefficients from the $TMCM₂SnCl₆$ 20 was firstly experimentally measured, then verified by DFT study, and finally confirmed by the 21 piezoelectric output of the PNGs. The piezoelectric composite films were prepared by dispersing 22 TMCM₂SnCl₆ NPs into the PDMS matrix. The peak values of V_{oc} and I_{sc} of the fabricated PNGs 23 reached up to 81 V and 2 μ A at an applied mechanical excitation of (4.9 N, 40 Hz). The SWCN

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19 **References**

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