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Organometallic Chemistry of New Carbon Materials IV. Structure and Dynamic Behavior of Group 6 Metal Tricarbonyl Complexes of Graphene and Perforated Graphene: A DFT Study

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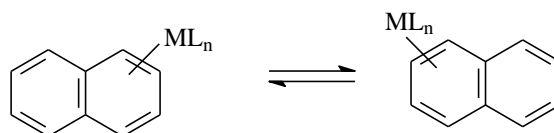
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Abstract. The mechanism of inter-ring haptotropic rearrangements (IRHR) was investigated by DFT for tricarbonyl η^6 -complexes of group 6 metals (M = Cr, Mo, W) of coronene (**I-M**), kekulene (**II-M**) and a model graphene (**III-M**). The computed η^6, η^6 -IRHR activation barriers in the *middle* size PAHs **I-M**, and **II-M** were calculated to be substantially lower than in the case of complexes of relatively *small* size PAHs such as naphthalene chromium tricarbonyl ($\Delta G^\ddagger \approx 20$ -25 kcal mol⁻¹ vs. ≈ 30 kcal mol⁻¹). The barrier is further lowered in the case of the model graphene complex **III-Cr** ($\Delta G^\ddagger \approx 13$ kcal mol⁻¹). An even lower barrier is found for **III-Mo** ($\Delta G^\ddagger \approx 10$ kcal mol⁻¹), whereas it slightly increases for **III-W** ($\Delta G^\ddagger \approx 14$ kcal mol⁻¹).

Keywords: density functional theory, metal tricarbonyl complexes, PAHs, graphene, interring haptotropic rearrangements.

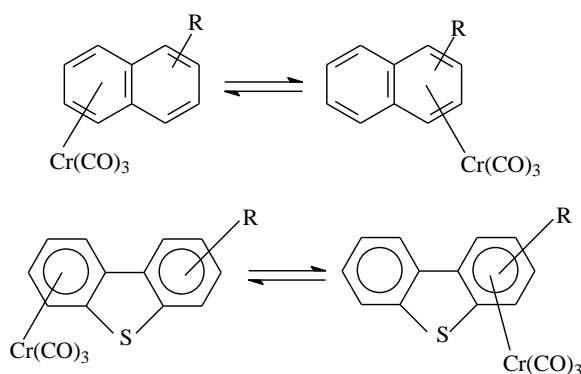
Introduction

Transition metal complexes of polycyclic aromatic hydrocarbons (PAHs) are of considerable interest due to their structural peculiarities and dynamic properties.¹ Their ability to catalyze various organic reactions provide them with many practical applications.² They are also widely used for the activation of specific positions in the coordinated PAH ring. This approach facilitates the preparation of synthetically important organic derivatives from aromatic ligands. This includes industrial production of optically pure substances of pharmaceutical utility.³⁻⁵ PAH transition metal complexes are characterized by variable fluxional behavior. Inter-ring haptotropic rearrangements (IRHRs) (η^n, η^n -IRHR, $n = 2, 4, 6$) were among the most actively studied molecular dynamic processes in the recent years. These transformations consist of the migration of an ML_n organometallic group (OMG) along the PAH molecule from one six-membered ring to another one,⁶⁻⁸ as illustrated in Scheme 1 in the case of naphthalene.



Scheme 1

The most extensively studied energy barrier associated with such IRHR processes corresponds to R-substituted η^6 -chromium tricarbonyl complexes (η^6, η^6 -IRHR). In this case, the chromium tricarbonyl group migrates along the polyaromatic carbo-^{6,7} or hetero-cycle⁸ (Scheme 2).



R = Me, D, SiMe₃, SnMe₃

Scheme 2

The experimental thermodynamic parameters of these η^6, η^6 -IRHR obtained on the basis of kinetic data are in the range of $\Delta G^\ddagger \sim 27\text{-}33 \text{ kcal mol}^{-1}$. They occur at temperatures around 90-130°C in inert non-coordinating solvents (*e.g.* decane, decaline, hexafluorobenzene). They are determined on π -complexes containing one R group on the ligand, thus eliminating isomer degeneracy. The latter can be quantitatively determined in the reaction mixture mainly using the NMR method.^{4,6-8} Factors such as the PAH and OMG natures and the presence of substituents or heteroatoms, exert a considerable effect on the direction and rate of the process.⁶⁻¹²

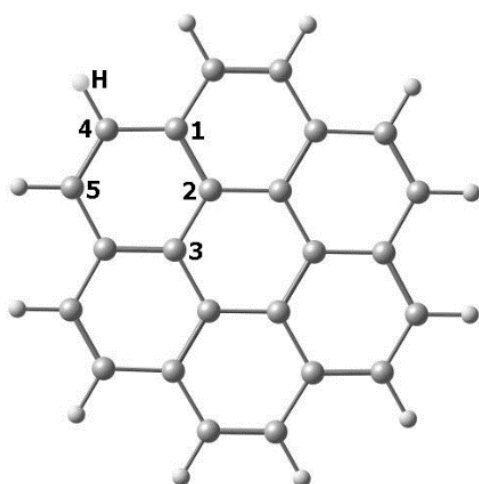
Mechanistic density functional theory (DFT) investigations of these η^6, η^6 -IRHRs have been previously performed by us, mainly for chromium tricarbonyl complexes of naphthalene, biphenyl, biphenylene, dibenzothiophene, carbazole, fluoranthene).^{10,11} Chromium tricarbonyl complexes of anthracene, phenanthrene or pyrene have been investigated by other authors.¹² For such complexes of small PAHs, the calculated activation barriers ΔG^\ddagger were all of the order of 30 kcal mol⁻¹. DFT calculations have been also performed for large PAHs, *i.e.* graphene¹³ and nanotubes.¹⁴ For all these π -complexes, DFT was able to reproduce not only the experimental structural parameters, but also their thermodynamic and η^6, η^6 -IRHR kinetic data.^{10,11,15-18} In addition, DFT allows calculating spectral parameters at a reasonable accuracy ($\sim 10\%$), such as ¹H and ¹³C NMR chemical shifts^{19,20} as well as IR vibrational frequencies of carbonyl groups.²¹ Such computed spectral parameters should facilitate compounds identification in the future.

Investigations of η^6, η^6 -IRHRs for molybdenum complexes are much rarer. For example, experimental²² and theoretical²³ rearrangement studies in heterocyclic PAH complexes of Mo(POMe₃)₃ have been performed. As far as we know, η^6, η^6 -IRHR activation barriers of tungsten complexes were neither experimentally nor theoretically studied, with the exception of a DFT investigation devoted to group 6 tricarbonyl complexes of exotic hydroxyl- and methoxysubstituted phenanthrenes.²⁴ Thus, there are no studies that allow correct comparison of experimental and theoretical data with respect to the nature of the metal within the group 6 triad. Therefore, the dependence of the mechanisms and activation energies on the metal nature remains yet unclear. Thus, theoretical modeling of the structure and dynamic behavior of group 6 transition metals complexes of PAH appears to be an interesting and urgent task. Indeed, it is directly related to many aspects of practical use of these complexes in materials science, medicine, and catalysis.¹⁻⁵

In the followings, we investigate the IRHR processes of M(CO)₃ (M = Cr, Mo, W) complexes of coronene²⁵ **I**, kekulene²⁶ **II** and the larger model **III** (C₉₆H₂₄) (Fig. 1). These fairly large PAHs can be considered as reasonable models for graphene and imperfect (defected or perforated) graphenes, as well as various new carbon materials (NCMs).^{13,14} Whereas group 6 complexes of **I** are unknown, their isoelectronic FeCp⁺ and RuCp⁺ counterparts were briefly

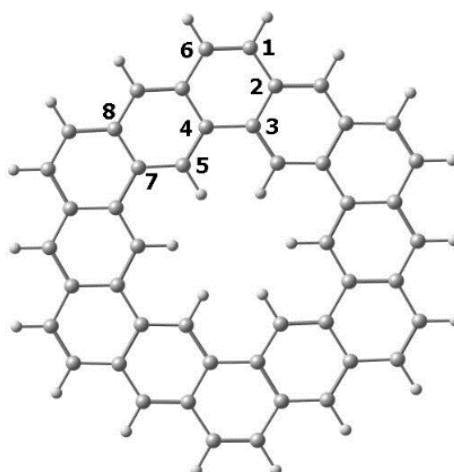
decribed.^{27,28} On the other hand, no complex of **II** has been reported so far. With respect to high molecular weight PAHs such as NCMs like fullerenes, graphene, graphite, and nanotubes, there are some reports of the synthesis and characterization of chromium complexes.²⁹⁻³¹ The Cr/C ratio in these polymetallic complexes varies in a wide range, the maximum value being Cr/C~1/18. Not only η^6 -molybdenum or η^6 -tungsten tricarbonyl complexes of such NCMs are not documented so far, but even their complexes of small PAHs are poorly studied.³²

Coordination of unligated M(0) group 6 metal atoms at the surface of NCMs and migration between six-membered rings have been investigated by electron microscopy.³³⁻³⁵ The M(0) migration has been shown to be very fast because of its weak bonding interaction with the ligand. Eventually, the M(0) atoms aggregate upon collisions to form less mobile clusters, mainly on the edges of the NCM.³⁵ On the other hand, although the unusually high rate of $\text{Cr}(\text{CO})_3$ η^6, η^6 -IRHR over a graphene surface was theoretically^{13,14} and experimentally proved,^{31,35} it is nevertheless slower than that of an M(0) atom.



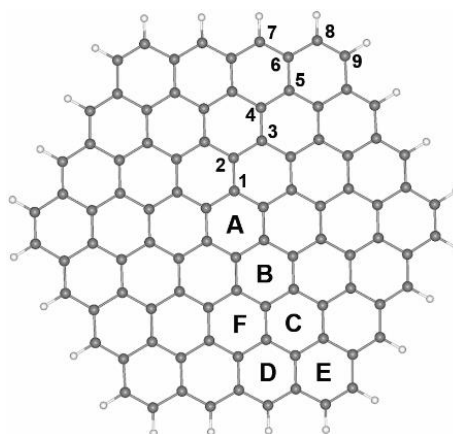
C1-C4 1.423; C1-C2 1.425; C2-C3 1.426; C4-H 1.092

I



C1-C2 1.439; C1-C6 1.366; C3-C4 1.457; C1-H 1.085

II



C1-C4 1.422; C2-C3 1.424; C3-C4 1.427; C4-C5 1.423; C5-C6 1.439;
C6-C7 1.394; C6-C8 1.440; C8-C9 1.364; C8-H 1.0924

III

Fig. 1. Optimized geometries of **I** (coronene), **II** (kekulene) and **III** ($C_{96}H_{24}$) with selected bond distances (Å).

Calculation Methods

The geometries of stable complexes, transition states, and intermediates were optimized at the DFT level, using the PRIRODA-04 program³⁶ at the MBC-100k cluster of the Joint Super Computer Center (JSCC) of the Russian Academy of Sciences (Moscow). The nonempirical nonlocal PBE functional³⁷ and scalar-relativistic theory were used, together with the **L1** extended

basis set, including Gaussian functions contracted according to the following schemes: {2,1}/{6,2} for H, {3,2,1}/{10,7,3} for C and O, {6,5,3,1}/{19,15,11,5} for Cr, {7,6,4,1}/{26,23,16,5} for Mo, and {8,7,5,2}/{30,29,20,14} for W,^{38,39} for the decomposition of one-electron wave functions to atomic orbitals. Stationary points on the PES were identified by analysis of the Hessians. In addition to the total energy (E) for stationary points on the PES, the thermodynamic functions (G) (free Gibbs energy) at 298.15 K were calculated by statistical equations of rigid rotator and harmonic oscillator. The correlation of the transition states with the corresponding minima on the PES was checked by the construction of the internal reaction coordinate (IRC).⁴⁰ The scanning procedure over one or several parameters (bond lengths, bond and torsion angles) was used to determine the transition state region. After the structure corresponding to the maximum on the scanning curve was found, the Hessian was calculated for the determination of the vibrational mode corresponding to the process with the further optimization of the saddle point. This procedure is provided in the PRIRODA program. ¹H and ¹³C NMR spectra were calculated with the GIAO (gauge including atomic orbitals)⁴¹ using L1 full-electron basis set (except for W). The calculated chemical shifts are expressed as differences between shielding of tetramethylsilane as a standard and the compound under study. Aromaticities NICS(0)/NICS(1) data were calculated for the center of the corresponding aromatic ring and for the point on the height of 1 Å (Table S4).⁴² The functional and PBE/L1 basis set were chosen due to systematic studies of the geometry of organometallic compounds and activation barriers of various processes, in particular, IRHR and by comparative calculations in the course of this work.

The interactions between the PAH and M(CO)₃ fragments were investigated also within the Morokuma-Ziegler energy decomposition analysis (EDA) framework^{43,44} through single-point calculations with the ADF program⁴⁵⁻⁴⁷ on the PRIRODA04-optimized structures, employing the PBE functional, and using the standard TZ2P basis set.⁴⁸ Within the EDA analysis, the total bonding energy (TBE) between two fragments is expressed as the sum of three components, the Pauli (or exchange) repulsion (E_{Pauli}), the electrostatic interaction energy (E_{Elstat}), and the orbital interaction energy (E_{orb}). The E_{Pauli} destabilizing component complies with the electronic antisymmetry conditions. E_{Elstat} is the electrostatic energy resulting from the superposition of the unperturbed fragment densities. The E_{orb} component originates from the relaxation of the molecular system and is associated with the mixing of occupied and unoccupied orbitals, *i.e.*, with covalency. Some ADF single-point test calculations were also performed with different (hybrid) functionals to substantiate the stability of the PRIRODA04-computed activation energies.

Results and Discussion

Coronene complexes. The structure of coronene or superbenzene **I** (C₂₄H₁₂) consists of six fused benzene rings.²⁵ It is the smallest reasonable model for graphene. Its optimized geometry as a free ligand (Fig. 1 and Table S1) is in good agreement with previous X-ray^{49,50} and electron diffraction analyses.^{51,52} The C-C distances deviate by less than 1-2% from their experimental counterparts⁴⁹⁻⁵² and are also in a good agreement with previous DFT investigations.^{53,54} Our computed ¹H and ¹³C NMR chemical shifts are also consistent with previous theoretical data^{55,56} and comparison of them with the experimental NMR and IR ν_{CO} data (Tables S3, S4) also confirm the conclusions about the aromaticity of **I** (see NICS(1) values in Tables S5). Minor deviations of the calculated chemical shifts from the experimental values can be explained by the fact that the DFT calculations were performed for the gas phase, not considering solvent effect (*e.g.* aromatic solvent induced shift ASIS-effect).⁵⁷

Calculation on the M(CO)₃ (M = Cr, Mo, W) complexes of **I** led to two optimized isomers **Ia-M** and **Ib-M** (Fig. 2 and Table S6) which differ by the nature of the six-membered ring which is complexed (outer and inner, respectively). They have a nearly planar structure of the ligand and differ significantly in energy (by more than 10 kcal mol⁻¹). As expected, complex **Ia-M**, with the metal coordinated to a more electron-rich peripheral ring, is more stable.

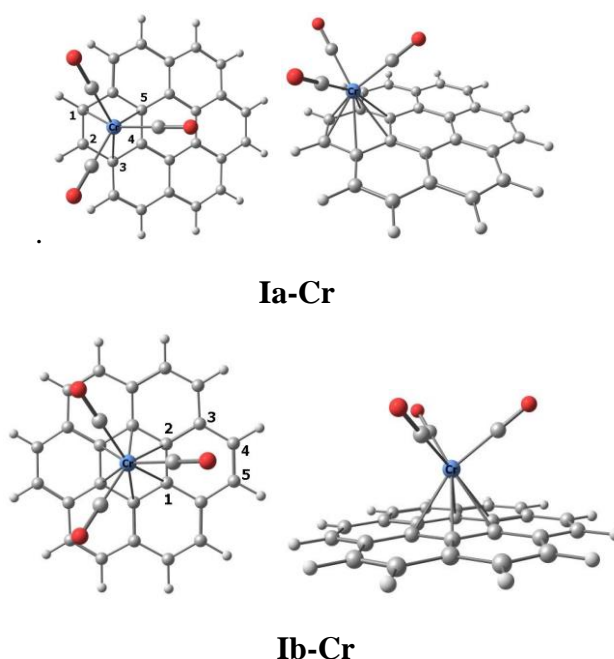
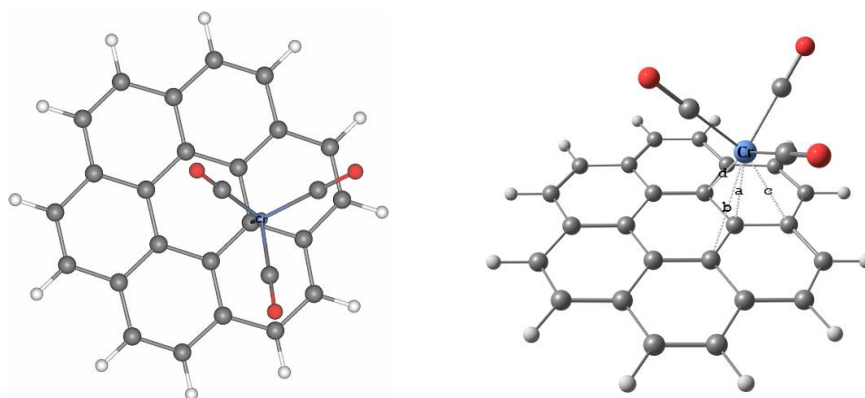


Fig. 2. Top and side views of the two isomers of Cr(CO)₃(η^6 -coronene).

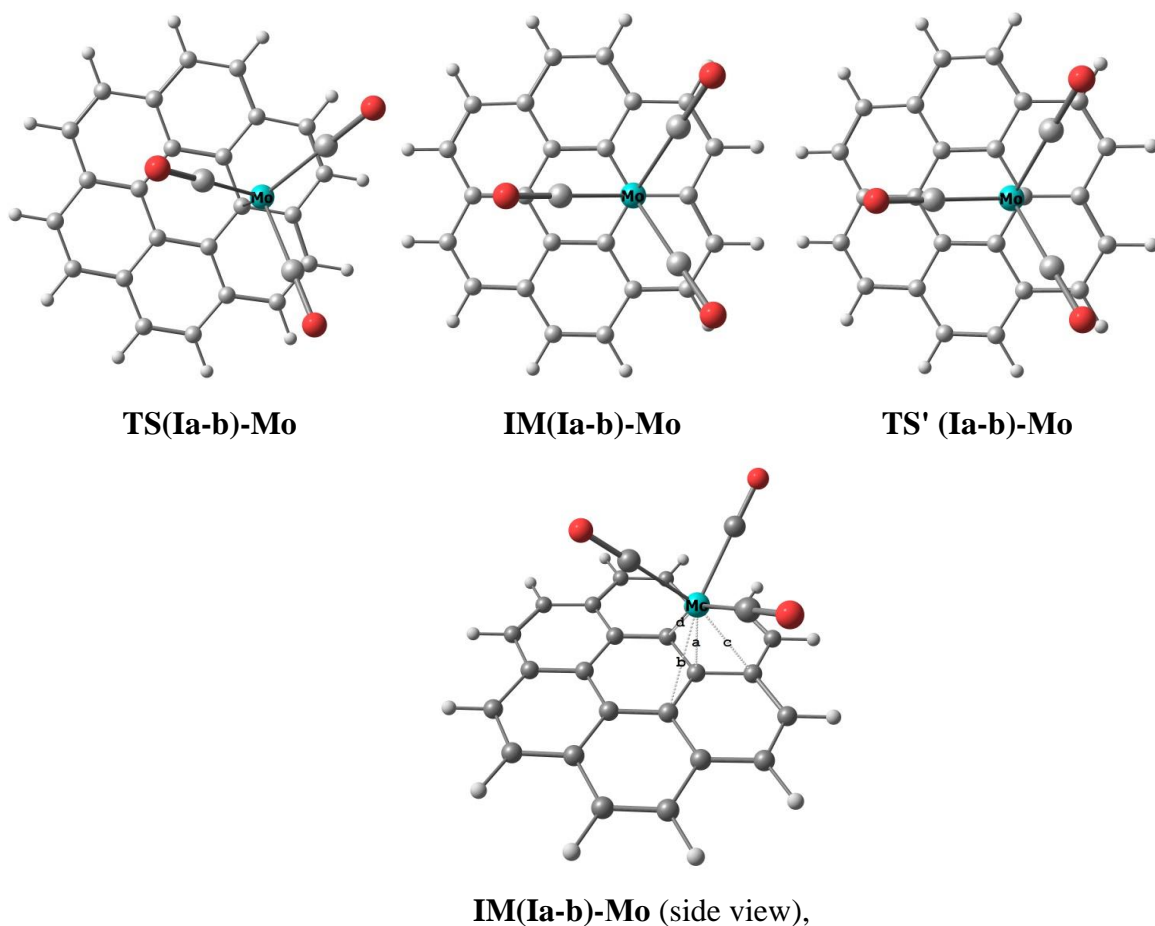
Thus, **Ia-M** should be formed predominantly in the course of thermodynamically controlled high temperature synthesis, for example, in the reaction of **I** with $M(\text{CO})_6$ ($M = \text{Cr}, \text{Mo}, \text{W}$), in dibutyl ether at 150°C . This agrees with the experimental data for the aforementioned isoelectronic RuCp^+ and FeCp^+ complexes for which the metal coordination to the central ring is not experimentally observed.^{27,28} The calculated energy difference between **Ia-Cr** and **Ib-Cr** is high ($\Delta G = 12.7 \text{ kcal mol}^{-1}$) and is close to the corresponding value obtained by *Türker et al.* for **Ia-Cr** and **Ib-Cr** ($13.2 \text{ kcal mol}^{-1}$)⁵⁸ and *Sato et. al.* for complexes of RuCp^+ ($13.52 \text{ kcal mol}^{-1}$) by DFT.⁵⁹

Two types of dynamic behaviors are possible for complexes **Ia,b-M** ($M = \text{Cr}, \text{Mo}, \text{W}$). One is a degenerate rearrangement $\text{Ia-M} \rightleftharpoons \text{Ia-M}$, *i.e.*, a shift of the OMG from one outer six-membered ring to another outer ring. The other one is the nondegenerate rearrangement $\text{Ia-M} \rightleftharpoons \text{Ib-M}$, *i.e.*, the OMG migrates between non-equivalent outer and inner rings. The later process has been studied earlier by DFT⁵⁹ for the RuCp^+ complex, and the computed activation barrier from the periphery to the center of coronene in the gas phase turned out to be very high ($\sim 40 \text{ kcal mol}^{-1}$). In our $\text{Cr}(\text{CO})_3$ complex, the non-degenerate process $\text{Ia-Cr} \rightleftharpoons \text{Ib-Cr}$ is found to involve no intermediate and proceeds through a single transition state (Scheme 3 and Fig. 3) **TS(Ia-b)-Cr** with an activation barrier $\Delta G^\ddagger = 25.1 \text{ kcal mol}^{-1}$. Hence, the backward process of Scheme 3 has a low activation energy $\Delta G^\ddagger = 12.4 \text{ kcal mol}^{-1}$, *i.e.*, the process should occur very rapidly at temperatures of kinetic measurements inherent to the η^6, η^6 -IRHR (90 – 130°C).



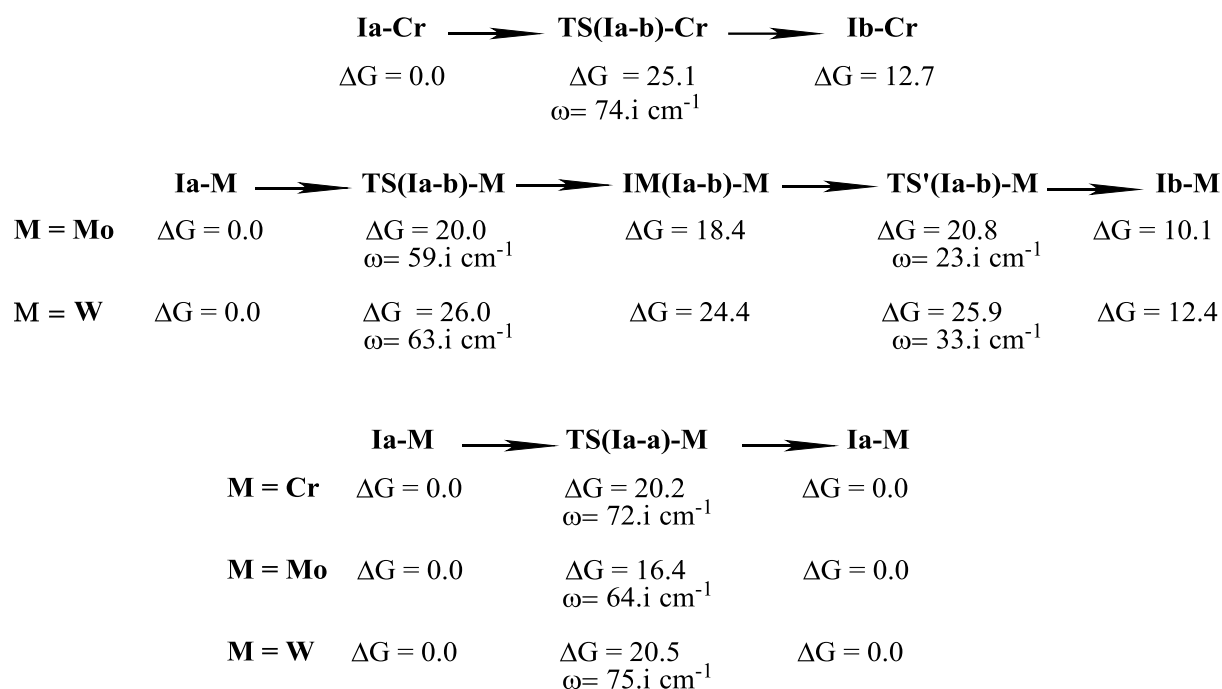
TS(Ia-b)-Cr

Bond lengths (\AA), a-d: a=2.226; b=2.784; c=2.599; d=2.511



Bond lengths (Å): a=2.373; b=2.735; c=2.842; d=2.736 (labelling in **IM(Ia-b)-Mo** and **TS' (Ia-b)-Mo** similar to that in **TS(Ia-b)-Mo**).

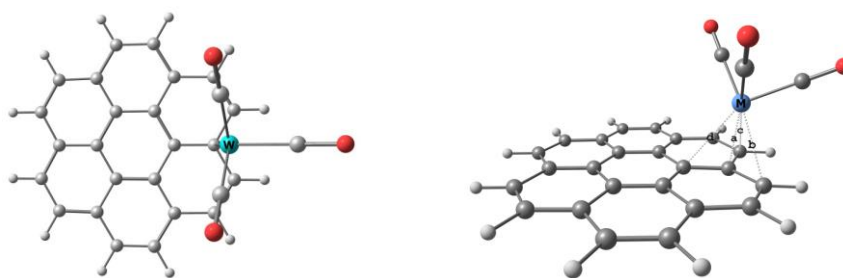
Fig. 3. Top and/or side views of the transition states and intermediates involved in the IRHR **Ia-M** \rightleftharpoons **Ib-M** (M = Cr, Mo) process. The case of M = W (not shown) is similar to that of M = Mo. All stationary points for all metals are presented in Tables S8-S10.



Scheme 3

Unlike in the chromium case, when $M = \text{Mo}$ or W the $\text{Ia-M} \rightleftharpoons \text{Ib-M}$ IRHR proceeds via a slightly different mechanism. It goes through an intermediate IM(Ia-b)-M and *via* two nearly mirror-symmetric transition states TS(Ia-b)-M and TS'(Ia-b)-M both with η^1 hapticity (Scheme 3 and Fig. 3). It is noteworthy that TS(Ia-b)-M , IM(Ia-b)-M and TS'(Ia-b)-M ($M = \text{Mo}, \text{W}$) are close in structure and energy. They are also quite close to TS(Ia-b)-Cr . During the whole rearrangement, the M(CO)_3 ($M = \text{Cr}, \text{Mo}, \text{W}$) rotational orientation almost does not change (compare Fig. 2 and 3). It should be noted that in these coronene systems, the rotational energy changes are negligibly low.

The second IRHR process, namely the degenerate $\text{Ia-M} \rightleftharpoons \text{Ia-M}$ rearrangement between two outer equivalent six-membered rings, proceeds for all metals *via* the transition state TS(Ia-a)-M without any intermediate formation. With a peripheral η^3 coordination mode of M (Fig. 4), it reminds strongly that computed for the η^6, η^6 -IRHR in naphthalene species.¹⁰ To summarize, the activation barriers of both degenerate and non-degenerate rearrangements in I-M obey the same trend: $\Delta G^\ddagger \sim 20 \text{ kcal mol}^{-1}$ for both Cr and W and a value lower by $\sim 3\text{-}5 \text{ kcal mol}^{-1}$ for Mo (Scheme 3).



TS(Ia-a)-M

Bond lengths for **TS(Ia-a)-W** (Å): a=2.273; b=2.605; c=2.605; d=2.722 for W

Fig. 4. Top and side views of the transition state **TS(Ia-a)-W**. Similar transition states are found for Cr and Mo (Fig. 4).

In order to get a deeper understanding of the bonding within the stationary points of **I-M** a Morokuma-Ziegler energy decomposition analysis (EDA) was carried out, considering the interaction between the PAH and $M(\text{CO})_3$ fragments (see Calculation Methods). The results are exemplified below in the $M = \text{Cr}$ case. The total bonding energy (TBE) between **I** and $\text{Cr}(\text{CO})_3$ in the major stationary points is given in Table 1. TBE is expressed as the sum of three components, the Pauli repulsion (E_{Pauli}), the electrostatic interaction energy (E_{Elstat}), and the orbital interaction energy (E_{orb}). The Pauli repulsion decreases with metal connectivity. It is overbalanced by the stabilizing E_{Elstat} and E_{orb} components, of which E_{orb} is prevailing, indicating covalency predominance. Unsurprisingly, the total bonding energy variation follows that of ΔG (Scheme 3). Thus, the highest activation barriers correspond to the less bonding transition state, *i.e.* **TS(Ia-b)**. The particularly small E_{orb} component in this transition state is associated with low electron transfers between the fragments. In **TS(Ia-b)-Cr**, the electron transfers corresponding to PAH \rightarrow metal donation and metal \rightarrow PAH backdonation are 0.32 and 0.20 respectively. For comparison, they are 0.32 and 0.26 in **TS(Ia-a)-Cr** and 0.88 and 0.65 in **Ia-Cr**, respectively.

Table 1. Morokuma-Ziegler energy decomposition analysis (EDA) in stationary points of **I-Cr** and **II-Cr**, all values in eV. E_{Pauli} = Pauli repulsion; E_{elstat} = electrostatic interaction; E_{Orb} = orbital interaction. TBE = total bonding energy = $E_{\text{Pauli}} + E_{\text{elstat}} + E_{\text{Orb}}$.

	<i>I-a-Cr</i>	<i>TS(Ia-b)-Cr</i>	<i>TS(Ia-a)-Cr</i>	<i>I-b-Cr</i>	<i>II-a-Cr</i>	<i>TS(IIa-b)-Cr</i>	<i>TS'(IIa-b)-Cr</i>	<i>IM(IIa-b)-Cr</i>	<i>II-b-Cr</i>
E_{Pauli}	5.64	2.15	2.70	4.78	6.04	3.40	3.15	4.28	5.23
E_{elstat}	-3.18	-1.30	-1.74	-2.48	-3.41	-2.09	-1.95	-2.45	-2.95
E_{Orb}	-4.75	-1.91	-2.28	-4.03	-5.14	-2.85	-2.64	-3.64	-4.50
TBE	-2.29	-1.06	-1.32	-1.73	-2.51	-1.54	-1.44	-1.81	-2.22

Kekulene complexes. The structure of kekulene **II** ($\text{C}_{48}\text{H}_{24}$) consists of six fused benzene rings.⁵⁰ It constitutes an interesting simple model for perforated graphenes. Its computed structural and spectral parameters (Fig. 1 and Table S2) are in good agreement with previous experimental data.⁵¹ As for **I**, two isomers were also found for the $\text{M}(\text{CO})_3$ ($\text{M} = \text{Cr}, \text{Mo}, \text{W}$) complexes of **II** (Fig. 5). They correspond to the coordination of either the terminal six-membered ring of the phenanthrene-like fragment (**IIa-M**), or the terminal six-membered ring of the anthracene-like fragment (**IIb-M**). The former is more stable for all metals (Table S2), but the difference in energy of the isomers is not as high as in the case of **Ia,b-M**, because the two types of C_6 rings in **II** are chemically quite similar. This means in the course of syntheses both isomers can be formed simultaneously in comparable quantities.

Two possible reaction routes were found for the non-degenerate rearrangement **IIa** \rightleftharpoons **IIb**: one of them proceeds through the middle of an inner C-C bond, without any exit to the ligand periphery. This pathway is unusual and was not previously observed for similar IRHR. Indeed, it has been established as a rule for a long time⁶⁰ that the IRHR proceeds with a metal shift to the periphery of the ligand but in our case, no migration over the external periphery is observed. In this process, the migration occurs *via* an η^4 -transition state **TS(IIa-b)-M** ($\text{M} = \text{Cr}, \text{Mo}, \text{W}$) (Scheme 4 and Fig. 6). No intermediate was observed in the course of the reaction. A minor deviation from planarity of the kekulene cycle can be observed in **TS(IIa-b)-M**, allowing a weak agostic bonding interaction in this transition state.

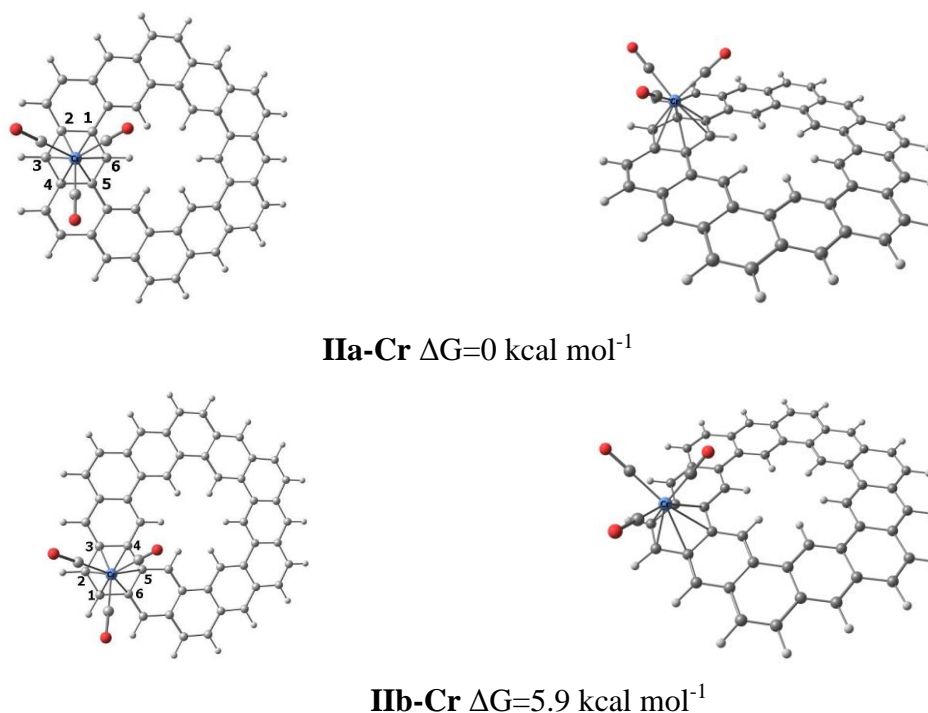
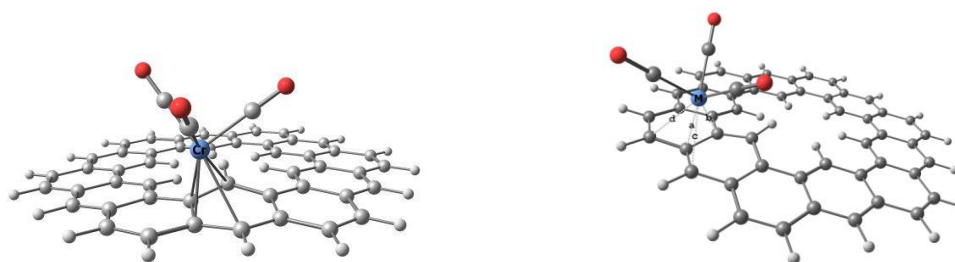


Fig. 5. Top and side views of the two isomers of Cr(CO)₃(η⁶-kekulene). Similar structures are found for Mo and W (Tables S6 and S7).



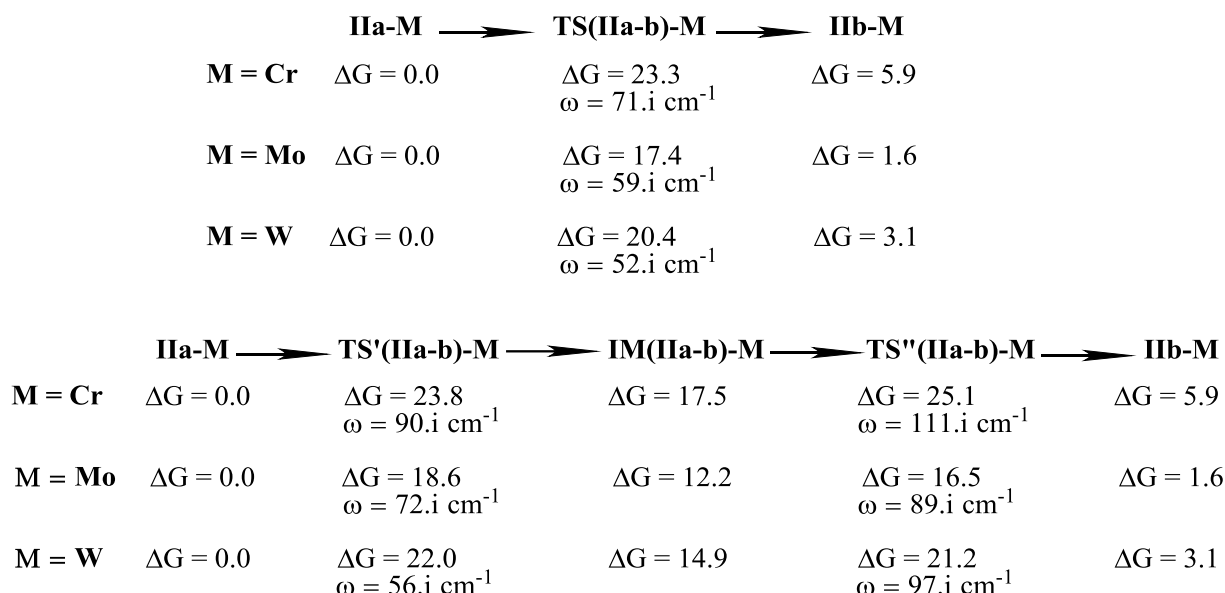
Bond lengths in **TS(IIa-b)-Cr** (Å): a = 2.159; b = 2.296; c = 2.485; d = 2.745

Fig. 6. Side and top views of the η⁴-transition state **TS(IIa-b)-M**

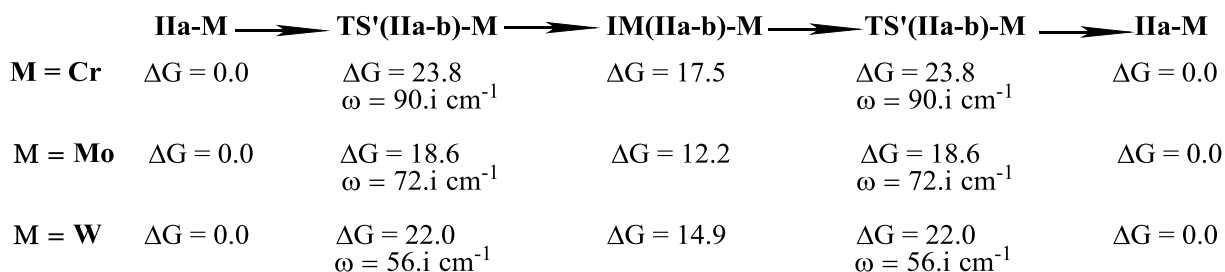
The other pathway corresponds to a usual metal shift towards the ligand periphery, but that of the ligand internal *perimeter*. An intermediate is formed, with subsequent migration to the next six-membered ring. In this process, the reaction pathway goes from **IIa-M** to the transition state η²-**TS'(IIa-b)-M**, then the intermediate η⁴-**IM(IIa-b)-M**, the transition state η²-**TS''(IIa-b)-M** and finally **IIb-M** (Scheme 4 and Fig. 7). Starting from the symmetric intermediate **IM(IIa-b)-M**, an alternative route shows up, leading to a degenerate **IIa-M** isomer through a degenerate **TS'(IIa-b)-M** transition state. This **IIa-M** ⇌ **IIa-M** rearrangement is analogous to one found earlier by DFT calculations on (η⁶-phenanthrene)Cr(CO)₃.²⁴ It should be noted that **TS'(IIa-b)-M** and **TS''(IIa-b)-M** have very similar energies for all metals and thus the IRHR energetic

picture for the degenerate rearrangement $\text{IIa-M} \rightleftharpoons \text{IIa-M}$ differs negligibly from the undegenerate $\text{IIa-M} \rightleftharpoons \text{IIb-M}$. Both $\text{IIa} \rightleftharpoons \text{IIb}$ IRHR mechanisms have similar activation energy barriers $\Delta G^\ddagger \sim 20\text{-}25 \text{ kcal mol}^{-1}$ in the case of Cr and W and 4-5 kcal mol^{-1} lower in the case of Mo. This trend within the metal triad was already found in the case of the coronene complexes.

Non-degenerate mechanisms

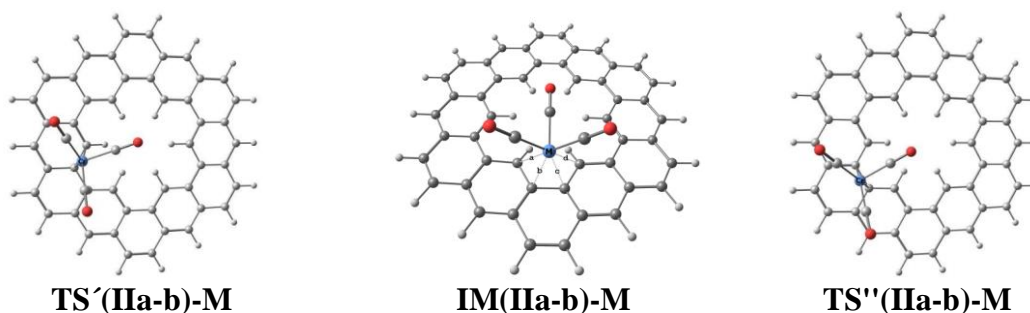


Degenerate mechanism



Scheme 4

The EDA analysis of the crucial stationary points in the case of $M = \text{Cr}$ (Table 1) provides TBE and E_{orb} values in full agreement with the computed ΔG values (Scheme 4). They also indicate stronger bonding than in the **I-Cr** system all along the reaction pathways.



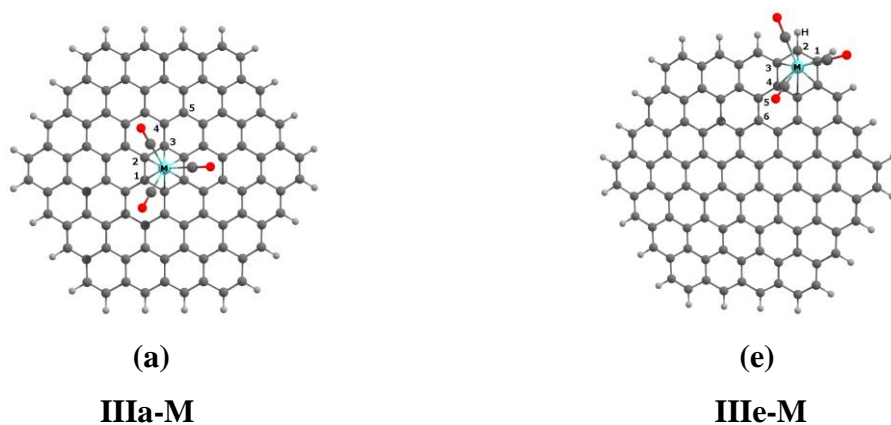
Bond lengths in **IM(IIa-b)-Cr** (Å): $a = d = 2.456$; $b = c = 2.420$

Fig. 7. The transition states and intermediate **TS(IIa-b)-M**, **IM(IIa-b)-M** and **TS''(IIa-b)-M**.

Model graphene complexes. As a reasonably large model for graphene, we chose the 4×4 structures **III** ($C_{96}H_{24}$, Fig. 1).¹³ **III** can be considered as a hexagonal planar molecule with 4-ring edges cut out from a sheet of graphene. Our computed structural data on **III** (Fig. 1 and Table 2) fit very well with neutron diffraction experiments which found the C-C distance to be 1.422 ± 0.001 Å.⁶¹ Data also agree with previous DFT by Moritz *et al.*⁶²

We now consider complexation of **III** by $M(CO)_3$ ($M = Cr, Mo, W$). There are six types of symmetry-equivalent rings available for the OMG coordination (labelled A-F in Fig.1). We investigated complexation of the two extreme positions: that of the central ring A (**IIIa-M**, Fig. 8a) and that of the more external ring E (**IIIe-M**, Fig. 8b). **IIIe-M** was found to be more stable than **IIIa-M**, consistently with the fact that the preferred localization of an OMG is always on the periphery of the PAH, rather than on an inner position. This is related to the fact that the π electron density is larger at the PAH periphery than at its center. The energy difference between both isomers is $\Delta G = 8.3, 9.2$ and 10.8 kcal mol⁻¹ for $M = Cr, Mo$ and W , respectively. Thus one can draw a conclusion that $Cr(CO)_3$ will be more uniformly dispersed on the graphene surface than Mo and W because ΔG is minimal for Cr .

The general structural features of the computed **IIIa-M** and **IIIe-M** isomers (Table. 2), are very similar to that of complexes **Ia,b-M** and **IIa,b-M**, and fit well with data from previous experimental and theoretical investigations for tricarbonyl complexes of various PAH.¹⁰⁻¹⁴ The only peculiarity of the larger **IIIa-M** and **IIIe-M** complexes is that they do not afford any noticeable out-of-plane ligand distortion, in contrast with **IIa,b-M**.



		Cr	Mo	W
IIIa	ΔE	7.8	8.9	9.8
	ΔG	8.3	9.2	10.8
IIIe	ΔE	0	0	0
	ΔG	0	0	0

Fig 8. Optimized geometries of **IIIa-Cr** and **IIIe-Cr** and their ΔG and ΔE values in kcal mol⁻¹. The M=Cr, Mo and W structures are very similar.

Table 2. Selected bond lengths (Å) for the optimized structures **IIIa-M** и **IIIe-M**

Parameter	IIIa-M			IIIe-M		
	Cr	Mo	W	Cr	Mo	W
M-CO	1.837	1.946	1.953	1.831 1.860	1.938 1.981	1.945 1.981
M-C1	2.322	2.492	2.452	2.223	2.367	2.344
M-C2	2.321	2.491	2.450	2.224	2.367	2.344
M-C3	2.322	2.493	2.452	2.368	2.531	2.498
C1-C2	1.429	1.430	1.432	1.400	1.400	1.406
C2-C3	1.436	1.438	1.440	1.443	1.448	1.449
C3-C4	1.430	1.428	1.430	1.446	1.448	1.449
C4-C5	1.417	1.417	1.416	1.431	1.428	1.431
C5-C6				1.428	1.428	1.428
C1-C2-C3 (°)	120.0	120.0	120.0	121.4	121.4	121.4

Graphene is known for its peculiar electronic properties such as electric conductance dealing with very low energy gap width. In Table 3, HOMO-LUMO gaps are presented for the ligands **III** and **I, II** (for comparison) and their computed transition metal complexes.

Table 3. HOMO-LUMO gaps (in eV) for **I**, **II**, **III**, and their metal tricarbonyl complexes.

Ligand	Uncomplexed	Cr	Mo	W
I	2.89	Ia-M -1.90	Ia-M -1.97	Ia-M -1.92
		Ib-M -1.54	Ib-M -1.63	Ib-M -1.56
II	2.41	IIa-M -1.82	IIa-M -1.77	IIa-M -1.76
		IIb-M -1.61	IIb-M -1.74	IIb-M -1.66
III	1.37	IIIa-M -0.99	IIIa-M -1.03	IIIa-M -0.99
		IIIe-M -0.86	IIIe-M -0.99	IIIe-M -0.89

From the data of Table 3, it is possible to draw following conclusions, which are important for electronics:

1. Increasing the PAH size leads to a decrease of the HOMO-LUMO gap.
2. Complexation leads to a further decrease of this gap and it is more pronounced for Cr, less pronounced for Mo and intermediate for W.
3. Materials decorated with OMGs will be more conductive when these groups will be localized on edge rather at the PAH center.

For further verification of the structure of graphene tricarbonyl complexes of the chromium triad, their computed IR ν_{CO} values are reported in the SI (Table S11). Such data can help in the future to control synthesis and production of such substances in the design of electronic materials and devices, as well as molecular switchers, sensors and machines.

η^6, η^6 -IRHR in metal tricarbonyl complexes of model graphene. The following two principal routes of the η^6, η^6 -IRHR mechanisms in **III-M** (M=Cr, Mo, W) were considered:

- a) A rearrangement consisting of the OMG shift between two inner rings of type **A** and **B**, respectively (**IIIa-M** \rightleftharpoons **IIIb-M**). These two six-membered rings are *practically* equivalent in our model where they play the role of two neighboring rings in pristine 2D graphene or very large PAH molecule. **IIIb-M** is thermodynamically slightly preferred over **IIIa-M** (Scheme 5), owing to the general tendency of the OMG to shift from center to periphery.
- b) A similar rearrangement, but involving two outer rings of type **D** and **E**, respectively (**IIIc-M** \rightleftharpoons **IIIe-M**). These two outer sites can be considered as models for practically equivalent but specific peripheral rings in very huge graphene molecules. This is supported by the very small energy difference between the **IIIc** and **IIIe** isomers (Scheme 5).

It should be noted that the unlimited number of inner rings in comparison with the limited number of edge rings in graphene flakes makes the first mechanism much more important for IRHR descriptions in graphene. The **IIIa-M** \rightleftharpoons **IIIb-M** (M = Cr, Mo, W) IRHR proceeds through an intermediate **IM(IIIa-b)-M** and via two practically equivalent transition states **TS(IIIa-b)-M** and **TS'(IIIa-b)-M** (Fig. 9, Scheme 5). The IRHR process occurs almost without any $M(CO)_3$ rotation, *i.e.* no additional energy spent for this motion.

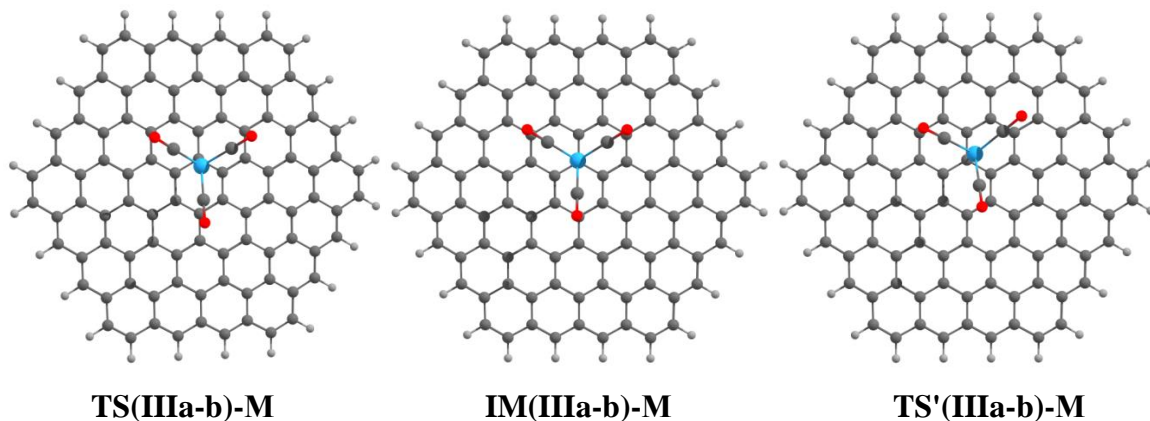


Fig. 9. The transition states and intermediates involved in the **IIIa-M** \rightleftharpoons **IIIb-M** (M = Cr, Mo, W) η^6, η^6 -IRHR

	IIIa-M	TS(IIIa-b)-M	IM(IIIa-b)-M	TS'(IIIa-b)-M	IIIb-M
M = Cr	$\Delta G = 3.5$	$\Delta G = 12.6$ $\omega = 58.i \text{ cm}^{-1}$	$\Delta G = 11.2$	$\Delta G = 12.8$ $\omega = 65.i \text{ cm}^{-1}$	$\Delta G = 0.0$
M = Mo	$\Delta G = 1.3$	$\Delta G = 10.3$ $\omega = 44.i \text{ cm}^{-1}$	$\Delta G = 9.3$	$\Delta G = 10.3$ $\omega = 50.i \text{ cm}^{-1}$	$\Delta G = 0.0$
M = W	$\Delta G = 1.1$	$\Delta G = 13.3$ $\omega = 46.i \text{ cm}^{-1}$	$\Delta G = 12.7$	$\Delta G = 14.0$ $\omega = 48.i \text{ cm}^{-1}$	$\Delta G = 0.0$

	IIIId-M	TS(IIIId-e)-M	IIIe-M
M = Cr	$\Delta G = -0.9$	$\Delta G = 14.8$ $\omega = 80.i \text{ cm}^{-1}$	$\Delta G = 0.0$
M = Mo	$\Delta G = 0.6$	$\Delta G = 12.6$ $\omega = 50.i \text{ cm}^{-1}$	$\Delta G = 0.0$
M = W	$\Delta G = 0.5$	$\Delta G = 15.3$ $\omega = 51.i \text{ cm}^{-1}$	$\Delta G = 0.0$

Scheme 5. The stationary points along the **IIIa-M** \rightleftharpoons **IIIb-M** and **IIIId-M** \rightleftharpoons **IIIe-M** (M = Cr, Mo, W) IRHR processes and their relative free energies in kcal mol⁻¹.

The second (outer) process, namely the **IIIId-M** \rightleftharpoons **IIIe-M** IRHR, proceeds for the three metals in a single step *via* **TS(IIIId-e)-M** (Fig. 10 and Scheme 5), and is quite similar to the **Ia-M** \rightleftharpoons **Ia-M** rearrangement in the coronene derivatives. The outer process **IIIId-M** \rightleftharpoons **IIIe-M** is less favorable than the **IIIa-M** \rightleftharpoons **IIIb-M** inner one, as also previously found for the coronene systems.

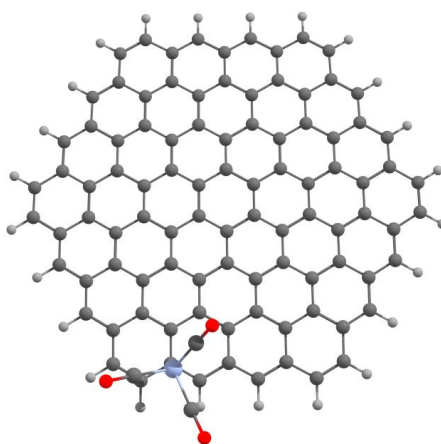


Fig. 10. The transition state **TS(IIIe,d)-M**, (M = Cr, Mo, W).

Table 4. Morokuma-Ziegler energy decomposition analysis (EDA) in stationary points of **III-Cr**, all values in eV. E_{Pauli} = Pauli repulsion; E_{elstat} = electrostatic interaction; E_{Orb} = orbital interaction. TBE = total bonding energy = $E_{\text{Pauli}} + E_{\text{elstat}} + E_{\text{Orb}}$.

	<i>IIIa</i>	<i>IM(IIIa-b)</i>	<i>IIIb</i>	<i>IIIc</i>	<i>TS(IIIc-d)</i>	<i>IIIe</i>
E_{Pauli}	4.73	2.57	4.87	5.14	3.25	4.91
E_{elstat}	-2.54	-1.55	-2.62	-2.88	-2.05	-2.84
E_{Orb}	-3.99	-2.13	-4.13	-4.46	-2.62	-4.23
TBE	-1.79	-1.11	-1.88	-2.20	-1.42	-2.16

The EDA analysis of the major stationary points in the case of M = Cr is given in Table 4. Comparing the TBE values of **IIIa-Cr** with that of **Ib-Cr** indicates similar Metal-PAH bonding strength when Cr binds to the central ring. The the PAH→metal metal→PAH electron transfers of **IIIa-Cr** (0.72 and 0.54 respectively, are also close to that in **Ib-Cr** (0.68 and 0.54, respectively). The TBE values of the high-energy stationary points **IM(IIIa-b)** and **TS(Ia-b)** are also similar. These results suggest that an IRHR process will be easier when occurring in the inner part of a graphene molecule, due to the weaker metal-graphene bonding in the η^6 -coordinated energy minima, whereas the bonding within the transition states is less affected by the inner vs. outer nature of the process.

General considerations. Comparing our computed values with related literature data available for (M = Cr, Mo, W) complexes of various PAHs,^{6-8,24} one can conclude that the activation barriers of metallotropic rearrangements decrease with the increasing of PAH size.¹¹⁻¹⁷ The available literature data for the most documented chromium tricarbonyl complexes are presented in Table 5.

Moving now down the group 6 column, it appears that all our computed IRHR energy barriers decrease systematically from chromium to molybdenum (up to 5 kcal mol⁻¹), then again increase again for tungsten (1-5 kcal mol⁻¹). A similar behavior of the group 8 triad was found earlier for DFT-calculated activation barrier in naphthalene of MCp⁺ complexes (M = Fe, Ru, Os).^{18,63} This dependence can be explained by the fact that when going down the triad, the atomic radius increases, allowing the larger metal to make bonding contacts with more distant atoms in the loosely connected transition states. In the case of the heavier metal, the f-compression effect takes place, thus reducing the above-mentioned atomic radius effect.

Table 5. DFT-computed ΔG^\ddagger activation barriers for η^6, η^6 -IRHR of chromium tricarbonyl complexes of middle-size arenes, graphene* and nanotubes* complexes are presented for comparison.

Arene	ΔG^\ddagger η^6, η^6 -IRHR	Number of different mechanisms found for η^6, η^6 -IRHR	DFT	References
Naphthalene	30.1	2	B86/SVP	[4, 15]
1-substituted naphthalenes	28.9-33.9	2	PBE/ TZ2P	[4, 15]
2 (4)-NH ₂ -byphenyl	29.9, 32.5	1	PBE/ TZ2P	[6]
Biphenylene	28.5	1	PBE/ TZ2P	[7]
dibenzothiophene	31.2	1	PBE/ TZ2P	[8]
Fluoranthene	31.3	2	PBE/ TZ2P	[11]
Phenanthrene	29.2, 30.6	1	B86/SVP	[24]
graphene*	15.4	2	PBE/ TZ2P	[13]
nanotubes*	16.2	1	PBE/ TZ2P	[14]

Finally, it should be noted that in any of our investigated IRHR processes, the rotation afforded by the M(CO)₃ group remains modest, in spite of the fact that the eclipsed vs staggered conformation can change. The rotational energy barriers in the **I-M**, **II-M** and **III-M** isomers being always low, the contribution to the activation barrier of such a small M(CO)₃ rotation in the IRHR course is negligible.

Conclusion

The structure of the $M(\text{CO})_3$ ($M = \text{Cr}, \text{Mo}, \text{W}$) complexes of coronene, kekulene and graphene (**I**, **II** and **III**) were calculated by the DFT method. The structural peculiarities of the stable isomers **Ia,b-M**, **IIa,b-M** and **IIIa,b/d,e-M** are similar for the three metals. The energy and IR ν_{CO} , ^1H and ^{13}C NMR spectral parameters of these so far uncharacterized complexes were determined. The investigation of the η^6, η^6 -IRHR dynamic process of **Ia,b-M**, **IIa,b-M** and **IIIa,b/d,e-M**, indicates that the corresponding activation barriers decrease when increasing the PAH size. It is due to the fact that the metal-PAH bonding is weaker in the case of the stable minima of large PAH complexes, whereas it remains approximately the same in the transition states and unsaturated intermediates. In the case of complex **III-Cr** the energy barriers are lower than $\sim 15 \text{ kcal mol}^{-1}$ in comparison with low size PAH. They decrease by $\sim 3\text{-}5 \text{ kcal mol}^{-1}$ when going from Cr to Mo then increases again by $\sim 3\text{-}4 \text{ kcal mol}^{-1}$ from Mo to W. Calculations on **IIa,b-M** strongly suggest that IRHR on perforated graphenes occurs preferentially around the internal rings.

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There are no conflicts of interest to declare.

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