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Photoisomerization of IC₃N. An experimental and theoretical study.

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Graphical Abstract



Highlights

- VUV photolysis of IC₃N in argon matrix leads to IC₂NC and ICNC₂ isomers
- IRC calculations show reaction paths leading to discovered isomers
- IC₃N ions and radical are not formed in detectable amounts for the IR absorption

ABSTRACT: Here we present the results of a joint spectroscopic and quantum chemical study on the photochemistry of 3-iodo-2-propynenitrile (iodocyanoacetylene, IC₃N). Vacuum-UV photolysis of the compound isolated in solid argon was explored. Calculations, carried out at the DFT and/or coupled-cluster level of theory, provided essential data concerning the thermodynamic stabilities, vibrational and electronic energy levels, ionization potentials, and electron affinities of IC₃N-stoichiometry isomers and ions. The photochemical formation of thus far unknown species IC₂NC and ICNC₂ has been evidenced by the experiment and rationalized based on a detailed theoretical approach, involving the excited-state potential energy surfaces.

KEYWORDS: photochemistry, cyanoacetylenes, matrix isolation, IR spectroscopy, computational chemistry

1. Introduction

The direct motivation for this study stems from a larger project devoted to the photochemistry of small alkynylcarbonitriles. In particular, former investigations on

cryogenically isolated, UV-irradiated HC₃N and HC₅N have shown the wealth of photoproduced isomers, some of which being highly energetic and feature unusual chemical bonding.^{1–3} Isomerization could be anticipated also for IC₃N. Moreover, it was of interest to check whether photolyzed IC₃N could be an efficient source of C_3N^- , opening the way towards better characterization of this anion (HC₃N and HC₅N photolyses yielded C₃N⁻ and C₅N⁻, respectively, but in small amounts).

Kloster-Jensen et al. published on the location of IR and UV absorption bands of IC₃N, ⁴ as well as on the electric dipole moment (4.59 D) of this nitrile, inferred from microwave spectra. Its Raman spectroscopy has also been studied.^{5,6} Linearity of the molecule was confirmed with X-ray diffraction measurements.⁷ Bieri et al. and Kuhn et al. measured the photoelectron spectra.^{8,9} To our knowledge, no photochemical studies of IC₃N have ever been reported.

2. Experimental and computational methods

IC₃N was synthesized on a preparative scale by the method of Kloster-Jensen, involving the reaction of HC₃N with I₂/KI alkaline solution.⁴ The product was characterized by ¹³C NMR spectroscopy: ¹³C NMR ((CD₃)₂CO, 100 MHz) δ 22.4 (s, C-I), 67.8 (s, C-CN); 106.1 (s, CN). The chemical shift of the sp carbon next to the iodine atom, observed at high-field, is characteristic of iodoalkynes. Traces of HC₃N could not be completely removed by freeze-pumpthaw cycles.

Gaseous IC₃N was mixed with Ar (Multax s.c.) in ratios ranging from 1:600 to 1:1100, in a vacuum manifold equipped with capacitance manometers (MKS Instruments). Mixtures were deposited at a rate of about 2 mmol/h onto a CsI substrate window, the latter attached to the cold finger of an Air Products *Displex* DE-202S cryostat. In photochemical experiments, the window was continuously VUV-irradiated (see below) throughout the sample deposition. The window was held either at 15-17 K or, exclusively during VUV irradiations, at 19 K. For electronic spectroscopy measurements, gaseous IC₃N (0.13 Torr at room temperature) was mixed with 100 Torr of argon in a 10-cm glass cuvette equipped with quartz windows. Spectra were recorded in the 205-700 nm range with a resolution of 1.0 nm, using a Shimadzu UV-3100 instrument. A Nicolet *Magna* 560 FTIR spectrometer was used for IR absorption measurements in the 4000 -400 cm⁻¹ range with a resolution of 0.5 cm⁻¹.

For far-UV photolysis experiments, a microwave-driven Kr discharge lamp (Opthos Instrument Comp.) was attached to the vacuum shroud of the cryostat, about 5 cm from the sample substrate. The lamp, operated at 20 W, had a broadband spectral characteristic (127-160 nm) with a maximum near 145 nm; the emitted beam was not focused. The slightly elevated deposition temperature (19 K) was applied to boost the photochemical processes.

The Becke's three-parameter hybrid exchange functional including the correlation functional of Lee, Yang, and Parr (B3LYP) was selected for the majority of DFT calculations.¹⁰ ¹¹ The Dunning's correlation-consistent polarized valence triple-zeta basis set,¹² augmented by s. p, d, and f functions (aug-cc-pVTZ) was used for N and C, while for the I atom a basis set with pseudopotentials (hereafter: PP) was employed.¹³ Geometric structures, energies, and harmonic frequencies were calculated for IC₃N, its isomers, and several other potential photolysis products. The harmonic frequencies were scaled with a factor of 0.96 to correct for anharmonicity and for other inherent deficiencies of the DFT approach.^{14,15} Transition states were localized on the electronic ground state potential energy surface (PES) using the Berny algorithm.¹⁶ Their nature was verified by inspection of vibrational frequencies (exactly one imaginary frequency was expected). Further IRC (intrinsic reaction coordinate) calculations indicated the geometries of species linked by a given transition state.¹⁷ In order to describe the mechanism of a photo-induced reaction, knowledge of excited PESs is crucial. Both triplet and singlet excited electronic state energies were calculated using the time-dependent DFT (TD-DFT) version.^{18–20} The Tamm-Dancoff approximation was applied, as it was reported to perform well for numerically unstable cases and was found to provide reliable singlet-triplet separations.²¹⁻²⁵ Excited-state PES landscapes were approximated using vertical excitations for a series of geometries obtained with ground-state IRC calculations.

The geometries of selected lowest excited states of IC₃N and IC₂NC were optimized with the B3PW91 functional and the same basis sets as for the ground state. B3PW91 (differing from B3LYP by a PW91 non-local correlation term)^{26–28} has already been successfully used to describe some excited electronic states of cyanoacetylene family molecules, including HC₉N,²⁹ CH₃-C₃N,³⁰ and CH₃-C₅N.³¹ We verified the reliability of B3PW91-derived vertical excitation energies by applying also the CAM-B3LYP functional ³² with large basis sets (aug-cc-pVQZ for C and N, and aug-cc-pVQZ-PP for the I atom).^{12,33}

Second order perturbation theory (VPT2) was used to predict anharmonic vibrational frequencies and intensities for fundamental,^{34–36} overtone, and combination bands, at the B3LYP level of theory. All DFT calculations were carried out using the Gaussian 09 Rev. E. 01 set of programs.³⁷ The vibrational self-consistent field (VSCF)^{38,39} and correlation-corrected vibrational self-consistent field (VMP2)^{39,40} methods were also used to compute anharmonic fundamental frequencies using a frozen-core version of CCSD(T) (coupled-cluster singles and doubles with perturbative treatment of triples).^{41–46} These two different approaches, VSCF and VMP2, were chosen to check for consistency of the resulting frequencies (variational VSCF computations are not as sensitive to resonance-related artifacts as those carried out with the perturbation methods VMP2 and VPT2. An efficient approach was used in VSCF and VMP2 calculations to generate potential energy surfaces around the equilibrium structures.^{47–49} Fundamental harmonic frequencies were numerically derived, ^{50,51} for the previously optimized ⁵² structures. One- and two-body terms were included. For CCSD(T) computations, Dunning's correlation-consistent polarized valence triple-zeta basis set (cc-pVTZ)¹² was selected for N and C, while a basis set with pseudopotentials (cc-pVTZ-PP)¹³ was used for the I atom. All ab initio calculations of anharmonic vibrational frequencies were performed with the Molpro 2012.1 software.^{53,54} Computational results were visualized using ChemCraft.⁵⁵

3. Results and discussion

Electronic spectroscopy

Table 1 contains vertical electronic excitation energies and the respective oscillator strengths computed by CAM-B3LYP/aug-cc-pVQZ for IC₃N. Optimization of excited-state structures starting from a linear geometry using B3PW91/aug-cc-pVTZ failed for most states or led to imaginary frequencies. This suggested a non-linearity of these states. Indeed, our search for the equilibrium geometry of the lowest excited singlet (S₁) and triplet (T₁) revealed the bent structures (Figs. S1 and S2). Similar computations were performed for IC₂NC, a likely photochemical product arising from UV irradiation of IC₃N (see Table S1 and Figs. S1-S2 of the Supplementary material). The predicted 0-0 (i.e. vibrationless) energies of S₁-S₀ excitations are 3.3 eV and 3.2 eV (TD-B3PW91/aug-cc-pVTZ) for IC₃N and IC₂NC, respectively.

A UV-Vis spectrum of gaseous IC₃N is presented in Fig. 1. The molecule is a strong absorber. One can discern a structure resembling that previously reported, namely the bands at 222, 233, 245 and 262 nm, measured by Kloster-Jensen et al. in a cyclohexane solution.⁵⁶ The band pattern is similar to the one known for HC₃N,^{57–59} but IC₃N bands are distinctly broader.³⁰ The spectrum shows clear maxima at 44 400, 41 400, 39 300, and 36 600 and cm⁻¹ (225, 242, 254, and 273). According to the predictions of Table 1, there are three closely spaced electronic states around 250 nm. While at least one of these is responsible for the said (presumably vibronic) bands, no reliable assignments can currently be proposed.



Figure 1. UV-Vis absorption spectrum of IC₃N, as measured for a gas-phase sample. Bars indicate the predicted location and oscillator strength of vertical electronic transitions, derived with TD-CAM-B3LYP/aug-cc-pVQZ.

IR spectroscopy

An IR absorption spectrum of IC₃N in Ar obtained for a sample deposited without photolysis (see Fig. 2) served as a reference for comparison with a photo-processed matrix. Fundamental vibrational bands have been identified based on data published (Table 2).^{6,60} A high S/N ratio allowed assignment of new combination bands and of several spectral features belonging to ¹³C and ¹⁵N isotopologues of IC₃N present at their natural abundance (see Tables S2-S3 and the insets of Fig. 2); this was assisted with DFT predictions of isotopic shifts. Furthermore, it was possible to identify, for the first time, various combination and overtone vibrational bands of the molecule (Table 2). Several weak, unassigned bands may either be due to more complicated combinations of IC₃N vibrational modes or to impurities.



Figure 2. IR absorption spectrum of IC₃N in solid Ar at 15 K. Asterisked bands arise from incomplete compensation of CsI substrate absorption.

Theoretically considered IC₃N photolysis products

IC₃N is a halogenated analogue of the extensively studied molecule HC₃N. UV irradiation of the latter in inert cryogenic matrixes leads to HCCNC, HCNCC, and CCCNH.^{1,62–64} We computed the electronic energies, structures, and vibrational IR absorption parameters of the 3 analogous isomers bearing iodine in place of hydrogen (Fig. 3). The ensuing relative energy values and molecular geometries turn out to be similar to those characterizing the HC₃N family of isomers. The most stable is the nitrile IC₃N, the isonitrile IC₂NC being more energetic by about 137 kJ/mol. Much less stable are two other chain structures: C₃NI and C₂NCI. No bound Y-shaped carbene structure C₂(I)CN, analogous to the cyanovinylidene molecule,^{65–67} has been found in our B3LYP study, but CAM-B3LYP computations suggest the presence of a corresponding shallow potential energy minimum. IR spectroscopic parameters predicted for the lowest-energy IC₃N molecules can be found further in the text and in the Tables S3 – S5 of the Supplementary material.



Figure 3. Relative ZPE-corrected electronic energies and structures (Å, deg) of the selected IC₃N isomers (NBO charges are given for each atom), as derived with B3LYP/aug-ccpVTZ (aug-cc-pVTZ-PP for I atom). See Table S5 for the computed ionization energy and electron affinity values.

Table S5 of the Supplementary material presents ionization energies predicted for the four considered isomeric variants of IC₃N. What follows from these results is that none of the respective cations should be expected to appear out of IC₃N irradiated with the presently applied far-UV lamp ($\lambda > 127$ nm). Neither was any obvious source of electrons, leading to the corresponding anions, available in our experiment. Nevertheless, Table S5 provides also the

electron affinities, as these, together with the structural parameters given in Fig. S3, might be useful in the interpretation of future experiments. The calculated ionization energy of IC₃N (989 kJ/mol) is a good match to the value found experimentally, i.e. 974.5 kJ/mol (10.1 eV or 123 nm).⁹

Our theoretical treatment of permanent dissociation processes was limited here to the consideration of net thermodynamic effects confined to the ground-state potential energy surface. Reaction endothermicities were compared to the amount of supplied energy (equivalent to λ > 127 nm). What follows is that one cannot rule out the dissociation of IC₃N into either I+C₃N, ICN+C₂ or IC₂+CN, while the heterolytic formation of ionic pairs, *e.g.* I⁺+ C₃N⁻, is not likely (Table S6). This simplistic approach does not allow us to conclude on the actual feasibility of the said homolytic cleavages, as it neglects the involvement of excited electronic states, possible energy barriers, and, of importance in cryogenic matrices, the *cage effect*.⁶⁸ The latter often curbs any durable separation of photofragments, giving preference to the stoichiometry-preserving molecular rearrangements.

Additionally, the possible formation of noble gas compounds IArC₃N and IC₂ArCN has been computationally examined. The respective theoretical predictions are provided in Fig. S4 and Table S7 of Supplementary materials (present experiments gave no supportive evidence for these species).

Photolysis results and discussion

Far-UV photolysis of Ar matrix-isolated IC₃N led to the appearance of new bands in the IR absorption spectrum (see Fig. 4). The evolution of these was monitored over the course of the experiment. Thermal cycling (annealing) steps, involving the sample warm-up to approx. 22 K for about 6 minutes, followed by re-cooling to 15 K, was performed after the irradiation.





In the case of IC₃N, the observed IR band intensity pattern was correctly reproduced by our DFT calculations (see Table 2). We assume a similar reliability of this theoretical approach applied to other closed-shell IC₃N-stoichiometry species.

Data gathered in Table 3 concern the identification of IC₂NC. A prominent IR absorption feature at 2050.2 cm⁻¹ can be identified as v_2 . It is predicted for this isomer as the strongest one. The second strongest band, the one due to v_6 bending, is expected at around 270 cm⁻¹, well beyond our available spectral range. The v_3 stretching is predicted as 2 orders of magnitude less intense in IR than v_2 . Band strength data, collected throughout the photolysis, were helpful in the search for this weak feature. Correlations of integrated intensities, involving the postulated v_2 band and

the v₃ candidate bands at 1054.5 cm⁻¹ (Fig. 5a) and 1039.5 cm⁻¹ (Fig. 5b) were analyzed. Considering the extent of error bars, Fig. 5a shows (unlike Fig. 5b) quite a probable linear correlation. Common origin of the spectral features located at 2050.2 cm⁻¹ and 1054.5 cm⁻¹ is additionally suggested by their similar resistance to annealing (in contrast to the 1039.5 cm⁻¹ band, slightly decreasing upon thermal treatment). Noteworthy, a broad hump centered at 1040 cm⁻¹ is much too strong to be due to the v₃ mode; its integrated intensity is almost the same as for the 2050.2 cm⁻¹ band. In view of the above considerations, the sharp 1054.5 cm⁻¹ feature is interpreted here as the v₃ band of IC₂NC.



Figure 5. Correlation of the measured integrated optical density values for selected pairs of IR bands relevant to the identification of IC₂NC (*a*, *b*) and ICNC₂ (c), arising in the course of far-UV photolysis of IC₃N isolated in solid argon. Red lines represent least-squares linear fits.

The predicted frequency of the v₅ bending vibration of IC₂NC formally falls within the spectral range of the instrument used, although a S/N decrease near 400 cm⁻¹ prevented detection of the respective absorption. On the other hand, the IR spectrum is crowded with absorption features around the predicted IC₂NC v₁ frequency. The anticipated separation of v₁ and v₂ is about 200 cm⁻¹ at DFT and 220 cm⁻¹ at CCSD(T) levels of theory. One can therefore expect to find the v₁ band of IC₂NC at around 2260 ± 20 cm⁻¹. Theory describes it as even weaker than v₃, and comparable in intensity to the bands of ¹³C-isotopomers of the parent species, located in the same region (cf. Fig. 2). Without further studies, involving isotopically labeled precursor species or the application of Raman spectroscopy (indeed, the expected Raman activity and IR absorption intensity patterns significantly differ one from the other; see Table 3), it is impossible to assign any specific band to the v₁ mode of IC₂NC.

We have carried out a theoretical study on the IC₃N \rightarrow IC₂NC reaction, starting from the groundstate rearrangement, analyzed within the IRC scheme. Energies of the lowest excited singlet and triplet states were derived assuming vertical excitations from discreet IRC-obtained ground-state structures along the reaction coordinate. The result is depicted in Fig. 6. The computed groundstate activation energy (including zero-point energy contributions) amounts to 267 kJ/mol, a value similar to that reported for the analogous isomerization of HC₃N (272 kJ/mol).⁶⁹ This is well below the energies of any excited electronic state of IC₃N. When the parent molecule absorbs a vacuum-UV photon from a Kr lamp, the reaction may proceed by overcoming the barrier in an excited singlet state, with a possible crossing to the lowest triplet-state surface.



Figure 6. Ground- and selected excited-state cross-sections of IC₃N potential energy surfaces depicting the isomerization towards IC₂NC. The reaction coordinate comes from ground-state IRC calculations at the B3LYP/aug-cc-pVTZ (aug-cc-pVTZ-PP for I atom) level of theory. Molecular geometries corresponding to the two energy minima and a transition state are shown. Excited state energies were derived by TD-DFT calculations, departing from the respective ground-state geometries.

As demonstrated in Fig. 5c, spectral features observed at 2172.5 cm⁻¹ and 1934.8 cm⁻¹ are well correlated throughout the photolysis time. It was at first tempting to consider the C₃N⁻ anion as their possible common carrier. However, rather than a single 2172.5 cm⁻¹ feature, the reported spectrum of C₃N⁻ coming from HC₃N photolysis in solid Ar revealed a Fermi-resonance doublet manifesting as two bands located at 2173.0 and 2178.0 cm⁻¹.^{70,71} Secondly, a mismatch of about 10 cm⁻¹ between 1934.8 cm⁻¹ and the frequency of another HC₃N-derived anion band, 1944.0 cm⁻¹, is large, even considering some differences in local microenvironments (matrix cages) inherent to IC₃N and HC₃N experiments. Finally, the intensity ratio of 2172.5 cm⁻¹ and 1934.8 cm⁻¹ bands differs greatly from that reported for the two strongest C₃N⁻ IR absorption features produced upon the photolysis of HC₃N. The currently observed ratio matches decently, however, what is predicted for the ICNC₂ isomer, as do the frequencies of the two bands (see Table 4). In

view of the above, and considering that a similar rearrangement, $HC_3N \rightarrow HCNC_2$, has been observed for cyanoacetylene, we assign the said bands to $ICNC_2$.

According to our computations, UV photons available in this experiment are energetic enough (their spectrum extending from 9.76 eV to near-IR) to overcome all energy barriers separating IC₃N from ICNC₂. The initial step most likely involves IC₃N \rightarrow IC₂NC isomerization, described above (Fig. 6). The path then proceeds (Fig. 7a) via a cyclic isomer. The latter species appears with a large excess of energy, which likely facilitates an efficient conversion towards ICNC₂ (Fig. 7b). While the exact identity of the involved potential energy surfaces is not known, the path may lead either through singlet or triplet states of the participating molecules.



Figure 7. Ground- and selected excited-state cross-sections of IC₃N potential energy surfaces depicting a 2-step (panels *a* and *b*) isomerization path from IC₂NC to ICNC₂. The reaction coordinate comes from ground-state IRC calculations at the B3LYP/aug-cc-pVTZ (aug-cc-pVTZ-PP for I atom) level of theory. Molecular geometries corresponding to the energy minima and transition states are shown. Excited state energies were derived by TD-DFT calculations departing from the respective ground-state geometries.

Conclusions

There is experimental evidence, backed with theoretical predictions, for the formation of chain species IC₂NC and ICNC₂ upon far-UV photolysis of iodocyanoacetylene isolated in solid argon. The two isomeric species have not, to our knowledge, been observed before. IR-spectroscopic description of the parent IC₃N molecule has been appended with the assignment of additional combination bands. Several fundamental IR transitions of ¹³C- and ¹⁵N-isotopologues of IC₃N, present in their natural abundances, were identified. No IC₃N-related ions, in particular no C₃N⁻, have been detected in the course of these experiments.

Declaration of interests

□ The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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State ^a	Vertical energy eV (nm)	f
$A^{-1}\Sigma^{-1}$	4.95 (251)	0
$B^{-1}\Delta$	5.07 (245)	0
$C^{1}\Delta$	5.11 (243)	0
$D^{1}\Pi$	6.90 (180)	0.07
$E^{-1}\Sigma^+$	7.10 (175)	1.3

Table 1. Vertical excitation energies and oscillator strengths for electronic transitions of IC₃N, as computed with CAM-B3LYP/aug-cc-pVQZ.

^{*a*} States designated as the linear ones, irrespective of their equilibrium geometries, to ease the comparison with other cyanoacetylene family molecules.

	B3I VP a					$CCSD(T)^b$			Experime		
	0.0011						D(1)		frequency, cm ⁻¹		
	-	IR					Vaharm	, cm ⁻¹			
		int.									Re
Mode	Vharm	harm.	Ram v_{ahar} an m , activi cm ⁻	Vahar	I _{harm} /	Vhar					1.
symmetry	, ^c	,				m.			IC3N,	C .	int
5 5	scal	km/		km	cm ⁻	VS	VM	gas ^e	IC ₃ N:Ar		
	ed,	mol	ty^d	1	mol ⁻¹	1	CF	P2	(solid)		Ar
	cm ⁻¹	(relat									
		ive									
	226	Int.)		222		22	227	227	2270		10
νι, σ	220	(100)	1381	1	161	16	6	227	(2262)	2266.4	0
	212	21		218		21	211	211	(2202)	2123.1	Ū
ν2, σ	6	(12)	80 5	5	15	55	8	5	(21130		13
	983 ¹ (0.6)	1	1084	994	0.19	10		100 9	1031	1031.1	0.
ν3, σ		(0.6)				00	988		(1037)		8
ν4 σ	347	0	645	359	0.01	35	353	353	364	out of	
V4, U						6	555 555	(349)	range		
ν5, π	517	4	1.5	532	5	49	498 4	498	496	493.9+4	5.
		(2.3)			_	4			(483)	94.9	2
		24				30	30 3 308	304	309	out of	
ν6, π	309	(14)	341	327	15	3			(323,32	range	
						11			2sh)	and a f	
ν7, π	108	0	271	114	1.0	11	124	116	(118,13	out of	
	279					0			4)	Talige	0
$v_2 \pm v_4 \pm v_6$ Of	278									2753.8	0. 1
V1+2V7+V6	1									2717.0	
V1+V4+V7	271									2/1/.9	U.

Table 2. Comparison of computed and experimental IR absorption data for IC₃N.

	7								05
N1 N4	260		268	0.000				2619.0	0.
VITV4	9		0	8				2017.0	7
	258								
v2+v4+v7 or	1							2562 1	0.
v1+3v7	258							2302.4	05
	6							6 .	
	247		254	0.000					
V2+V4	3		234	5				2476 1	0.
or v_1+2v_7	247		5	5				2470.1	5
	8								
2.10	196		201	0.10			(2087)	2074.0	0.
2V3	8		6	0.19			(2087)	2074.0	7
			106						
2	103		3,	0.2;		, v	(958,	055 1	0.
2V5	5		108	0.12			959)	955.1	06
			4 <i>g</i>						
244245	864		800	2 ·				846.8	0.
V4+V5	004	090	10-6				040.0	13	

^a Basis sets aug-cc-pVTZ for C and N atoms, aug-cc-pVTZ-I-PP for iodine.

^b Non-augmented versions of the basis sets used in DFT calculations were applied.

 c Harmonic frequencies scaled by 0.96 (see Table S4 in the Supplementary material for an

extended list of calculated overtone and combination bands).

^{*d*} Harmonic approximation.

^e Ref. 60.

^{*f*}Ref. 6, values from Raman spectra if IR data not available, sh – shoulder band.

^g Ref. 61.

	B3I VD^a					$CCSD(T)^b$			Experiment	
Mada	DJLII								IC ₃ N:Ar	matrix
symmetr	v_{harm} ,	IR int.	Raman	v_{aharm}	I.		$\nu_{aharm} \; , \;$	cm ⁻¹		Relativ
y	scale harm.	harm.	act. (relative)	, cm ⁻¹	harm km/mo l	v_{harm}	VSC	VMD	1	e IR
	d	(relative)				cm ⁻¹			v, ciii	intensit
	cm ⁻¹ ^c	, km/mol	d				1	2		у
			2054			229			not	
1, σ	2228	2 (1.1)	(100)	2287	2 (1.2)	5	2250	2253	observe	~
			(100)			5			d	
2, σ	2039	179	144 (7.0)	2090	168	207	2036	2033	2050.2	100
_,		(100)	· · ·		(100)	9	9			
3, σ	1038	6 (3.4)	1238	1078	3 (1.8)	105	1042	1048	1054.5	1.3
			(60)			6				
4, σ	349	0 (0)	423 (21)	358	0.4	361	358	357	out of	
					(0.2)				range	
_					0.12				not	
5, π	442	2 (1.1)	64 (3.1)	449	(0.1) 422	422	422 426	424	observe	
									d	
6, π	270	17 (9.5)	397 (19)	278	11	264	277	274	out of	
					(6.5)				range	
7, π	112	0 (0)	324 (16)	116	0.3 (.2)	113	131	116	out of	
				-		-	-	_	range	

Table 3. Comparison	of theoretically pred	licted and experiment	tal IR absorption	data for
IC2NC.				

^{*a*} Basis set aug-cc-pVTZ with aug-cc-pVTZ-PP for I atom.

^b Basis set cc-pVTZ with cc-pVTZ-PP for I atom.

^{*c*} DFT frequencies scaled by 0.96 (see Table S8 in the Supplementary material for the predicted overtone and combination bands).

^{*d*} Harmonic approximation.

	Computa	ations	Experiment				
Mode	B3LYP ^a			CCSD(T) ^b	Ar matix		
	Vharm, ^c cm ⁻¹	I _{harm} , km mol ⁻¹ (rel. int.)	Vanharm, cm ⁻¹	I _{anharm} , km mol ⁻¹ (rel. int.)	V _{harm} , cm ⁻¹	Freq, cm ⁻¹	I _{rel}
ν_1	2189	660 (100)	2231	590 (89)	2203	2172.5	71
V 2	1950	630 (96)	1988	660 (100)	1966	1934.8	100
V 3	1030	16 (2.4)	1074	14 (2.1)	1043	not observed	
ν4	357	1.4 (0.2)	378	0.7 (0.1)	366	not observed	
V 5	457	5.7 (0.9)	458	6.7 (1.0)	449	not observed	
ν ₆	198	6.8 (1.0)	177	5.8 (0.9)	185	not observed	
ν ₇	63	6.6 (1.0)	32	6.6 (1.0)	61	not observed	

Table 4. Comparison of the theoretically predicted and experimental	IR absorption data
for ICNC ₂ .	

^{*a*} Basis set aug-cc-pVTZ with aug-cc-pVTZ-PP for I atom.

^b Basis set cc-pVTZ with cc-pVTZ-PP for I atom.

^{*c*} DFT frequencies scaled with 0.96 (see Table S9 in the Supplementary material for the list of calculated overtone and combination bands).