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The complex nonstoichiometry of wüstite Fe_{1-z}O : review and comments

Jean-Raymond Gavarri¹, Claude Carel²

(version 2)

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ABSTRACT – Thermodynamic properties and structural aspects of the nonstoichiometric wüstite Fe_{1-z}O , and its modifications or allotropic varieties W_i and W'_j (i or $j = 1, 2, 3$) – the so-called pseudo phases – as function of departure z from stoichiometry, equilibrium oxygen pressure and temperature are reviewed until 2015. The complexity of the equilibrium phase diagram is analyzed in some details. The first order transition $W \rightleftharpoons W'$ is specified on the iron/wüstite boundary near 1 185 K. The phase diagram is re-examined for this transition in the temperature interval $T \in [1\ 183 - 1\ 191]$ K. The structural determination based on study of the stabilization of the point defects clusters is reviewed. The pseudo-phases transitions are re-considered concerning the transformations of defect clusters and/or their distribution modes (percolation modes/superstructures) including changes in the electronic charge carriers.

KEYWORDS: wüstite, alpha-gamma transition of iron, point defects, defect clusters, pseudo phases, percolation, superstructures.

RÉSUMÉ – Les propriétés thermodynamiques et les aspects structuraux de la wüstite Fe_{1-z}O et de ses modifications ou variétés allotropiques W_i and W'_j (i ou $j = 1, 2, 3$) – dites pseudo-phases – en fonction de l'écart à la stœchiométrie, de la pression d'oxygène d'équilibre et de la température sont l'objet d'une mise au point jusqu'en 2015. La complexité du diagram de phases à l'équilibre est analysée en détail. La transformation du premier ordre $W \rightleftharpoons W'$ est spécifiée sur la frontière fer/wüstite au voisinage de 1 185 K. Le diagramme de phase est ré-examiné pour cette transition dans l'intervalle $T \in [1\ 183 - 1\ 191]$ K. La détermination structurale basée sur l'étude de la stabilisation des amas ou 'clusters' de défauts ponctuels est également ré-examinée. Les transitions entre pseudo-phases sont re-considérées sur la base des transformations et modes de distribution des amas de défauts ponctuels (modes de percolation ou surstructures) compte tenu des changements dans les porteurs de charges électroniques.

MOTS-CLÉS : wüstite, transition alpha-gamma du fer, défauts ponctuels, amas de défauts, pseudo-phases, percolation, surstructures.

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1. INTRODUCTION

The nonstoichiometric iron monoxide $Fe_{1-x}O$, also written equivalently as Fe_yO or FeO_x after the German mineralogical appellation 'wüstite(e)'¹.

The wüstite is the main constituent of traditional blast furnace slags. Historically this oxide was the subject of numerous studies in the general context of iron metallurgy improvement, and for its role in the fields of recycling of industrial waste, catalysis and nanoparticles. It is also considered of particular interest by the geophysicists because it is a main constituent of the lower mantle of earth.

¹ There is no equivalent in the English language for the German phoneme 'ü', nor for the French one 'u'.

Because of so many papers published before, numerous reviews concerning the nonstoichiometric wüstite are available in the literature. So the present authors make choice of some of them or equivalent works ([1] to [25], period 1956 – 2015, chronological order) used independently in the series of the successive references:

Kröger, Vink [1], Wagner C [2], Goodenough [3], Kofstad [4], Vallet [5], Burgmann [6], Gokcen [7], Men' *et al.* [8], Spencer, Kubaschewski [9], Bauer, Pianelli [10], Toft Sørensen [11], Hazen, Jeanloz [12], Gleitzer, Goodenough [13], Gokcen [14], Lykasov *et al.* [15], Mrowec, Podgorecka [16], Long, Grandjean [17], Sundman [18], Wriedt [19], Collongues [20], Gleitzer [21], Smyth [22], Desré, Hodaj [23], Worral, Coley [24], Stokłosa [25].

Multiple studies have addressed the nature of the point defects (iron vacancies and interstitials) in wüstite, and the concept of defect clustering was systematically developed with a large variety of structural models. The NaCl face centered cubic structure of wüstite is characterized by a high rate of iron vacancies, and the existence of a certain proportion of iron ions occupying interstitial tetrahedral sites. Crystallographic determinations showed that these defects are not disordered in the lattice. Diverse models of clusters of iron vacancies and interstitials have been proposed.

In the present paper are reviewed observations in historical order and concepts for wüstite at equilibrium and in quenched samples. Then, it is showed how the complex nonstoichiometry of wüstite can be better described using mixtures of point defects and their associations in clusters which fall into order progressively at increasing z .

Because it simplifies the structural description of defects, the oxygen site is assumed to be fully occupied (*i.e.* there is no significant amount of oxygen vacancies). Adopting the notation of Kröger and Vink [1] the chemical formula of this monoxide can be developed as

$$\text{Fe}_{1-z}\text{O} = (\text{Fe}_{\text{Fe}})_{(1-3z)} (\text{Fe}_{\text{Fe}}^{\circ})_{(2z-t)} (\text{Fe}_{\text{i}}^{\circ\circ})_t (\text{V}_{\text{Fe}}^{q(r)})_{(z+t)} \text{O}_{\text{O}} \quad (1)$$

The ionic species Fe^{2+} and Fe^{3+} , in octahedral sites, are noted respectively Fe_{Fe} and $\text{Fe}_{\text{Fe}}^{\circ}$. The Fe^{3+} cations in tetrahedral (interstitial) sites are noted $\text{Fe}_{\text{i}}^{\circ\circ}$. Iron vacancies noted $\text{V}_{\text{Fe}}^{q(r)}$ can be charged negatively corresponding to $q = 2$ or $q = 1$, or neutral corresponding to $q = 0$. Electric neutrality would result from electrons jumping in the conduction band or interacting with holes in the valence band. It can be recalled that the band gap of wüstite was recently evaluated at about 1.0 eV at 25 °C, which means that a large population of electrons in the conduction band and/or holes in the valence band can be expected at high temperature, depending on the oxygen partial pressure [26].

Such a structure with a high concentration of defects led to introduce, in addition to the phenomenological thermodynamics which permits the detection of new (pseudo) phases Wi , the relative determination of the concentrations of simple and complex defects in connection with their free energy of formation ΔG linked to the equilibrium oxygen pressure p' . It results a new kind of diagrams relating to the defects concentrations to be related to the classic phase diagram relating to the atomic and molecular concentrations. Let us reference to Hauffe and Pfeiffer in 1953 [27], and Carl Wagner since 1940 in his book in 1962 [2] and as an example in Ref. [28], pioneers in the approach of the departure from the stoichiometry of iron oxides (See *Addendum 3* – p. 41 for A. Stokłosa, Ref. [25]: **Identification of the cluster types**).

2. THE COMPLEX BEHAVIOR OF WÜSTITE

2.1. Early work (1921-1962)

In 1921, **G. Chaudron** works in H. Le Chatelier's lab when he determines experimentally the external boundaries of the stable domain of ferrous oxide considered to be the daltonide 'FeO' [29].

Point C is at the intersection of the two boundaries, where three solid phases (Fe, Fe₃O₄ and 'FeO') coexist at equilibrium at 570 °C, characteristic temperature named Chaudron's point, that is revised many times, particularly by Raccah, Vallet, Carel at the value 591,6 °C ([30] [31] [32] [33]).

When measuring the cell parameter and the density, **Jette and Foote** demonstrate that wüstite is an iron deficient solid solution. They deduct that for the ionic compound, a missing Fe²⁺ is necessarily compensated by the oxidation of two other Fe²⁺ giving rise to electron holes Fe³⁺ [34].

Darken and Gurry study in detail the phase diagram when determining chemically the composition of wüstite after slow quenching from Θ °C = 1 100, 1 200, 1 300 and 1 350 extrapolated to 1 400. The boundaries with iron and magnetite are determined by extrapolation down to Chaudron's point. Additionally, the authors determine the molar partial and integral thermodynamic properties of wüstite (Ref. [35] parts I and II).

By comparison of calculated and experimental entropy values, **Todd and Bonnickson**, and **Humphrey, King and Kelley** deduct that the lacunary iron lattice envisaged by Jette and Foote [34] is ordered ([36] [37]).

Brynstad and Flood observe that some (electron holes) Fe³⁺ are located in tetrahedral sites written as [Fe^{III}O₄] [38].

Using neutron diffraction analyses, **Roth** determine the proportion of Fe³⁺ occupying the tetrahedral sites of the fcc structure and conclude to the formation of clusters (2:1) constituted of two iron vacancies and one tetrahedral site occupied by a cation Fe³⁺ [39]. Tarte *et al.* show directly by I.R. spectroscopy the existence of these interstitial Fe_i^{ooo} ions [40].

Using the formalism of the mass action law initially proposed by Brynstad and Flood [38], **Carl Wagner et al.** [2], Hauffe et Pfeiffer [27], and subsequently **Smyth** [22] [41] contribute to establish a simplified expression linking the composition z to the equilibrium oxygen pressure p' and the defect concentrations s following

$$z \propto p'^{1/5} \quad (2)$$

In a first attempt Smyth characterizes the defect structure of both Mn_{1-z}O and Fe_{1-z}O based on the experimental compositional data from Darken and Gurry [35]. His work often referenced has been continued by many authors in different areas of the defect chemistry. Let us reference particularly the study of the clustering (See **Stokklosa** [25]: *Addendum 3* –, p. 41: New applications of the thermodynamics to equilibrium defect structure).

2.2. Modifications or pseudo-phases W₁, W₂, W₃

2.2.1. Raccah-Vallet's team

During years 1962-65, three supposed « allotropic varieties » or « modifications » noted W_i (i = 1,2,3) named later 'subphases' then 'pseudo phases' are discovered by **Raccah** (doctoral thesis [42]) and Vallet *et al.* ([43] [47]), additionally with Kléman in [47], from the Raccah's experimental measurements. In 1964, an initial layout of an historical phase diagram is published, based on coordinates (x , Θ °C) [30].

During their studies from 1961 to 1989, **Vallet et al.** establish correlations between mass variation, *i.e.*, departure from stoichiometry expressed by x in FeO _{x} , temperature T and oxygen partial pressure p' under equilibrium (see also Ref. [5] pp. 277-80, p. 294, p. 295). They describe the phase diagram of wüstite written FeO _{x} , above 911 °C through the following three equations connecting log₁₀ p' noted

l' with composition x and temperature T , and expressing the three thermodynamic behaviors or modifications or pseudo phases W1, W2 and W3, successively [33]:

$$W1: \quad l' = (46\,753.4T^{-1} - 7.378\,1).x + (-78\,825.3T^{-1} + 16.061\,3) \quad (3a)$$

$$W2: \quad l' = (-9\,568.9T^{-1} + 31.172\,8).x + (-18\,413.3T^{-1} - 25.256\,9) \quad (3b)$$

$$W3: \quad l' = (-33\,238.9T^{-1} + 48.366\,9).x + (6\,883.9T^{-1} - 43.566\,9) \quad (3c)$$

To relate better the non-stoichiometry and equilibrium pressure p' to the presence of defects (mainly cation vacancies and interstitials), we express l' as a function of z instead of x and we report l' as a function of $\log z$ values. It can be recalled that these relationships derive from the general expression of equilibrium constants $K(T)$ linking the composition z to the oxygen partial pressure p' , following the basic setup:



$$\text{with } [h^{\circ}] = q.[V_{Fe}^{q(\cdot)}] = q.z,$$

$$K(T) = \{[O_O].[V_{Fe}^{q(\cdot)}].[h^{\circ}]^q\} / (p')^{1/2} = C.(z^s / p')^{1/2} \quad (4b)$$

$$\rightarrow \log z = (1/s).\log(K(T)/C)^2 + (1/s) l' \quad (4c)$$

where C depends on q [41]. For iron vacancies expressed V_{Fe} , exponent $q(\cdot)$ indicates a charge q times minus; h° corresponds to an electron hole. In the expression of $K(T)$, we have neglected the activity coefficients, considering that they are slowly varying with z in a limited composition range. The activity $[O_O]$ of oxygen on its site fcc is considered equal to 1 because no significant amount of vacancies can be present in the oxygen lattice. Finally, we have taken into account the relationship $-q.[V_{Fe}^{q(\cdot)}] + [h^{\circ}] = 0$ (electroneutrality condition) with $[V_{Fe}^{q(\cdot)}] = z$.

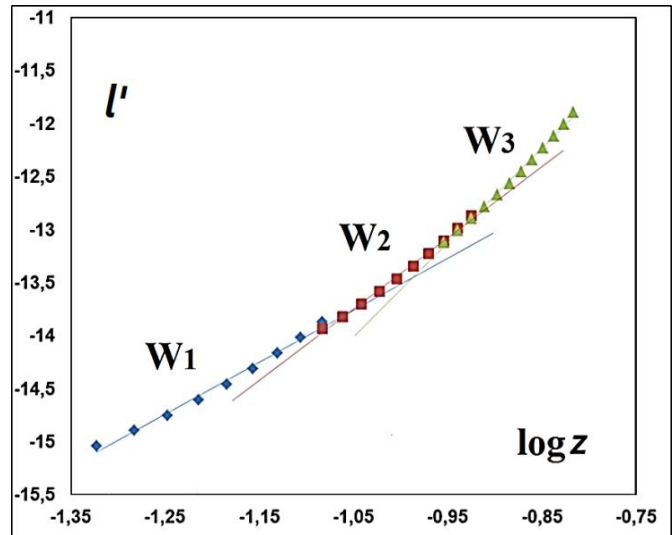
The equilibrium constants depend on the values 2, 1, and 0 of q . The three hypothetical corresponding values of exponent s should be 6 ($q=2$), 4 ($q=1$), and 2 ($q=0$) (more likely close to these values, taking into account the various approximations made). In the remainder of this review, exponent s is taken as characteristic of the specific equilibrium equation connecting the fractions of point defects between them.

Let us note that the relationships between activities and compositions z are not strictly linear. They assume the existence of three domains and allow a comparison with classical models of defect equilibria. In each domain, a unique model of defect equilibrium is supposed to be valid, which is not strictly exact because of the continuous evolution of interactions between defects, as composition and temperature vary.

2.2.2. Fixed temperature and variable composition

To illustrate the different behaviors or transitions between one and another in Fig. 1 as shown opposite, we detail the relation $\log z$ vs l' (see relation 4c) in the specific case of the Vallet *et alii*'s data as obtained at 1 000 °C ($T = 1\,273$ K).

Fig. 1. Logarithm l' of oxygen equilibrium partial pressure p' vs logarithm of composition z derived from data [33] at 1 273 K.



We use the relationships (3a, b, c) above, and transform them into the following relationships

$$W1: \quad \log z = 0.2036 l' + 1.7511 \quad \rightarrow \quad s = 4.9 (\pm 0.10) \quad (5a)$$

$$W2: \quad \log z = 0.1369 l' + 0.8352 \quad \rightarrow \quad s = 6.7 (\pm 0.12) \quad (5b)$$

$$W3: \quad \log z = 0.1124 l' + 0.5246 \quad \rightarrow \quad s = 8.9 (\pm 0.15) \quad (5c)$$

From the three segments observable on the curve l' vs $\log z$ in Fig. 1, three different values: $s(W1) = 4.9$, $s(W2) = 6.7$, and $s(W3) = 8.9$ can be derived. These values are similar to the ones found by Toft Sørensen and El Sayed Ali, and Rekas and Mrowec [48] [49]. In reality, there is no reason for assuming linear relationships given the probable existence of mixtures of defects, and of their continuous evolution conditioned by the changes in charge of $V_{Fe}^{q(t)}$ ($q=0,1,2$).

2.2.3. Fixed composition and variable temperature

For a fixed value $x = 1.130$ ($z = 0.115$), we report in Fig. 2 the experimental values of l' versus $1/T$ using the revised data of Vallet and Carel [33]. Three quasi linear domains can be distinguished suggesting the existence of three Gibbs free energies of reaction in relation with the equilibrium constant $K(T)$, with the classical relationship $\Delta G(\text{reaction}) = \Delta H - T\Delta S = -RT \ln K$. The linear correlations are:

$$W1: \quad l' = -2.5994(10^4 \cdot T^{-1}) + 7.724 \quad (6a)$$

$$W2: \quad l' = 2.9226(10^4 \cdot T^{-1}) + 9.9684 \quad (6b)$$

$$W3: \quad l' = -3.0676(10^4 \cdot T^{-1}) + 11.088 \quad (6c)$$

Using the simplified relationship (4b) without activity coefficients, we can derive the following relations for the fixed value $z = 0.115$ following:

$$\log K(T) = (s/2) \cdot \log z - \frac{1}{2} l' + \log C = - (2/2.3026) \Delta G/RT \quad (7a) \quad \rightarrow \quad dl'/d(1/T) = (2/2.3026) \Delta H/R \quad (7b)$$

From relation (7b), three values of enthalpies can be determined, linked to the changes in slopes and attributable to the pseudo-phases:

$$W1: \quad \Delta H_1/R = -25\,526 \text{ K} \quad \rightarrow \quad \Delta H_1 = -2.20 \text{ eV}$$

$$W2: \quad \Delta H_2/R = -29\,322 \text{ K} \quad \rightarrow \quad \Delta H_2 = -2.52 \text{ eV}$$

$$W3: \quad \Delta H_3/R = -31\,008 \text{ K} \quad \rightarrow \quad \Delta H_3 = -2.67 \text{ eV}$$

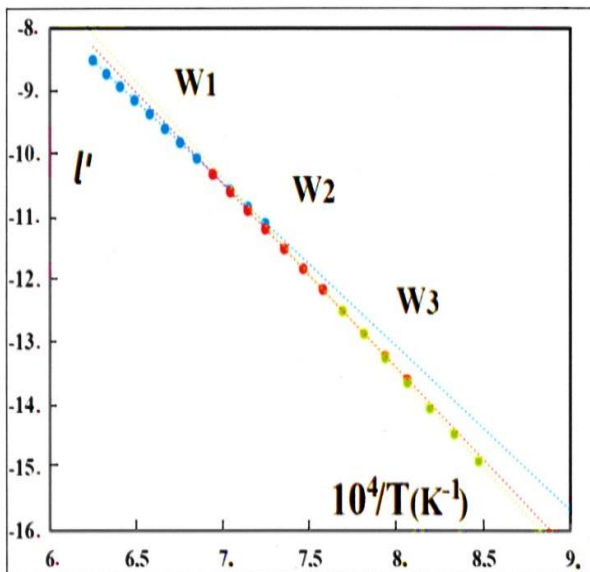


Fig. 2. Variation of l' vs $1/T$ for a fixed value $z = 0.115$: three quasi linear correlations corresponding to the initial determinations in the three domains W1 (high T), W2 (intermediate T) and W3 (low T) give rise to three activation energies ΔH_i .

At the same time, a qualitative dilatometric study under CO₂/CO indicates weak changes of slopes in curves $\ell(T)$ with transitions located on the boundary W₂/W₃ [50]. In addition, several correlations between the cubic cell parameter obtained after quenching or *in situ* and the composition x in FeO _{x} led correspond to the separation in three subdomains of three pseudo phases W₁, W₂ and W₃ in the equilibrium phase diagram ([31] [51] [52]).

In the equations above, we have neglected the activity coefficients, assuming that the ionic interactions remained constant as z varies. Another point of view would be to consider that a significant evolution of the activity coefficients can be observed and that the variation is associated with molar fractions $[i]$ of point defects at equilibrium (the activity a_i being proportional to the molar fraction: $a_i = \gamma_i x[i]$). In the case of ionic solids, these activity coefficients γ_i are generally used to express deviations from ideal models of isolated defects at equilibrium.

2.3. Abnormal behavior near (1 184 -5) K

2.3.1. A splitting of the phase diagram

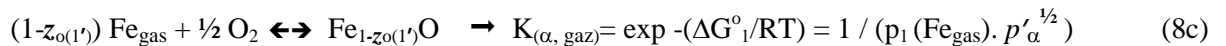
Vallet proposes a separation of an upper subdomain W of the phase diagram from the lower subdomain W' on each part of the first order transformation α -Fe \rightleftharpoons γ -Fe [53]. Each subdomain is constituted of three pseudo-phases W_i above and W'_j below 1 184 K. Because the properties of W' and W are distinct, the boundaries respectively with α - or γ -iron of the stability domains are also distinct. That is the reason for the existence of points A and B, at the lower and upper end respectively of the corresponding curves in the phase diagram (see [32], lines arbitrarily at 1 184.15 K, Tables IV et V respectively p. 723 et 724 in *French paper* or p. 19 and 20 in *English translated paper*). The limit compositions on these boundaries $y_{o(1)} = 0.942\ 7$ or $z_{o(1)} = 0.057\ 3$, and $y'_{o(1')} = 0.952\ 8$ or $z_{o(1')} = 0.047\ 2$ respectively of W₁ and W'₁ in equilibrium with γ Fe-O _{ϵ} and α Fe-O _{ϵ'} are proposed at points A and B, respectively, both assumed at 1 184 K, which cannot strictly occur as discussed below (see eq. 8 a-c and 9 a-c).

2.3.2. Thermodynamic inconsistency

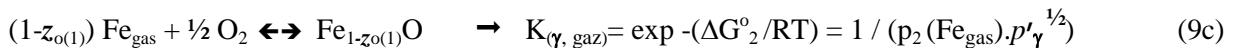
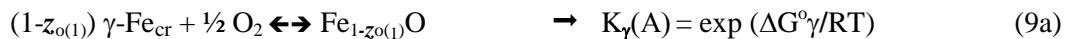
In an extended and careful review published in 1991, **Wriedt** [19] asks questions about the isotherm at 1 184 -5 K and its assumed continuity, while it seems to correspond with transitions W_i/W'_j, *i.e.*, with five sub-boundaries for the indexes 1/1, 2/1, 2/2, 2/3 and 3/3. Are they linked to the $\alpha \rightleftharpoons \gamma$ Fe transformation? These pending questions have no answer until the 1980s.

The basic equations governing the various equilibria between gases and solid phases (see equilibrium diagram in Fig. 3) are listed below. Particularly at the two boundaries involving α -Fe_{cr} or γ -Fe_{cr} (subscript 'cr' for 'crystallized solid'), equilibria can be expressed:

1) on the boundary α -Fe/W'₁ from point B to point C,



2) similarly, on the boundary γ -Fe/W₁ from point A to melting,



At the arbitrarily unique temperature of 1 184 K, it is thus necessary to apply two different partial pressures p_1 and p_2 of Fe(gas), at equilibrium with the solid at points B and A. Obviously, at the unique $\alpha \rightleftharpoons \gamma$ Fe equilibrium temperature, it exists only one value of the pressure $p(\text{Fe}_{\text{gas}})$, which is $= 3.005 \times 10^{-11}$ atm. Thus points A and B cannot be at the same temperature. Correlatively, in the case of points A and B differing in temperatures $T(A)$ and $T(B)$, the compositions $z_o(A)$ and $z_o(B)$, the equilibrium partial pressures $p'(A)$ and $p_2(A)$, $p'(B)$ and $p_1(B)$ of oxygen and iron gas are necessarily different.

Finally, a small difference in temperatures, probably a few degrees, remains to be determined experimentally with high accuracy in order to adjust this set of parameters, which determines the phase diagram (Fig. 3) close to the stable temperature of the $\alpha \rightleftharpoons \gamma$ iron transition. In this situation, the bottom row in Table IV should be changed to a value of $T(A)$ lower than 1 191 K, and the header line of Table V to a value of $T(B)$ also higher than 1 184 -5 K, but lower than $T(A)$ in (Ref. [32] p. 723 and 724 or p.19 and 20 in English translation in [hal-03083695](#)).

Note: If $T(A)$ and $T(B)$ differ from a typical value of 4 °C (See § 2.3.4. *Fe-O phase diagram near to 1 185 K*, Fig. 3), for example $T(A) = 1\,191$ K and $T(B) = 1\,187$ K, the equilibrium partial pressures of iron are necessarily $p(A) = p_1 = 3.692 \times 10^{-11}$ atm and $p(B) = p_2 = 3.220 \times 10^{-11}$ atm following tabulated data [54], [55], [56]. Thus, the difference between the two equilibrium partial pressures of iron gas is on the order of 0.5×10^{-11} atm. The corresponding equilibrium partial pressures of oxygen are $p'(A) = 4.829 \times 10^{-17}$ atm and $p'(B) = 4.229 \times 10^{-17}$ atm, i.e. a difference of 0.6×10^{-17} atm. The compositions $z_o(A) = 0.057$ 1 and $z_o(B) = 0.046$ 8 correspond to a gap of near to 0.010 in wüstite chemical composition.

2.3.3. Phenomena observed close to the isotherm at 1 184-5 K

Experimental confirmations of this hypothetical separation exist in the literature. Several of them are interesting to be examined. Löhberg and Stannek separate the pre-exponential term a and the energy term $b = \Delta G/RT$ of an expression functionally similar to a Boltzmann distribution $p_{\text{H}_2\text{O}}/p_{\text{H}_2} = a \cdot \exp(b \cdot z)$ expressed as a function of $1/T$ in two segments intersecting in the vicinity of the ($\alpha \rightleftharpoons \gamma$) iron transition temperature in (Ref. [57] Fig. 8 and 9 p. 250). Similarly, Janowski, Mrowec, Stokłosa separate in two segments the variation of the self-diffusion coefficient D^*_{Fe} of iron in wüstite as a function of $1/T$ from their own results and several sets from the literature (see Ref. [58] Fig. 7 p. 99). The intersection occurs for $T^{-1} = 8.28 \times 10^{-4} \text{ K}^{-1}$, i.e., at 935 °C, of the segments attributed to W and W'.

In years between 1981 and 1993, Jacobsson and Rosén, Guillermet and Per Gustafson, Sjöden *et al.*, Grønvold *et al.* publish new experimental thermodynamic data of highly improved accuracy about the molar thermodynamic properties of iron and wüstite ([59] [60] [61] [62]). Sundman, in an entire assessment of the Fe-O system, formulates a model for the ‘pure iron corner’ (solid solutions $\gamma\text{Fe-O}_e$ and $\alpha\text{Fe-O}_e$) [18]. The assessment of Spencer and Kubaschewski is also relevant in the same context [9].

A departure from a monotonous variation of the electromotive force of galvanic cells, near the $\alpha \rightleftharpoons \gamma$ Fe transition, is evidenced from measurements at the iron/wüstite electrode, in several galvanic cells [59] [61]. In addition, authors [61] publish their numerous determinations of the electromotive force between 866.7 and 1 339.3 K. An anharmonic transitional zone between two separate continuous variations, not in continuation of each other, can be identified in the temperature range [1 186 - 1 194] K (See in Ref. [32] [hal-03083695](#) (2021) English, Thermodynamic Properties of the W_i and W'_i from Thermogravimetric Data at Equilibrium, p. 18-23: Tables III to VIII, *ADDENDA*, 3 –The transition $W'_1 \rightleftharpoons W_1$ at 1 185 K: boundaries with α - and γ -Fe).

2.3.4. Consequences relevant to the Fe-O phase diagram

The $\alpha \rightleftharpoons \gamma$ Fe first order transition in the ‘pure iron corner’ of the phase diagram Fe-O is not documented sufficiently. The left side in Fig. 3 displays literature data (relative to FeO_x) concerning the domain of this transition ([9] [18] [63]). The temperature of the strictly pure $\alpha\text{Fe} \rightleftharpoons \gamma\text{Fe}$ transition is tabulated at $1\,184 \pm 3$ K in the data base NIST-Janaf (1999-2011) [56]. At present, it is generally taken to be 1 185 K (point Ø in Fig. 3).

The $\alpha \rightleftharpoons \gamma$ Fe transitional enthalpy is known with a twenty percent dispersion from assessed and experimental data (see Ref. [60] Table 3 p. 604). Very few experimental studies are available concerning the P(eritectoid) invariant point (α Fe- O_ε , W'_1 , γ Fe- O_ε) close to the temperature of the pure α -Fe \rightleftharpoons γ -Fe transition [63]. Sundman (see Ref. [18]) takes up the molar fractions $\cong \varepsilon$ at point b, $\cong \varepsilon'$ at point P, as displayed in Fig. 3, of dissolved oxygen 6.8×10^{-6} and $N_O = x/(1+x) = 16 \times 10^{-6}$ in γ - and α -Fe respectively at 1 185 K or 912 °C. The separation of the boundary iron/wüstite into parts γ Fe- O_ε / W_1 from point A and α Fe- O_ε / W'_1 from point B arises distinctly in the temperature range [912-921] °C characterized thanks to authors [61], following the suggested phase diagram in the right part of Fig. 3.

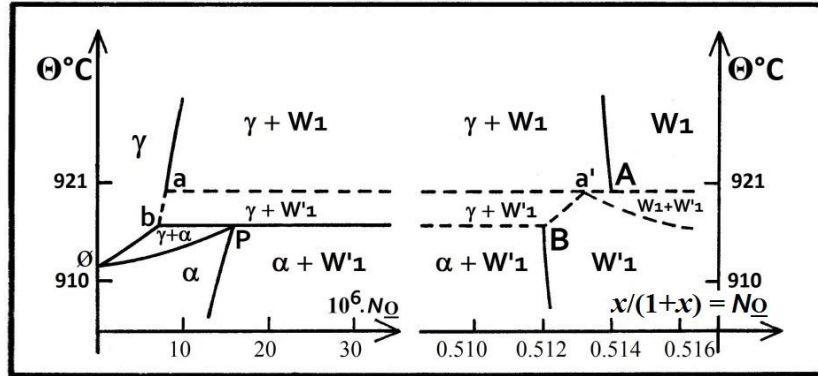


Fig. 3. Equilibrium diagram (FeO_x) between W_1 and W'_1 and ‘pure iron (point Ø only)’, with both limit points A and B of γ -Fe/ W_1 and α -Fe/ W'_1 boundaries [33]. Dashed lines correspond with hypothetical limits.

Two types of two-phase domains ($\gamma + W'_1$) and ($W_1 + W'_1$) have been added hypothetically as well as an invariant point a' (W'_1 , γ Fe- O_ε , W_1). As of now, the location of points A and B remains to be determined more precisely. New experimental data is necessary to confirm the preceding suggestions.

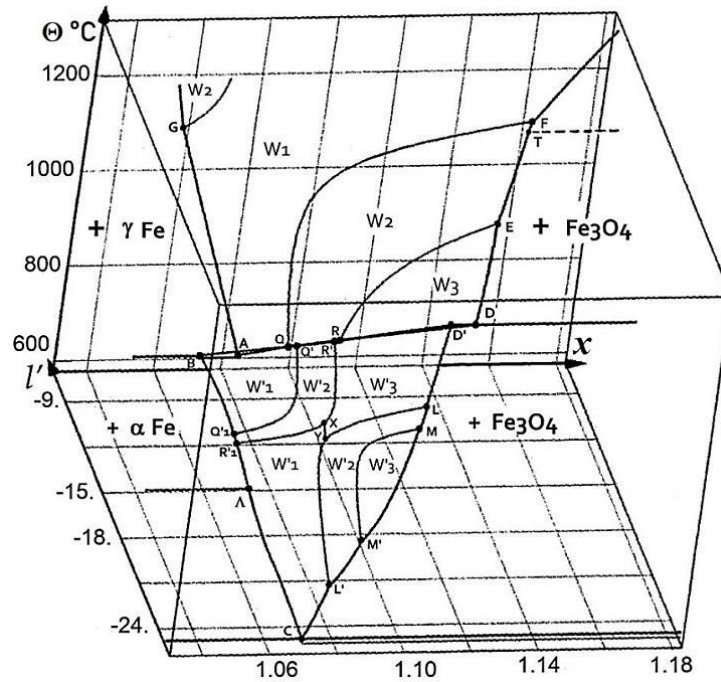


Fig. 4. Phase diagram Θ °C (l' , x) of solid wüstite FeO_x under equilibrium, following Raccach-Vallet-Carel; upper temperature = 1 250 °C, lower one = 592 °C at point C; coordinate l' for $\log_{10} p(O_2)$. Points noted by a capital letter are tabulated in [31], except point Λ in [33] and point T in ([64] [65] [66]).

The sets of numerical relations detailed above for l' as functions of x and T^{-1} allow the determination of the transition lines or boundaries between domains of the pseudo-phases W_1 , W_2 and W_3 , from γ -iron to magnetite between 912 and 1 400 °C, and W'_1 , W'_2 and W'_3 , from α -iron to magnetite between 912 and 592 °C, at point C, as drawn in Fig. 4 (see precedent drawings in Refs. [30] [31] [32] [33]).

2.4. First experimental evidences of phase modifications or pseudo phases in the literature

2.4.1. Wagner jr's team - 1966

In the course of (thermodynamics - defect structure) approaches by Carl Wagner, **Wagner jr et alii** study the electrical conductivity σ at equilibrium as function of $\log_{10} p_{O_2}$ [67]. With the help of their own thermogravimetric data $l'(x,T)$ [68], they state that the composition of the anomalies on their curves $\sigma(x)$ well agree with the transitions $W_i \rightleftharpoons W_{i+1}$ determined in Raccach's work. So, the pseudo phases are identified at equilibrium for the first time outside of Vallet's group in the literature (see [67] Table IV p. 955; also ANNEX 1 – Fig. A1 p. 38). The further studies of Wagner jr's group about wüstite were concerned with transport phenomena in the bulk and on the surface, in relation with the p to n transition, and never more with the transitions $W_i \rightleftharpoons W_{i+1}$.

2.4.2. Fender and Riley - 1969

From numerous accurate e(lectro)m(otive)f(orce) measurements in galvanic cells, **Fender and Riley** [69] brought the second experimental confirmation, outside of Vallet's team, of the existence of three pseudo phases, and the location of the boundaries of the subdomains in the range 700-1 350 °C. Their plot of the equilibrium diagram is in fairly good agreement with that of Vallet *et al.* (see ANNEX 2 – Fig. A2 p. 39). Here again the varieties no more exist in the further papers of the group.

2.5. Partial molar enthalpies and entropies. Convergence point Ω

Fender and Riley then later Vallet and Carel assessed the partial molar properties of the oxygen solid solution ($\Delta \bar{H}_O = \bar{H}_O - \frac{1}{2} H^\circ_{O_2}$, $\Delta \bar{S}_O = \bar{S}_O - \frac{1}{2} S^\circ_{O_2}$) as functions of $x \in [1.058-1.172]$ for the three pseudo phases W_i ($i=1,2,3$) also named wI , wII and $wIII$ (see [69] Table I p. 796, and *infra* in Figs. 5 and 6).

The following relations were formulated in [32] for the three W_i ($i=1,2,3$) at $T > 185$ K

$$\Delta \bar{H}_O = 19.144\,75(a_i \cdot x + c_i) \quad (10h)$$

$$\Delta \bar{S}_O = -19.144\,75(b_i \cdot x + d_i) \quad (10s)$$

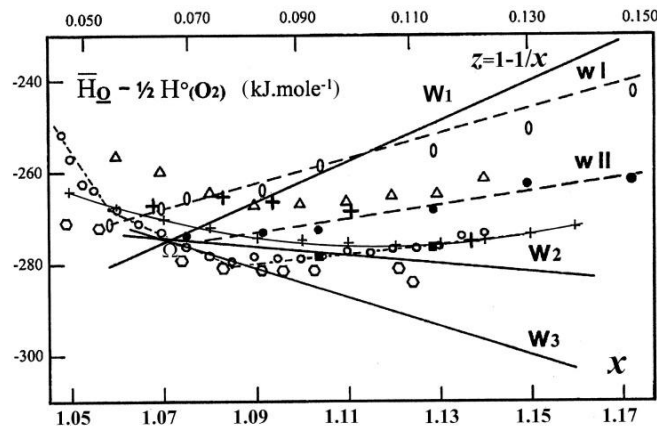


Fig. 5. $\Delta \bar{H}_O$ from some authors. —: derived for the W_i ($i=1-3$) i.e. at $T \in [1\,193 - 1\,523\text{ K}]$ from thermogravimetry + CO_2/CO [30]; $\bigcirc\text{---}\bigcirc$: wI , $\bullet\text{---}\bullet$: wII , and \blacksquare : $wIII$ (two points only) by emf at $T \in [1\,023-1\,623\text{ K}]$ [69]; dash-dot line $\bigcirc\text{---}\bigcirc$: by calorimetry at 1 348 K under vacuum (+ δn_{O_2}) [70]; $+$: by emf at $T \in [873-1\,273\text{ K}]$ [71]; Δ : by thermogravimetry + CO_2/H_2 at $T \in [1\,373-1\,673\text{ K}]$ [72]; $\bigcirc\otimes$: by emf at $T \in [823-1\,323\text{ K}]$ [73]; $+$: by emf data at $T \in [873-1\,373\text{ K}]$ ([15] [74]).

and, for the three W'_j ($j=1,2,3$) at $T < 1\ 185\ K$

$$(\Delta\bar{H}_O)' = 19.144\ 75 (a'_j x^2 + c'_j x + e'_j) \quad (11h)$$

$$(\Delta\bar{S}_O)' = -19.144\ 75 (b'_j x^2 + d'_j x + f'_j) \quad (11s)$$

Relations (10) and (11) for W_i , and W'_j respectively are independent of T , in accordance with the property of regularity (Ref. [14] p. 81-85) of the wüstite solid solution identified in [43] but contested in [70].

$\Delta\bar{H}_O$ is represented in Fig. 5 from numerous available data sets. A convergence in the vicinity of the common intersection point Ω at $x = 1.071\ 3 (\pm 0.001)$ of the straight lines $\Delta\bar{H}_O(x)$ from relation (10h) is globally noticeable.

Contrary to this contrived amalgamation process, the numerical adjustments presented next are related principally to Raccach's results ([42] [43] [44] [45] [46]) which form a well populated and coherent data set at equilibrium, for which it exists interpretations by Vallet's group ([30] [31] [32] [33]), and experimental confirmations out of the group ([67] [69] [72] [75]). (See also *infra* Refs. [85], and [25] in Addendum 3 - p. 41).

The graphical representation of the relations (10h) and (11h), (10s) and (11s), in the case of the isotherms at $1\ 000$ and $820\ ^\circ C$, is given in Fig. 6. Along the isotherm at $1\ 000\ ^\circ C$, the three W_i are represented by successive linear segments. For the two transitions $W_i \rightarrow W_{i+1}$ at decreasing iron content a jump is observed toward the lower molar heat and entropy of solution which are exothermic and exo-entropic: for $x_{1/2} = 1.078\ 2$: $\delta(\Delta\bar{H}_O) = -3.5\ kJ.mole^{-1}$, $\delta(\Delta\bar{S}_O) = -2.3\ J.K^{-1}.mole^{-1}$, and $x_{2/3} = 1.115\ 4$: $\delta(\Delta\bar{H}_O) = -10.5\ kJ.mole^{-1}$, $\delta(\Delta\bar{S}_O) = -8.3\ J.K^{-1}.mole^{-1}$ corresponding to an increase in the ordering.

Above $1\ 184\ K$, the exothermic $\Delta\bar{H}_O$ for W_1 increases with x and z corresponding to oxygen dissolution. For each O a vacancy, and zero, one or two electronic holes are created simultaneously as the clustering is in progress (see equation (4a)). For W_3 , $\Delta\bar{H}_O$ and $\Delta\bar{S}_O$ are decreasing with x , probably because of the formation of long-range ordering. For W_2 , $\Delta\bar{H}_O$ varies in a manner intermediate between those for W_1 and W_3 , which corresponds to the observation that W_2 is a solid solution of W_1 and W_3 [69].

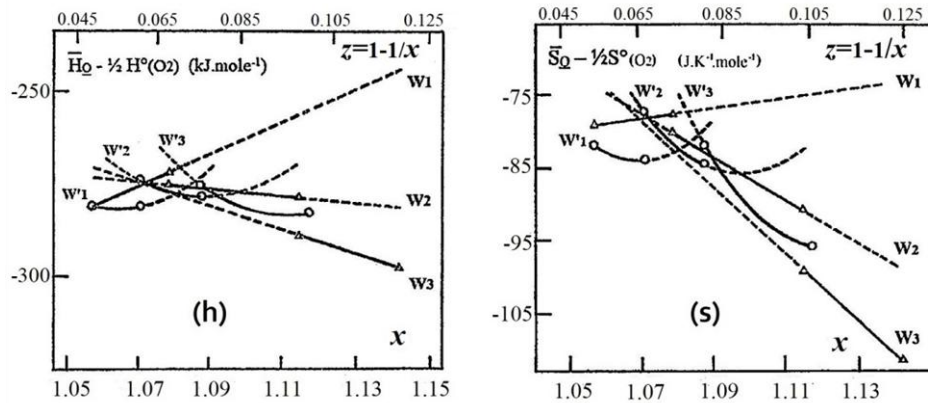


Fig. 6. Partial molar enthalpy ($\bar{H}_O - \frac{1}{2} H^0_{O_2}$) and entropy ($\bar{S}_O - \frac{1}{2} S^0_{O_2}$) of solution of oxygen, through the stability field of wüstite, whatever the temperature. Transitional jumps are for x at $820\ ^\circ C$ (W'_i) and $1\ 000\ ^\circ C$ (W_i). Dashed lines are drawn in extrapolated metastable zones.

Along the isotherm at $820\ ^\circ C$, the three W'_i are represented by parabolic curves. Jumps (weaker in amplitude than the ones at $1\ 000\ ^\circ C$) are observed toward higher values for the molar exothermic transitions $W'_j \rightarrow W'_{j+1}$, $x_{1/2} = 1.071$: $\delta(\Delta\bar{H}_O)' = +10.9\ kJ.mole^{-1}$, $\delta(\Delta\bar{S}_O)' = +6.3\ J.K^{-1}.mole^{-1}$ and $x_{2/3} = 1.115\ 4$: $\delta(\Delta\bar{H}_O)' = +2.9\ kJ.mole^{-1}$, $\delta(\Delta\bar{S}_O)' = +2.7\ J.K^{-1}.mole^{-1}$. A convergence of parabolas $\Delta\bar{H}_O'(x)$ is

observed for $x = 1.085\ 2$ or $y = 1 - z = 0.921\ 5$ (± 0.004). This convergence occurs for a composition different ($\delta y = 1.3 \cdot 10^{-2}$) from that of point C ($y = 0.935$), and is not defined as precisely for W' as it is for W .

The Gibbs-Duhem relation leads from $\Delta \bar{H}_O$ to the iron solution property $\Delta \bar{H}_{Fe}$ (Fig. 7).

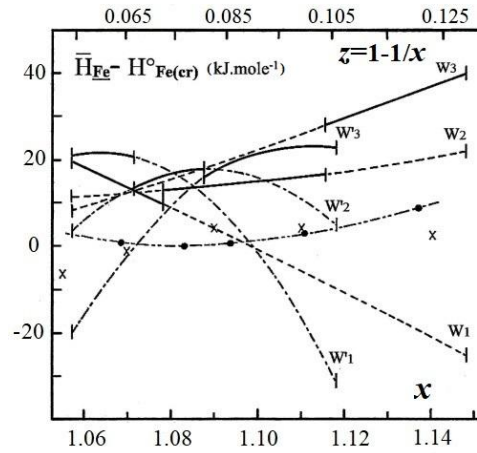


Fig. 7. Variations of the partial molar enthalpy ($\bar{H}_{Fe} - H_{Fe}^0$) vs x for the isotherms at 1 000 °C (W_i) and 820 °C (W'_j). For comparison: $\bullet\text{---}\bullet$: $\Delta \bar{H}_{Fe}$ from e.m.f. measurements at 1 000 °C by Asao *et al.* (See [71] Table 3 p. 75); \times : Darken and Gurry determinations (See [35] \bar{H}_1 Table X p. 1408).

In the domain of W , exothermic transitional jumps are observed toward the higher values for the two transitions $W_i \rightarrow W_{i+1}$ at equilibrium. At 1 000 °C, the transition W_1/W_2 ($x_{1/2} = 1.078\ 2$) is observed for $\delta(\Delta \bar{H}_{Fe})_{1/2} = 3.25\ \text{kJ.mol}^{-1}$, $\delta(\Delta \bar{S}_{Fe})_{1/2} = 2.56\ \text{J.K}^{-1}.\text{mol}^{-1}$, the transition W_2/W_3 ($x_{2/3} = 1.116\ 4$) for $\delta(\Delta \bar{H}_{Fe})_{2/3} = 12.06\ \text{kJ.mol}^{-1}$, $\delta(\Delta \bar{S}_{Fe})_{2/3} = 9.47\ \text{J.K}^{-1}.\text{mol}^{-1}$. In the domain of W' , the equivalent jumps are weaker and exothermic as shown in Fig. 7 drawn for transitions at 820 °C.

As iron deficiency increases, the partial molar enthalpy of solution of iron $\Delta \bar{H}_{Fe}$ decreases for W_1 and increases for W_3 , variations that are opposite to that of $\Delta \bar{H}_O$. The five determinations of Asao *et al.* at 1 000 °C are distributed on a curve noticeably parallel (see Fig. 7 from Ref. [71]) to the set of the successive segments concerning W_1 , W_2 and W_3 .

Equations $x_{i/i+1}(T)$ of the lines -the so-called boundaries- separating the domain of stability of single pseudo phases W_i and W_{i+1} were adjusted following the analytical equations:

$$x_{i/i+1} = [(c_{i+1} - c_i) + T.(d_{i+1} - d_i)] / [(a_i - a_{i+1}) + T.(b_i - b_{i+1})]$$

(see Ref. [30] p. 3680, equations (6)). Numerically (See Ref. [33]), they are adjusted now to

$$x_{1/2} = (60\ 412.0 - 41.318\ 2\ T) / (56\ 322.3 - 38.550\ 9\ T) \quad (12)$$

$$x_{2/3} = (25\ 297.2 - 18.310\ T) / (23\ 670.0 - 17.194\ 1\ T) \quad (13)$$

A metastable boundary can also be envisaged

$$x_{1/3} = (85\ 709.2 - 59.628\ 2\ T) / (79\ 992.3 - 55.745\ 0\ T) \quad (14)$$

The authors in Ref. [45] remarked that the corresponding terms in the numerator and the denominator [of equations (12), (13) and (14)] are almost exactly proportional in a same ratio of 1.07, *i.e.*, $(c_{i+1} - c_i) / (a_i - a_{i+1}) = (d_{i+1} - d_i) / (b_i - b_{i+1}) \cong 1.07$, also close to Chaudron's point composition. It results that the three pseudo-phases are under equilibrium ($\forall T$) under the same oxygen pressure for a composition $x \cong 1.07$ (*e.g.*, $p' \cong 8 \times 10^{-17}$, 4×10^{-15} or 8×10^{-11} atm at 1 184, 1 273, or 1 573 K respectively), irrespective of index i . This result constitutes the reason why the property $\Delta \bar{H}_O$ ($= -274.8\ \text{kJ.mol}^{-1}$) of equation (9h) is shared by the three W'_i with the same composition $x \cong 1.07$ under the same equilibrium

oxygen potential. The convergence at point Ω results of that fact. The iron partial molar properties converge also for $\Delta\bar{H}_{\text{Fe}} = +87.8 \text{ kJ.mol}^{-1}$ at 1 000 °C. At invariant point C ($\alpha\text{-Fe}$, W'1, Fe_3O_4) of abscissa $x_C = 1.070_0$, these properties are merged ($\Delta\bar{H}_{\text{O}} = -290.4 \text{ kJ.mol}^{-1}$ and $\Delta\bar{H}_{\text{Fe}} = +21.4 \text{ kJ.mol}^{-1}$).

It is noticeable that a common « kinetic origin » ω is characterized by $x = 1.070 \pm 0.001$ in W' (see Annex 2 « Re-analysis of kinetic data at 850 °C », derived from the experimental results [75]).

2.6. Triple point number 8 (W1, W2, W3) out of equilibrium

Point 8 (W1, W2, W3) was determined at the intersection of the extrapolated boundaries W1/W2 and W2/W3 at temperature 302 °C and composition $y = 0.932_2$, directly below point Ω and point C ($y_C \cong 0.935_0$) (see Ref. [30] Table III and Figure p. 3681).

Manenc *et al.* observed a disproportionation near 300 °C, in a miscibility gap, by inducing a slow drop in temperature ([76] [77]). Fender and Riley remarked that wII (\equiv W2) behaves $[\partial(\Delta\bar{G}_Q \text{ vs } 1-x)/\partial T]$ in its specific region as a two-phase system above its critical temperature. The same as Vallet and Raccach [30], they remarked that the extrapolation into the metastable domain below point C shows the phase separation into wI (\equiv W1) and wIII (\equiv W3) in the vicinity of 315 °C » [69].

As early as 1977, Andersson and Sletnes studied conjointly the ordering and the preceding behavior below 300 °C, in reference with the spinodal decomposition using electron diffraction and microscopy (dark field images in Refs. [78] [79]). Typical patterns were attributed to the phase P' (periodicity of $2.6a_0$), in the composition interval y (of Fe_yO) $\in [0.95\text{-}0.92]$. After heat treatment for 20 min at 300 °C, a tetragonal or orthorhombic symmetry was observed which can be compared to that of ordered P'-type regions. For compositions less than $y = 0.91$ ($z \geq 0.09$) in Fe_yO , the more ordered structure identified as P'' was observed even after quenching. Its parameters are $a = b = c = 5a_0$. These different observations indicate that W2 can be represented as a solid solution of W1 and W3, consistent with the determinations of $[\Delta\bar{H}_{\text{Fe}}(x)]_i^{\text{ex}}$ as formalized below.

2.7. Trends to ordering at increasing z . Excess properties

The difference $[\Delta\bar{H}_{\text{Fe}}(x)]_i^{\text{ex}}$ between the partial molar enthalpy of iron in $(\text{FeO}_x)_i$ and $\text{FeO}_{x_{o(i)}}$ on the boundary with Fe was envisaged as an excess type property in order to tentatively establish links between thermodynamic and defect structure properties [47]. The corresponding relations, using the previous formalism established for $[\bar{H}_{\text{O}}(x)]_i$, $[\bar{H}_{\text{Fe}}(x)]_i$, $[\bar{S}_{\text{O}}(x)]_i$ and $[\bar{S}_{\text{Fe}}(x)]_i$ leading to §-2.5. Figs. 6 and 7 above, become simply

$$[\Delta\bar{H}_{\text{O}}(x)]_i^{\text{ex}} = [\bar{H}_{\text{O}}(x) - \bar{H}_{\text{O}}(x_o)]_i = 19.14475(a_i/2(x - x_{o(i)})) \quad (15h)$$

$$[\Delta\bar{S}_{\text{O}}(x)]_i^{\text{ex}} = [\bar{S}_{\text{O}}(x) - \bar{S}_{\text{O}}(x_o)]_i = -19.14475(b_i/2(x - x_{o(i)})) \quad (15s)$$

$$[\Delta\bar{H}_{\text{Fe}}(x)]_i^{\text{ex}} = [\bar{H}_{\text{Fe}}(x) - \bar{H}_{\text{Fe}}(x_o)]_i = -19.14475(a_i/4(x^2 - x_{o(i)}^2)) \quad (16h)$$

$$[\Delta\bar{S}_{\text{Fe}}(x)]_i^{\text{ex}} = [\bar{S}_{\text{Fe}}(x) - \bar{S}_{\text{Fe}}(x_o)]_i = 19.14475(b_i/4(x^2 - x_{o(i)}^2)) \quad (16s)$$

Substituting the coefficients a_i , b_i and $x_{o(i)}$ on the $\gamma\text{-Fe}/\text{W}_{o(i)}$ boundary by their numerical values yields Table 1. The sign of the term $[\Delta\bar{H}_{\text{Fe}}(y)]_i^{\text{ex}}$ as well as the sign of its derivative $\partial/\partial x$ is the sign of $-a_i$. This term corresponds to the energetic effect δU on the iron lattice for the interaction between each zero, one or two electronic hole(s) being created (see *supra* relations 4a-4c). The values of $[\Delta\bar{H}_{\text{Fe}}]_i^{\text{ex}}$ in Table 1 below show that

- 1) for W1, the internal energy U_m of the iron lattice (the term $P\Delta V_m$ is negligible compared to the contribution ΔU_m in ΔH_m) decreases as z increases, which corresponds to the exothermic formation energy of clusters (m:n) that are thus favored (see Gokcen [7] p. 241-44, and [12])

p. 42-44, and Desré and Hodaj [23] p. 91-96).

- 2) conversely, W3 is characterized by an increase in the lattice energy, which corresponds to the completion of the cluster percolation and the formation of the superstructure (2.55a x 2.55a x 2.55a) named P'' by Manenc and characterized using HREM by Iijima [77] [80] [81] [82].
- 3) likely because W2 is a solid solution of W1 and W3, the variation of U_m is intermediate.

Such trends were discussed previously by Fender and Riley when considering $\Delta\bar{G}_O$ and $\Delta\bar{H}_O$ in (Ref. [69] p. 797).

Table 1 - Variations at 1 273 K of excess properties $[\Delta\bar{H}_O]_i^{ex}$, $[\Delta\bar{H}_{Fe}]_i^{ex}$ and $[\Delta\bar{S}_O]_i^{ex}$, $[\Delta\bar{S}_{Fe}]_i^{ex}$

FeO _x /Fe _{1-z} O	$x=1.095\ 0$ or $z=0.086\ 8$	$x=1.142\ 0$ or $z=0.124\ 3$
W1: $a_1= 46\ 753.4$ $b_1= -7.378\ 1$ $x_{o(1)}= 1.057\ 2^*$ $z_{o(1)}= 0.054\ 1$	$[\Delta\bar{H}_O]_1^{ex}$ 16.90 $[\Delta\bar{S}_O]_1^{ex}$ 2.67 $[\Delta\bar{H}_{Fe}]_1^{ex} **$ -18.19 $[\Delta\bar{S}_{Fe}]_1^{ex} **$ -2.87	$[\Delta\bar{H}_O]_1^{ex}$ 37.94 $[\Delta\bar{S}_O]_1^{ex}$ 5.99 $[\Delta\bar{H}_{Fe}]_1^{ex}$ -41.72 $[\Delta\bar{S}_{Fe}]_1^{ex}$ -6.58
W2: $a_2= -9\ 568.9$ $b_2= 31.172\ 8$ $x_{o(2)}= 1.052\ 2$ $z_{o(2)}= 0.049\ 6$	$[\Delta\bar{H}_O]_2^{ex}$ -3.92 $[\Delta\bar{S}_O]_2^{ex}$ -12.78 $[\Delta\bar{H}_{Fe}]_2^{ex}$ 4.21 $[\Delta\bar{S}_{Fe}]_2^{ex}$ 13.76	$[\Delta\bar{H}_O]_2^{ex}$ -8.23 $[\Delta\bar{S}_O]_2^{ex}$ -26.80 $[\Delta\bar{H}_{Fe}]_2^{ex}$ 9.03 $[\Delta\bar{S}_{Fe}]_2^{ex}$ 29.41
W3: $a_3= -33\ 238.9$ $b_3= 48.366\ 9$ $x_{o(3)}= 1.048\ 2$ $z_{o(3)}= 0.046\ 0$	$[\Delta\bar{H}_O]_3^{ex}$ -14.90 $[\Delta\bar{S}_O]_3^{ex}$ -21.68 $[\Delta\bar{H}_{Fe}]_3^{ex}$ 15.97 $[\Delta\bar{S}_{Fe}]_3^{ex}$ 23.24	$[\Delta\bar{H}_O]_3^{ex}$ -29.85 $[\Delta\bar{S}_O]_3^{ex}$ -43.44 $[\Delta\bar{H}_{Fe}]_3^{ex}$ 32.69 $[\Delta\bar{S}_{Fe}]_3^{ex}$ 47.57

* For $x= x_{o(i)}$, all terms are equal to zero. ** The terms $[\Delta\bar{H}_O]_i^{ex}$ and $[\Delta\bar{H}_{Fe}]_i^{ex}$ are expressed in kJ.mol⁻¹, and the terms $[\Delta\bar{S}_O]_i^{ex}$ and $[\Delta\bar{S}_{Fe}]_i^{ex}$ in J.K⁻¹.mol⁻¹.

The positive and negative values of the term $[\Delta\bar{S}_O]_i^{ex}$ (see Table 1) correspond mainly to a vibrational entropy change in the oxygen sublattice [20] [21]. In the case of the initially disordered pseudo-phase W1, the excess entropy term weakly increasing with z shows that the disorder due to the formation of point defects is predominant. Reversely, in the case of the pseudo-phase W3, increasing short- and long-range orderings would be associated with increasing z values as indicated by a strong decrease of $[\Delta\bar{S}_{Fe}(x)]_3^{ex}$. These entropy variations are related to the increasing frequency of the Fe-O bond because of larger clusters and/or an increasing degree in their percolation until completion and superstructure. The pseudo phase W2 can be thought at the starting threshold of the percolation of clusters, when the lattice becomes more rigid.

3. SEVERAL ACCESS OF INVESTIGATION

3.1. Burgmann's review

As soon as 1968, after he has defended a thesis in 1967 in Berlin, this author reviews the most important experimental results (electrical conductivity and transition p to n) at this time and their interpretations when any [6]. Because of numerous analogies with pyrrhotine Fe_{1-z}S [83], he describes tentatively the defect structure in terms of clusters (see ref. [6] Fig. 5 p. 173). He proposes a conduction model involving the formation of an acceptor band or new states produced by the interaction of defects as z increases.

The lattice contraction at increasing z reduces the gap between the valence and conduction bands, inducing a p-n transition for compositions $z = 1/12 = 0.083$ at 900 °C and $z = 1/13 = 0.076$ at 1 300 °C [6], also Hillegas's Thesis [84]. Such an overlap is also considered later by Molenda, Stokłosa and Znamirowski [85].

3.2. Manenc's phases P, P', P''

As the most important result of a seminal work, **Manenc** describes three types of « phases » named P, P' and P''. He characterizes them after (rapid) quenching into water using X-ray diffraction. For an iron rich wüstite ($z < 0.08$), vacancies and ions Fe^{3+} seem to be randomly distributed. This corresponds to the phase P which has a weakly incomplete NaCl-type structure alone. Other diffraction patterns of more iron deficient samples ($0.08 < z < 0.11$) have additional reflections corresponding to a cubic cell with a repeating distance $2.6a_0$. Diffuse trails parallel to $\langle 100 \rangle$ directions connect these reflections. These features are characteristic of the phase P' alloyed with the phase P, and individually observed only in « oxygen rich » samples. Iron vacancies are ordered in the two phases [77], [80].

Diffraction for a heated single crystal such as $z = 0.08$ shows that a superstructure exists in the range 800 -1000 °C for this composition. The phase P'' is observed in wüstite such as $z \cong 0.10$ near 300 °C and when slowly cooled, at the threshold of disproportionation yielding Fe and Fe_3O_4 . New superstructure lines are then characteristic of a commensurate 5X cubic cell [77], [80].

3.3. Mössbauer spectroscopy

Greenwood and **Howe** [86] interpret Mössbauer spectra in modeling the asymmetrical doublet as evidencing the different defect structures of “*the pseudo-phases identified by Vallet and Raccach [30] [46], and Fender and Riley [69] at high temperature.*” The quenching process produces nuclei leading to Manenc phases P, P', [(P+P')] for $z \in [0.05-0.08]$, (P') for $z \in [0.08-0.10]$, and (P'') for $z > 0.10$ below point C. The phases are detectable from fine variations of quadrupole splitting resulting from cubic symmetry distortions. These distortions can be evaluated for the single cluster (4:1), and the Koch and Cohen cluster (13:4). From a sample such as $z \cong 0.05$, the disproportionation gives rise to $\text{Fe}_{0.918}\text{O}$ (defect rich phase P') and $\text{Fe}_{0.976}\text{O}$ (defect poor phase P), because of a mixture of the (13:4) and (4:1) clusters.

Among numerous other Mössbauer studies reviewed by Long and Grandjean [17], **Romanov**, **Checherskaya** and **Tatsienko** [87] [88], **Hryniewicz et al.** [89], **Pattek-Janczyk et al.** [90], this latter is of particular interest. The authors modeled the asymmetric doublet with a singlet attributed to Fe^{3+} in tetrahedral sites, and two doublets attributed to octahedral Fe^{2+} and Fe^{3+} . They associated the maximum observed in the variation of the quadrupole splitting to the p-n transition described by [89]. This transition is likely the same as the transition « metal \leftrightarrow Mott insulator » identified by **Molenda**, **Stokłosa** and **Znamirowski** [85].

Aldon and **Jumas** use the Mössbauer spectroscopy to evaluate precisely the reversible insertion of lithium and the effects on the local and long-range orderings in the wüstite lattice. They characterize the lithium as an « actor of electronic pressure » (see section 4.3. *Nanocomposites. Epsilon-Fe*).

3.4. First electronic-structural modeling according to J. B. Goodenough

Goodenough's modeling [3] is proposed as early as 1971. Considering relation (2) p. 4 associated with thermogravimetric analyses, and relation $\sigma \propto p'^{1/5}$ (17) corresponding to the conductivity [67], the result $s=6$ in the W1 subdomain above 1 060 °C characterizes a statistical distribution of vacancies and electronic holes (Fe^{3+}), possibly corresponding to Manenc's phase P. Below 1 060 °C, the changes in values of s correspond to an increasing association between vacancies and holes. In the W2 domain, triplets $\text{Fe}^{3+} - \text{V}_{\text{Fe}} - \text{Fe}^{3+}$ are partially ordered, improving Madelung energy. Numerous (4:1) clusters form

a superstructure according to Koch & Cohen (see section 7.2. p. 24). The *historical* cluster (13:4) should correspond to W3.

A mapping of the transfers of iron ions, and electronic exchanges is proposed as a mechanism needing small polarons related to energy differences between the top of the valence band and an acceptor level [3]. Later **Gleitner** and Goodenough will consider the wüstite again from the point of view of its electronic properties and crystallographic structure. They will express some doubt concerning the role of the cluster (4:1) as only building unit of larger clusters ([13]).

3.5. Zvintchuk's team

In 1973, **Tchiong Tki Khong** by **Zvintchuk's** team, in collaboration with J. Manenc, studies samples of wüstite of composition $z \in [0.065-0.108]$ quenched from 1 100 °C [91]. Electron diffraction reveals three forms I, II and III corresponding to the modifications wI, wII and wIII of the phase diagram by Fender and Riley [69], and to the phases P, P', P''. A correspondence with the three Wi is specified.

The authors describe a defect complex (16:6) of point symmetry $m\bar{3}m$, formed of six tetrahedral corner sharing clusters (4:1) distributed on the six faces of an empty cube. This cluster is represented in Fig. 8 below. A later work by Zvintchuk's team describes a superstructure with a quadratic cell $P\bar{4}3m$ for a wüstite $Fe_{1-0.10}O$ [92].

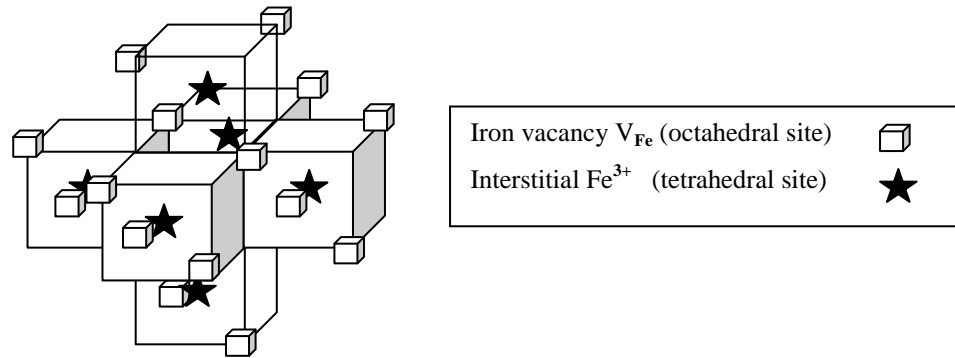


Fig. 8. The $m\bar{3}m$ cluster (16:6) of $\langle 110 \rangle$ type corresponding to $(z+t) / t = 2.66$ (from Ref. [91]).

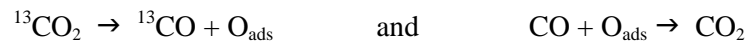
3.6. Worrall and Coley - 2010, 2013. Isotopic kinetic transitions

Except Wagner, jr *et al.* [67] in 1966, and Fender-Riley [69] in 1969 who never more worked positively on the topic, it was necessary to wait the Goodenough-Gleitner works from 1972, Men-Lykasov's works at Sverdlovsk-SSSR (around 1975), and Mrowec-Stoklosa's work at Krakow-Poland (after 1980) to confirm experimentally and theoretically Raccach-Vallet's discovery (1962) of three wüstite modifications [30] [44] .

Following the process first used by Temkin *et al.* at the surface of iron oxides ([93] [94]), and then applied by Grabke [95], Cramb and Belton [96], Mori, Morita and Sano [97], later by Xiaojun Hu *et al.* [98], Teng Zhang *et alii* [99], **Worrall and Coley** study the kinetics of the carbon isotope exchange in the equilibrium W/ (CO_2/CO mixtures) in 2010 [75]. A controlled proportion of $^{13}CO_2$ species is introduced in the environment made by the gas mixture in the furnace so that a specific equilibrium of the carbon isotopes is established following



which can be viewed as the sum of the two half-reactions



as reflecting oxygen exchange on the surface of the sample, and being conditioned by the electronic structure in the bulk. The apparent rate constant is expressed as a function of oxygen activity a_O as follows:

$$k_a = k_0 \cdot a_{\underline{O}}^{-m} \quad \text{or} \quad \log k_a = -m \cdot \log a_{\underline{O}} + \log k_0 \quad (18)$$

The variation of $\log k_a$ with the activity $a_{\underline{O}}$ defined as being the ratio $p(\text{CO}_2)/p(\text{CO})$ can be described by three successive linear segments at 995 °C, with $m = 0.51$ (w1 ?), 0.66 (w2 ?), 1.03 (w3 ?) (ref. [75] Fig. 8 and equations (36) (37) (38) p. 820-1).

The authors observe only one transition attributed to two pseudo phases (w'1- w'2) and w'3 along the isotherm at 850 °C, no change being observed in the vicinity of the boundary W'1/W'2 forecasted by Vallet and Carel [31]. The only observed transition appears as a discontinuity at the location of the transition W'2 \leftrightarrow W'3 forecasted for $x = 1.090$ at equilibrium. Nevertheless, the authors state that three pseudo phases w'1, w'2 and w'3 probably exist, but that two of them could not be separated, either because they were not detectable by electrical conductivity measurements as found before in [67], or because they were possibly « *obscured by the scatter* ». Finally, Worral and Coley declare that, in their study, transitions at 850 °C and 995 °C were observed « *correspond[ing] exactly with the proposed pseudo phase boundaries in quite good agreement with the phase diagram by Vallet and Carel* ». Because of the identity of coefficients m in the correlations concerning W_i and W'_j for $i=j$, they attributed in a following review [24] three predominant clusters (7:2) type II to w1 and w'1, (12:4) type II to w2 and w'2 and (16:5) type I to w3 and w'3.

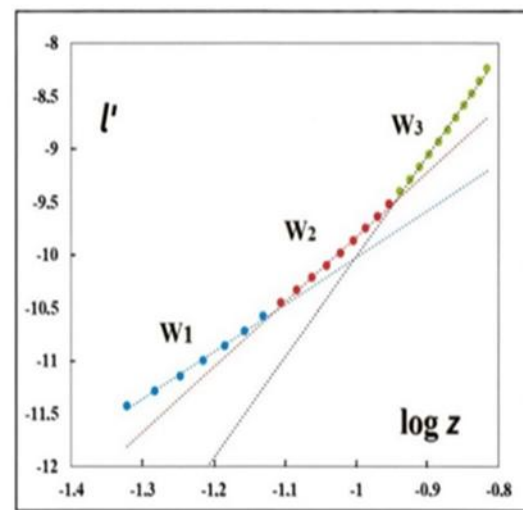
Note: In ANNEX 4 – Fig. A4. p. 40 a graphical and numerical re-analysis of Fig. 9 by (Ref. [75] p. 821) shows a possible separation of the data in three sets at 850 °C, which correspond likely to three w'_j , provided that some data or groups of data are located in subdomains not proposed as stable. In ANNEX 1 - Fig. A1 the location of the observed transitions is represented by the set f-) symbols $\Rightarrow \Delta$ corresponding with transitions W'1/ W'2 and W'2/ W'3.

3.7. Takayama and Kimizuka's phase modifications

Using thermogravimetric analyses under equilibrium oxygen pressure p' defined by CO_2/H_2 mixtures, Takayama and Kimizuka observe significant modifications in curves representing l' vs $\log x$ in FeO_x [72]. The authors conclude to the existence of a transition between pseudo-phases (see Fig. 9 and Réf. [72] Table I, and ANNEX 1 – Fig. A1 set d-) symbol $\Rightarrow +$).

Typical changes in slopes are clearly observed below in Fig. 9 from Takayama and Kimizuka's data.

Fig. 9. Variation of l' vs $\log z$ for Fe_{1-z}O at 1 250 °C. Three sets of data can be defined as three quasi linear relations



To illustrate them, we have transformed their representations for one temperature $T = 1\,523$ K. For this temperature, their data are interpreted in terms of two relations only.

– For low x values ($x = 1/(1-z)$): $l' = M_1 \cdot x + B_1 = 28.512 x - 41.366$ (19)

– For high x values: $l' = M_2 \cdot x + B_2 = 23.288 x - 35.715$ (20)

The new representation for l' vs $\log z$ are done in Fig. 9. Two coefficients S can be defined as a first step, using the hypothesis of a transition implying two pseudo-phases proposed by the authors: S close to 4 and 7.5. However, three correlations can be also defined as follows:

$$l' = 4.4539 \log z - 5.5685, \quad l' = 6.1378 \log z - 3.6876, \quad l' = 9.5059 \log z - 0.5046 \quad (21)$$

from which three values for S can be derived. A first domain for low z values yields a coefficient $S = 4.5$. So, it is possible to interpret the non-linear second domain in terms of two linear sequences with two coefficients $S = 6.1$ and 9.5 : these two last coefficients are quite close to those obtained from the data of Vallet *et alii* (at lower temperature).

The results of Takayama and Kimizuka [72] and those of Bransky and Hed [100] show a relative agreement in the same domain of temperature (see also ANNEX 1 – Fig. A1 sets d- and b- \Rightarrow + and ●) but disagree with results of Fender, and those of Vallet « *because these authors did not observe any breaking point on the isothermal lines above 911 °C* » (see *infra* section 10. Percolation approach, Nota: Metastabilities).

4. ELECTRICAL PROPERTIES. TRANSITIONS

In an extensive review about electrical properties, Gleitzer points out that the modeling of the electrical conductivity of wüstite is relevant of the percolation concept [21]. The superimposition of the two kinds of p-n and W_i / W_{i+1} transitions, almost independent of each other, and both function of z , is already referenced in Burgmann's review [6] (See Fig. 5 p. 173).

Because of this common stress, it is not easy to correlate the p-n transition to the structure. Since the percolation process is characterized experimentally by a steep variation of the studied property at the threshold composition, the curves for the electrical conductivity vs composition particularly in quenched wüstite can be interpreted in terms of upgradable clusters and long-range ordering in the successive subdomains of W_1 , W_2 and W_3 [67] [100] [101] [102].

Later in the present review (see Section 10 Percolation approach p. 35-6 and Fig. 20), a statistical distribution of clusters with two types of percolation modes is schematically represented through the phase diagram. This cluster distribution can justify structurally the singularities of the integral and partial thermodynamic properties (see above Figs. 1, 2, 5, 6, 7). Also, it can justify the distinct rate constants observed in kinetic processes (see section 3.6. Kinetics transitions, and ANNEX 3 – Figs. A3 and 4 – Figs. A4).

4.1. The p-n transition

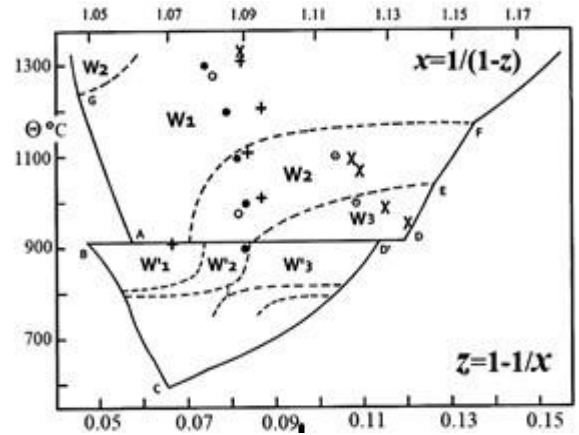
In 1962 **Tannhauser** observe a thermo-electronic p-n transition around 1 300 °C by thermal e.m.f. analyses of a highly nonstoichiometric wüstite [103]. More precisely with Bransky [104], a change in sign of the Seebeck coefficient is observed for polycrystalline samples near $O/Fe \cong 1.09$ or $z \cong 0.083$ at $T \in [1\ 010 - 1\ 310]$ °C and near $z \cong 0.066$ at 910 °C (Fig.10). The authors formalize the Seebeck coefficient α , in the case of completely ionized vacancies, when the hopping process is predominant

$$\alpha = (k/e) \cdot (A/kT + \ln(c_o - c)/c) \quad (22)$$

c = 'concentration' of charge carriers, c_o = 'concentration' of available sites for the charge carriers (see also [105]).

The same formalism is used by **Lafollet** and **Duquesnoy** in a *non-localized* case when the wüstite is considered as an intrinsic semiconductor, *i.e.*, doped by its own impurities [106]. They modeled $\alpha = 0$ for $z = 0.104$ at 1 100 °C, and $z = 0.108$ at 1 000 °C (see ref. [106] in Fig. 10).

Fig. 10 - Some thermal determinations of the p-n transition by emf measurements; X: Geiger *et al.* [67] not corrected for platinum; + : Bransky and Tannhauser [104]; ⊖ : Lafollet and Duquesnoy [106] by modeling; ○ : Hillegas jr and Wagner jr [107]; ● : Gartstein and Mason [108] by re-examination of Hillegas jr's thesis [84]



Like in previous studies, **Hodge** and **Bowen** state that the experimental sign change of the thermal electromotive force from positive to negative occurs at increasing values of z [109]. Their model considers principally clusters (4:1) and electron holes trapped in octahedral sites adjacent to cluster vacancies. The thermally activated hopping of these holes permits continuous paths through the crystal because of the high point defect concentration between the zones separating the clusters.

In their studies on electrical conductivity, **Ariya** and **Bratch** (1963), and **Kozheurov** and **Mikhailov** (1967) observe modifications to the relationship $\sigma \propto (p')^{1/n}$ closely related to the relationship $z \propto (p')^{1/n}$, where their coefficient n ($\equiv 5$) varied in such a way that three subdomains can be defined [101] [102]. They find an increase of the conductivity in two composition ranges, from $z \approx 0.048$ to 0.079 then from 0.079 to 0.097, in relation with the Fe/W and W/Fe₃O₄ boundaries.

4.2. Electrical conductivity

Neuschütz and **Towidi** establish a variation of the conductivity for $T \in [700 - 1220 \text{ °C}]$ from which they inferred that the defect electrons jump between Fe²⁺ and Fe³⁺, the activation energy of the process being 9.6 kJ.mol⁻¹ [110].

Following **Molenda**, **Stoklosa** and **Znamirowski's** paper the conductivity can be the sum of an « *extrinsic conductivity* σ_y » related to the nonstoichiometry, and of an « *intrinsic component* σ_o » independent of the defect concentration *i.e.* of p' [85]. Charge transport at high temperature would occur simultaneously in the valence and conduction bands, and the dopant acceptor band due to Fe³⁺ in octahedral and tetrahedral sites. Near the composition $z \approx 0.08$, the acceptor band would overlap the valence band giving rise to a broad metallic band close to the thermoelectric transition p-n for $z \approx 0.09$ and corresponding to the stability limit of W1 and W2-3. Experimental details are given in ref. [85] (Fig. 2 p. 519 and Fig. 4 p. 522). See also Annex A, Fig. A1 e- \rightarrow x, where successive segments and breaks corresponding to the Raccach-Vallet's isotherms are observed.

More recently **Toroker** and **Carter** examined the possible improving of the conductivity of Fe_{1-z}O by doping the p or n conduction mode, with the aim of stabilizing the nanoscale fabrication of wüstite, and then adapting the nanoparticles to create efficient solar light conversion materials [111]. They listed the procedures allowing the calculation of Marcus theory parameters for cluster models with hydrogen dopant and substitutional Cu⁺, Li⁺ and Na⁺. They concluded that « iron vacancies drastically limit the hole conductivity » while Li, H and Na dopants improve it. So, it would be interesting to suppress the formation of vacancies by “alloying FeO”, which would amplify the p-type conductivity.

Note: An extensive thermodynamic study of numerous substituted wüstites can be found in the work of **Lykasov et alii**, which is reviewed partly in [15] and [112]. Structural aspects of calcio- and magnesio-wüstites are available *infra* (see in Fig. 16 p. 27, and Table II p. 28)).

4.3. Nanocomposites. Epsilon-Fe

Aldon and **Jumas** examine the electrochemical redox phenomena at the negative electrode $\text{Fe}/\text{Fe}_{1-z_0}\text{O}$ formed by wüstite of limit-composition $z_0 (= 0.050 \pm 0.013)$ in Li-ion batteries [113]. The initial samples of hematite powder, reduced in H_2 at 800°C , are quenched outside the furnace. The analyses by Mössbauer spectroscopy and X-ray diffraction after the first discharge yield the nature and concentration of the products of reduction of wüstite by Li. Nanoparticles of iron ($\sim 2\text{nm}$) embedded in a Li_2O matrix are identified. A part of the iron atoms appearing in the disproportionation is made of the $\varepsilon\text{-Fe}$ phase which is stressed in the matrix under a pressure evaluated at 6.7 GPa , and which is not ferromagnetic. It can be due to the extrusion from the W structure. At the same time, small clusters of lacunary hematite are formed. These phenomena have been still identified in the formation of « planetary materials, in the Earth's inner core ». See also geological observations by Hazen and Jeanloz [12].

5. OTHER KINETIC AND DIFFUSIONAL WORKS

5.1. Iron self-diffusion

In 1963, **Desmarescaux** and **Lacombe** study the self-diffusion in using the iron isotopes ^{55}Fe and ^{59}Fe [114]. They state experimentally (see Fig. 11) that the self-diffusion coefficient D^*_{Fe} increases with z at a given temperature T . At increasing T , because the clusters are dissociated, the rate of free vacancies is increased, the self-diffusion coefficient too. They deduce that self-diffusion proceeded through a vacancy driven mechanism, especially noticeable at high temperature. At lower temperatures, fewer free vacancies and more complex vacancy-electron holes - *later naming clusters (amas in French)*- exist. Diffusion is then slower. The number of vacancies associated to Fe^{3+} cations [the clusters] increases with z , so a compensation with the free vacancies implies a quasi-constant D^*_{Fe} whatever z below 850°C .

These observations are interpreted in the Collongues's review [20] with the help of theoretical evaluations of the binding energies in clusters (see Sections 6.5. *Cluster stability*, and 6.6. *Diffusion and stability*).

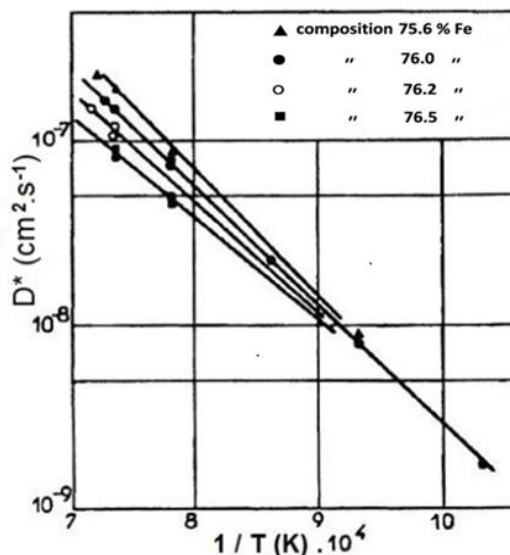



Fig. 11 – Fe diffusion (from ref. [114], Fig.4)
 $\log D^*_{\text{Fe}}$ vs $1/T$. Fe % = 75.6 $\square z = 0.112$;
 76.0 $\square 0.093$; 76.2 $\square 0.083$; 76.5 $\square 0.067$

5.2. Isothermal kinetics

Landler and **Komarek** measure the loss of mass vs time $\Delta M(t)$ in $\text{H}_2/\text{H}_2\text{O}$ mixtures by thermogravimetry when reducing polycrystalline samples [115]. In their calculations, the coefficients of iron self- and chemical diffusions decrease for increasing z at constant $T \in [800 - 1050]^\circ\text{C}$ because of

« the increased freezing due to repulsive interaction of the vacancies ». That corresponds to activation energy independent of T. This observation is related primarily to the formation of clusters, then to the ordering of clusters in the bulk (see Section 2.7. *Trends to ordering*).

Considering the relationship $\log(\Delta M/\Delta M_0) = n \cdot \log t + k$, in which ΔM_0 is the total mass loss, reduction kinetics data from the authors at 950, 1 000 and 1 050 °C can be separated into three regimes II, III, and IV (see Annex B, Fig. B1). At 1 000 °C, the calculated mean values of n are -0.72 (W3?), -0.49 (W2?) and -0.3 (W1?) respectively [116]. A brief anharmonic sequence is observed at the transitions between two reduction processes. The increase of 0.6 eV in the activation energies $E_{III} = 1.31$ eV and $E_{II} = 1.89$ eV for the reduction sequences III then II respectively is observed at decreasing z, i.e., when the vacancy concentration decreases. It corresponds to the decrease of the ordering process during the transformation W3→W2. This variation is opposite to other results in the literature, particularly to (ref. [114] Fig 5, see section 5.1. *Self-diffusion*). The three coefficients n are primarily the signature of three modifications in the bulk by Landler and Komarek, while the three kinetic coefficients m by Worral and Coley are primarily characteristic of surface phenomena [115] [75] (see [117]: Bulk and surface process in the course of the reduction).

Similarly, successive different reduction sequences using H₂O/H₂ gas mixtures that exhibit transitional anomalies can be observed by **Rieke** and **Bohnenkamp** [118] (see ANNEX Fig. A1. Set c) ⇒ ).

5.3. In situ diffusion

Rickert and **Weppner** combine a solid galvanic cell for measuring the ionic conduction by O²⁻ through the sample when the other face of the sample is in contact with a platinum foil [119]. The diffusion is measured *in situ* during the relaxation from a composition to another, driven by controlled impulses of electrical potential. Their main result is that the chemical diffusion coefficient increases with increasing z (i.e. with increasing vacancy concentration), which disagrees with Landler and Komarek [115] but not with any other authors. Transitional anomalies are clearly noticed (see [119] Fig. 6 and 7 p. 1856, and Fig. 8 p. 1857).

6. RELATIONSHIPS BETWEEN DEFECT STRUCTURES AND PHASE DIAGRAM

Key publications are reviewed from authors who have attempted to solve the challenge formulated by Carl Wagner in 1930-1940 [2] of linking more tightly defect structure and phenomenological thermodynamics, such as partial molar properties or the equilibrium constant of formation of the clusters. Kinetics (diffusional) relations were also developed (see Ref. [20]) but do not seem to be accurate enough. Electrical conductivity is particularly suited to the problem (see the big and constructive Review by Stokłosa in ref. [25]).

6.1. Per Kofstad and Zeev Hed

These authors remark that z increases at constant p' as T decreases, which is different than other oxides, and likely due to a complex defect structure [4] [120] (see section 5.1. *Iron self-diffusion* by authors [114]). They emphasize that there is some doubt that all the iron vacancies are doubly charged. The electron holes associated to a vacancy can jump into only one of the 12 next nearest octahedral neighbors ($\text{Fe}_{\text{Fe}}^{\circ}$) and one tetrahedral site ($\text{Fe}_i^{\circ\circ}$). So most of octahedral vacancies would be singly charged.

A model can thus be elaborated from the Roth complex $[\text{V}_{\text{Fe}}^{m(r)} - \text{Fe}_i^{n(\circ)} - \text{V}_{\text{Fe}}^{p(r)}]^{q(r)}$, n=1,3, (m,p) = 0,1,2. In the relationship $\sigma \propto p'^{1/s}$, the s value is larger than the value in the similar relationship connecting z to p' because vacancies would be neutral.

6.2. Toft Sørensen's works

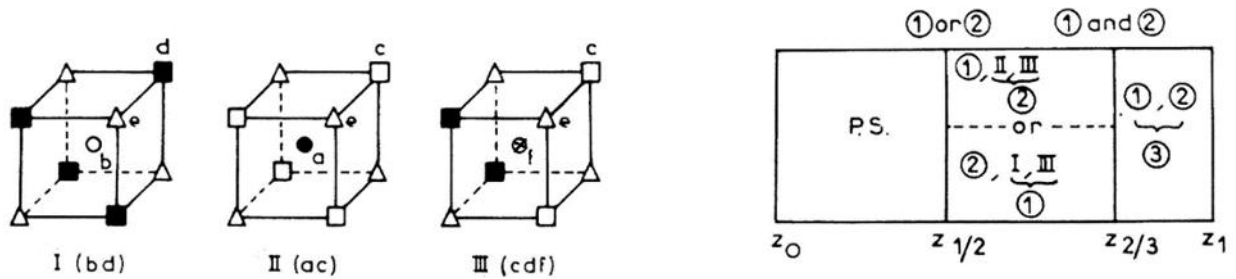
In a book devoted to nonstoichiometric oxides, Toft Sørensen then with El Sayed Ali [11] [48] tentatively identify the defect structure of the « subphases » in using thermogravimetric data [l' -z] at

1 000 – 1 300 °C exclusively by Bransky and Hed [100]. They consider three main defects (doubly charged vacancy V_{Fe}'' , tetrahedral complex defects (4:1)^{(m)-} and (6:2)) as defined by Catlow *et al.* (see ref. [11] chapter 2). They focus on the defect clusters differently ionized, which leads them to define a linear plot of l' vs $\log y (\equiv z)$ with minimal slopes 1/5 in the relation $z \propto (p')^{1/5}$. A second criterion concerns the so-called exclusion envelope, *i.e.* the long-range order as a layer or a close packing structure. They localize in the wüstite equilibrium area T(z) the predominant defect (4:1) plus a layer structure, (4:1) plus a close packing, (6:2) plus a layer structure, according to ΔT layers above 1 300 °C, between 1 300 and 1 200 °C and below 1 200°C, respectively. In addition, iron vacancies V_{Fe}'' are present on the right side of a line at composition z close to 0.09 (see [48], Fig. 7 p. 20). So, a structural characterization is roughly proposed, and the phase diagram of Vallet and Raccach [30] and Fender and Riley [69] is “approximately drawn again”.

6.3. Modeling by means of the C(luster) C(omponent) M(ethod)

In the course of their theoretical work about oxides properties, **Men' *et alii*** were concerned with the wüstite phase [8] [121]. Men' and Carel apply C.C.M. to modelling the wüstite solid solution.

The « cluster components » are defined from unit blocks that is crystallographic positions forming in the general spinel structure. For example, three cluster components ①, ② and ③, and their concentration



Structural **unit blocks S.S. I, II, III**: positions ■: (16d), □: (16c), ○: (8b), ●: (8a), ⊗: (48f), △: (32e); a and d are for iron, c and f for vacancies (see Ref. (6), Table 7.1., p. 186). Short range order **in terms of S.S. and cluster components C.C. ①, ②, ③** in the three sub-domains; C.C. ①: I-6III, C.C. ②: II-6III, C.C. ③: I-6III-II. In the range $[z_0 - z_{1/2}]$, II and III increase, I decreases; in $[z_{1/2} - z_{2/3}]$, II and III decrease, ① or ② increases; in $[z_{2/3} - z_1]$, ① and ② decrease, ③ increases.

Fig. 12: C.C.M. applied to wüstite (from Ref. [122]). The three pseudo phases from z_0 on the γ -Fe/W1 boundary to z_1 on the W3/Fe3O4 boundary described with the help of three cluster components ①, ②, ③.

as a function of z allow to build the lattice of the wüstite $Fe_{1-z}O$, and to model some of its physical properties, taking into account the pseudo-phases as allotropic varieties: Fig. 12 (see Refs. [122] [123] [124]).

6.4. Other modeling approaches.

Mrowec *et alii* work on defect clustering in the wüstite phase using a collection of data (z , $\log p'$) from the literature [16] [49]. They envisage all the defect clusters suggested by Catlow and Fender [125], *i.e.*, (4:1), (6:2), (8:3), (10:4), (12:4) and (16:5) with several ionization degrees including the highest one. Elaborate calculations of the equilibrium constant lead them to the assessment of the enthalpy and entropy of formation of all these ionized clusters, then to $\Delta \bar{H}_O$ independently of z and T, and to $\Delta \bar{S}_O$ for $z \in [0.055 - 0.145]$. They conclude that the (4:1) cluster is the most likely to occur and that the ionization degree increases with temperature to reach the maximum value 5- at 1 573 K.

In a meeting in Poland (1980), Gavarrri and Carel envisage for the first time a modeling of the distribution of two clusters in the three subdomains $z_0 - z_{1/2}$, $z_{1/2} - z_{2/3}$, $z_{2/3} - z_1$, from Fe to Fe3O4 [126] (see *Addendum 1*, p. 40).

Nowotny and **Rekas** investigate the clustering of point defects (iron vacancies and electron holes) across the entire domain using the Debye-Hückel theory, which allows to use a product $\gamma.[x_i]$ in place of $[x_i]$. With this approach, the activity coefficients are merged with the sole mean coefficient f_{\pm} determined as a function of the *ionic force* of the solution (not of its *square root*?). They modeled the interactions during the formation of the cluster (4:1) taking into account activity instead of concentration. Sets of data ($z, \log p'$) and ($\Delta\bar{H}_O, z$) were gathered from the literature. They evaluated the enthalpy and entropy of formation of a cluster (4:1)⁵⁻ at -396.7 kJ.mol⁻¹ and +267.4 J.K⁻¹.mol⁻¹, respectively. They concluded that the clusters (4:1)ⁿ⁻ would be the constitutive module of more extended clusters. An expression of the equilibrium constant for the cluster formation is determined across the whole domain independently of the existence of pseudo phase transitions [127].

6.5. Cluster stability

Catlow and **Fender** tried first to access the defect structure by means of large-scale computation of the binding energies following the Born model [125]. They envisage in a theoretical manner the clusters unity (4:1), edge-sharing (6:2) and (8:3), corner sharing (16:5) precursor of Fe₃O₄. The location (zone $[z, T]$ of the phase diagram) of these clusters is suggested qualitatively, but not in relation with the pseudo phases (yet well known since the paper by Fender and Riley [69]).

Later Catlow *et al.* assessed the formation energy and the relative stability of the clusters by means of the Mott-Littleton method. The cluster (12:4) <110> type by Lebreton and Hobbs is considered as being the most stable extended cluster, after clusters (6:2) and (8:3) [128] [129]. In a summary-paper Catlow envisaged four classes of clusters: the corner-sharing clusters (4:1), (7:2), and (16:5); three edge-sharing clusters (6:2), (8:3) and (12:4); the ZnS blende like cluster [129] [130] [131]; the cluster envisaged by Grimes *et al.* is not the most likely because it was characterized without taking into account the geometric relaxation, that they notice in reference [132]; lastly the clusters (13:4) and (5:2) [133] [134].

Tomlinson and **Catlow** attempt to relate the structural modeling of short-range order, and the macroscopic variation of z as a function of p' . The calculated binding energy per vacancy leads them to order the clusters as follows: [(8:3), (12:4), (6:2), (16:5)⁻¹, (12:5)⁻¹, (13:4)⁻², (4:1)⁻¹, (10:4)⁻⁴, (5:2 M), (5:2 G)] from 2.24 to 0.71 eV, successively. Finally, they propose a mixture of clusters [(12:4), (6:2), (4:1)] for the incommensurate sub-phase P' [136].

Then, **Grimes** *et al.* consider the structure and stability of clusters with different coordination types [135] [137]. They base their analysis on a generalized crystal field for the 3d electrons of Fe³⁺ and Fe²⁺, including the energy of the orbitals 2s and 2p of oxygen. Their new computation of the binding energy per vacancy favors the clusters <110> ZnS blende and <111> spinel type stacking. It rules out clusters (8:3) and (10:4), based on <100> stacking of (4:1) units, and the Koch and Cohen (13:4) cluster. Within the lower field of stoichiometry (P' phase), smaller clusters coexist including (4:1), (6:2), (7:2) <110>, (10:4) ZnS blende type, and (7:2) <111> spinel type. Larger clusters from the (12:4) <110> type, (18:6), (16:5) spinel <111> type, (16:5) <110> type and (16:7) <110> ZnS blende type clusters can be invoked at higher nonstoichiometry (P'' phase). Finally, they assert that the charge of the clusters should be zero in order to agree with experimental data.

6.6. Diffusion and stability

In a review about nonstoichiometric oxides, **Collongues** [20] envisages particularly the clustering as specified by Desmarescaux and Lacombe (see *section 5.1. Self-diffusion of iron*). The activation energy of the diffusion increases with increasing z , while the opposite effect is the principal rule. « Vacancy clusters » are the reason of this unusual behavior. The existence of three domains in the phase diagram is envisaged [20].

Catlow *et al.*'s calculations of (Fe^{3+} - vacancy) binding energy indicate that clusters such as (6:2) or (8:3) of size lower than that of (16:5) (precursor of magnetite) are more stable than this latter. This feature would indicate that even a cluster structurally far from magnetite is responsible of the reduced diffusion process.

NOTA: The classification of the clusters become more complicated because of the increasing number of proposals in the literature. Initially, only the criterion of corner or edge sharing of the basic (4:1) units was needed. Mixed cases were thought to be the exception. Lebreton and Hobbs [129] introduced three ways of clustering: *edge sharing along $\langle 100 \rangle$ (Type I)*, *corner sharing along $\langle 110 \rangle$ (Type II)*, and *corner sharing along $\langle 111 \rangle$ (Type III)*. Mixed ways are also to be envisaged. Clusters based on composition alone are possible exceptions (*e.g.*, cluster (10:3)). Labidi and Monty (see [138] p. 100-1) resumed the situation differently with the *types 1 (face shared) aligned in $\langle 100 \rangle$ direction*, *2 (edge shared) aligned in $\langle 110 \rangle$ direction*, and *3 (corner shared) aligned in $\langle 111 \rangle$ direction*. Worral and Coley (ref. [24] p. 24-6) classified the clusters in three types defined as either *type 1 (corner shared as in magnetite)*, *type 2 (edge shared sharing one octahedral vacancy)*, or *type 3 (edge shared sharing two octahedral vacancies)*.

7. STRUCTURAL APPROACHES. MODELS FOR SHORT- AND LONG-RANGE ORDERINGS.

7.1. Roth model: the cluster (2:1)

Considering the experimental uncertainties inherent to X-ray diffraction on polycrystalline oxides, it is more suitable to use neutron diffraction experiments to determine correctly the site occupancies relative to iron atoms in octahedral and tetrahedral sites of the NaCl-like lattice. Initial results relative to quenched wüstite are obtained using neutron diffraction by **W. L. Roth** who determines the octahedral and tetrahedral site occupancies of vacancies and Fe^{3+} cations, respectively represented by $(z+t)$ and t [39]. The characterization of the ratio $R = (z+t)/t$ was found close to $R=2$ for samples quenched from high temperatures. The neutron diffraction patterns obtained at 290 K and 4.2 K suggested the presence of defects constituted of two cation vacancies associated with one interstitial cation in the tetrahedral site. These two characteristics suggested a trend to the magnetite structure. In addition, the average magnetic moment per cation site was found to be much smaller than expected.

7.2. The historical Koch and Cohen cluster (13:4)

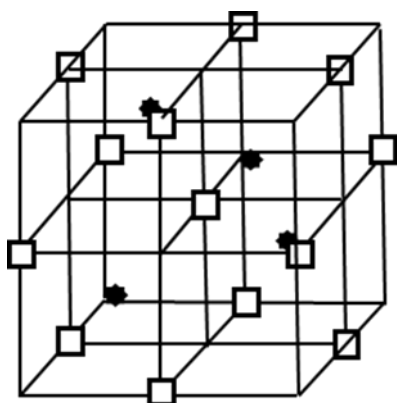


Fig. 13: The Koch and Cohen cluster (13:4) associated with a 3X commensurate structure valid for the sole wüstite $\text{Fe}_{0.902}\text{O}$ in Ref. [133].

To perform their structural refinements, the authors fix this cluster model assuming that the resulting superstructure should have a periodicity of 3X (the clusters being arranged only with distances of three cell parameters), in agreement with the known composition $z = 0.098$. The authors propose a new commensurate defect structure based on (13:4) large clusters (Fig. 13) distributed in the fcc lattice (cell parameter a) with a repetition distance ($3a$, $3a$, $3a$) noted 3X despite the fact that the additional peaks correspond to a superstructure 2.6X incompatible with their assumption.

For the first time, a crystallographic determination proposed a cluster corresponding to the corner sharing $\langle 110 \rangle$ type agglomeration of four basic clusters (4:1).

7.3. Neutron diffraction

Using *in situ* neutron diffraction experiments, **Cheetham, Fender and Taylor** [139] show that the ratio $R = (z+t)/t$ ranges between 3 and 4, with a mean value of 3.2, which can be compatible with a specific (13:4) cluster ($R = 3.25$). Using neutron diffraction at low and room temperature, coupled with magnetic scattering, **Battle and Cheetham** determine ratios $R = 2.78, 2.82$, and 3.01 and conclude that a cluster (8:3) can be consistent with their observations [140] and the preceding calculations of Catlow *et al.* (see ref. [11] Chapter 8 p. 407-410: summary of these results by Cheetham).

They couple neutron diffraction experiments to Mössbauer effect analyses and study the antiferromagnetic coupling of iron cations in octahedral and tetrahedral sites of the FeO lattice (see [140] Fig. 4 p. 343). For the first time, these authors show that atomic clustering observed in quenched samples (out of equilibrium) can be characterizing the study of magnetic properties of clusters in a lattice. This process will be used again by Saines *et al.* [141].

The numerous results obtained by **Radler, Cohen and Faber jr** from neutron diffraction patterns of polycrystalline samples under equilibrium yield values of $R > 3$ corresponding to mixtures of clusters at high temperature, and $R < 3$ to clusters larger than (10:3) below 1 173 K [142]. This observation could indicate a possible differentiation between the structures of W and W'.

Schweika et al. find a ratio of vacancies to interstitials $R = 4.0 (\pm 0.5)$ obtained under *in situ* conditions by neutron diffuse scattering from a single crystal with composition $z = 0.08$, at $T = 1\ 150\ ^\circ\text{C}$ [143, 144]. Their analysis of the diffuse scattering patterns leads them to conclude that 30% of the vacancies are free, while a further 15% of the defects are bound in isolated clusters (4:1). In their analysis, the authors state that their results invalidate the predominance of large clusters at elevated z . For example, the (13:4) cluster cannot fit the diffuse scattering that they characterized near the origin.

Saines et al. in Cheetham's team study by neutron diffraction a polycrystalline sample of wüstite $\text{Fe}_{0.902}\text{O}$ quenched from $900\ ^\circ\text{C}$ *i.e.* from the domain of W' [141]. They derive nuclear and magnetic clusters from an elaborate Reverse Monte Carlo modeling technique. The nuclear structure is described by islands of V4T units connected along the $\langle 110 \rangle$ directions « into a Koch-Cohen arrangement ». The magnetic structure in the bulk between the clusters is described as a « non-collinear variant » of the antiferromagnetic structure along the [111] axis as previously envisaged in [140]. Free vacancies are near the clusters. This structural approach can be compared with the paracrystalline description of the defect distribution by Welberry *et al.* [145] [146] [147].

7.4. X-ray diffraction

Welberry et al. present a detailed study of X-ray diffuse scattering for a single crystal with composition $z=0.057$, obtained after quenching from a temperature which is not specified. Their sample corresponds either to the Manenc's phase P (see [76] [77]), or to the W1 domain proposed by Vallet *et al.* [30][31][32]. The X-ray diffuse scattering analysis coupled with a paracrystalline modeling gives information on defect distribution, defect cluster size, number of interstitials, and lattice strain. The authors describe (statistically) a specific crystal constituted of ordered zones of clusters randomly distributed in a fcc lattice [145] [146] [147].

7.5. Gavarri et alii 's works

In 1979, using *in situ* neutron diffraction experiments, **Gavarri, Carel and Weigel** develop an analysis of the structural evolution of wüstite in CO_2/CO environments at high temperature [148]. A systematic determination of the ratio of the vacancy rate ($z+t$) divided by the interstitial rate t , $R = (z+t)/t$, is carried out. Quasi constant values of R close to 2.4 ± 0.4 , at two equilibrium temperatures (985

and 1 075 °C) and compositions z ranging between 0.058 and 0.120, are determined, with standard errors induced by the uncertainties on the separation between Bragg peaks and a complex diffuse scattering.

Fig. 14 displays the values of R obtained by different authors. In the case of studies made *in situ*, of which the ratio R is determined, most of the values of R range between 2.4 and 3.5. The authors observe two types of signals: (i) Bragg peaks from which the ratio R is determined, (ii) additional scattering corresponding to superstructure peaks, and diffuse scattering associated with defect clustering, disorder and static distortions (see [148][149] and *infra* Figs. 15 and 16).

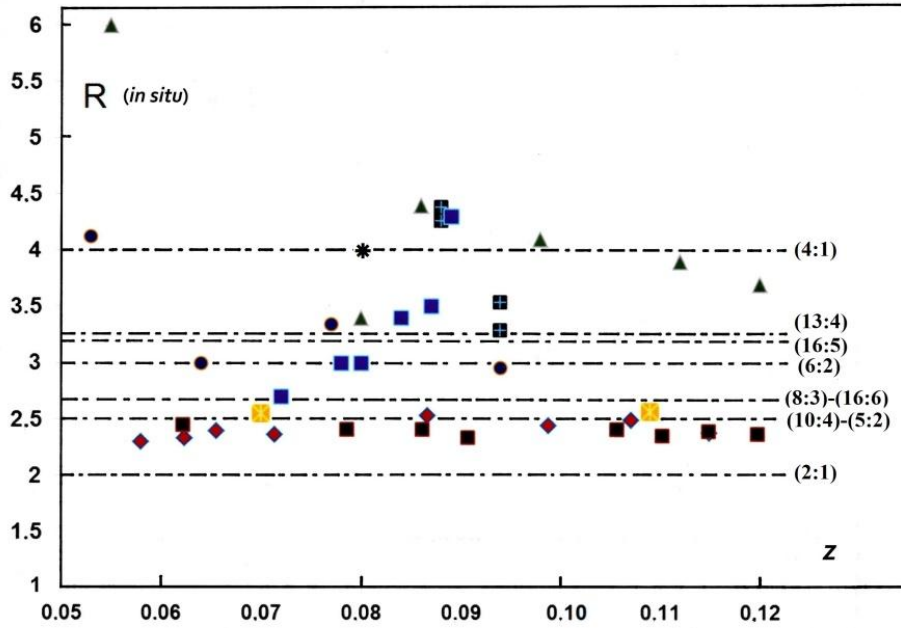


Fig. 14: $R = (z+t)/t$ values vs z , determined under equilibrium, from the literature.
 ■: at 900°C [134]; ●: at 800°C and ■: at 925, 1 050, 1 125, 1 180°C [139]; ▲: at 1 050°C
 and ■: at 850-1 100°C [142]; *: at 1 150°C [144]; ◆: at 1 075°C and ■: at 985°C [148].

The intensities $I_A(T, z)$ of this additional scattering increases linearly with composition z and decreases with temperature. The ratios I_A/I_{220} , where I_{220} is the intensity of the (220) Bragg peak, vary as follows:

$$I_A/I_{220} (985^\circ\text{C}) = 0.19 + 4.05(z-0.062) \quad (23)$$

$$I_A/I_{220} (1\,075^\circ\text{C}) = 0.15 + 4.25(z-0.058) \quad (24)$$

Therefore, these linear variations coupled with a small variation of the superstructure parameter k suggest a quasi-invariance of cluster size with clusters agglomerated in domains, and with the size of these domains increasing with z . The position of the centroid of this scattering allows the determination of an average superstructure corresponding to an irregular spacing of clusters quasi-constant with z or T . The average distance between clusters is evaluated to $2.7a_0$, a_0 being the cell parameter of ideal wüstite for a fixed temperature.

In later works by neutron diffraction on quenched samples of magnesio- and calcio-wüstites, using Rietveld refinements, Gavarrí *et al.* [148] [149] determined R values close to 2.5 for pure wüstites or for magnesio-wüstites $\text{Fe}_{(1-z-y)}\text{Mg}_y\text{O}$. In the case of calcio-wüstites $\text{Fe}_{(1-z-y)}\text{Ca}_y\text{O}$, as the composition y increased, the ratio R increased strongly with $R > 4$ and the superstructure vanished. Under these conditions, the cluster sizes would be strongly modified, and a majority of free vacancies would appear likely. These studies demonstrate that the formation of clusters is stabilized by Mg doping, and destabilized by Ca doping, producing free vacancies coexisting with residual clusters.

It is to be noted that the diffuse scattering due to long-range ordering is similar in diffraction patterns of quenched and *in situ* samples. Consequently, only a slight difference would exist between high temperatures and room temperature for the nature of defects and their organization in the fcc lattice. Superstructure and the structural disorder can be observed.

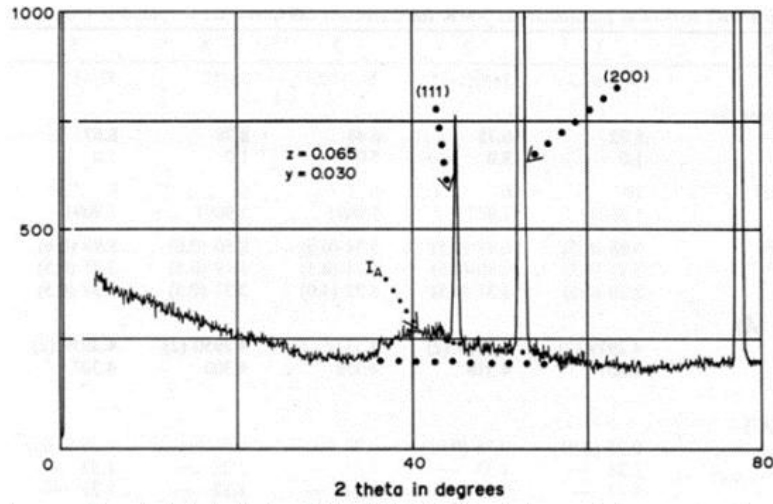


Fig. 15. Neutron diffraction pattern ($\lambda = 1.9 \text{ \AA}$) of a quenched polycrystalline calcio-wüstite $\text{Fe}_{1-z-y}\text{Ca}_y\text{O}$ in room conditions [149]. Presence of a large diffuse scattering (I_A) due to the superstructure (k close to 2.7), and small (110) peak ($36^\circ 2\theta$) due to vacancies and interstitial Fe^{3+} in the lattice. The intensity of I_A decreases with the calcium fraction y . Cell parameter $a(25^\circ\text{C}) = 4.3164 \text{ \AA}$. $R=(z+t)/t=3.3$.

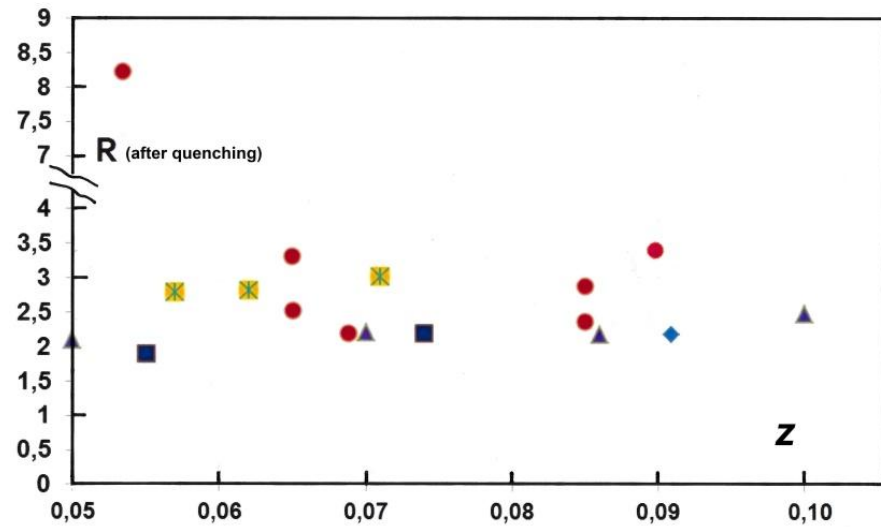


Fig. 16. R values vs z from literature determined after quenching.

Quenched Fe_{1-z}O : \blacklozenge : Gavarri, Carel, Weigel [148];
 WMg and WCa: \bullet : Carel and Gavarri [149];
 Quenched Fe_{1-z}O : \blacksquare : Roth [39]; \blacktriangle : Smuts [150]; \times : Battle and Cheetham [140]

Table 2 lists the values of R obtained in determination on quenched substituted magnesio- and calcio- wüstites, noted WMg and WCa respectively. The values of the ratio $R = (z+t)/t$ obtained from the literature on quenched, pure and Ca or Mg substituted, wüstites are displayed in Fig. 16. The utmost value $R = 8.2$ is obtained for a calcium rich calcio-wüstite ($y=0.05$).

Table 2.- $R = (z+t)/t$ values obtained from Rietveld analyses of neutron diffraction patterns of quenched substituted wüstites (see ref. [149]).

$\text{Fe}_{1-z-y}\text{Mg}_y\text{O}$	Fraction y (Mg)	$R = (z+t)/t$ (*)
$z=0.065$	0.075	2.53(0.5)
$z=0.085$	0.030	2.88(0.6)
$\text{Fe}_{1-z-y}\text{Ca}_y\text{O}$	Fraction y (Ca)	$R = (z+t)/t$
$z=0.068\ 8$	0.010	2.20(0.5)
$z=0.064\ 9$ (**)	0.030 (*)	3.31(0.5)
$z=0.053\ 4$	0.050	8.22(4.0)
$z=0.085$	0.010	2.37(0.5)
$z=0.089\ 9$	0.030	3.39(0.5)

(*) experimental uncertainty (due to diffuse scattering); (**) See Fig. 16.

7. 6. *T(ransmission) E(lectron) M(icroscopy) analyses*

In the past, microstructural analyses using electron diffraction and HREM images were performed (see **Andersson** and **Sletnes** [78]). The images obtained by **Iijima** for a sample $\text{Fe}_{0.92}\text{O}$ « provided by Professor Cohen » are often referenced [81]. The presence of the hypothetical cluster (10:4) <110> type with a ZnS-blende configuration has been particularly envisaged.

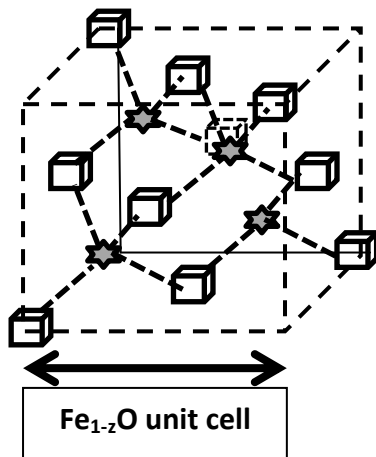
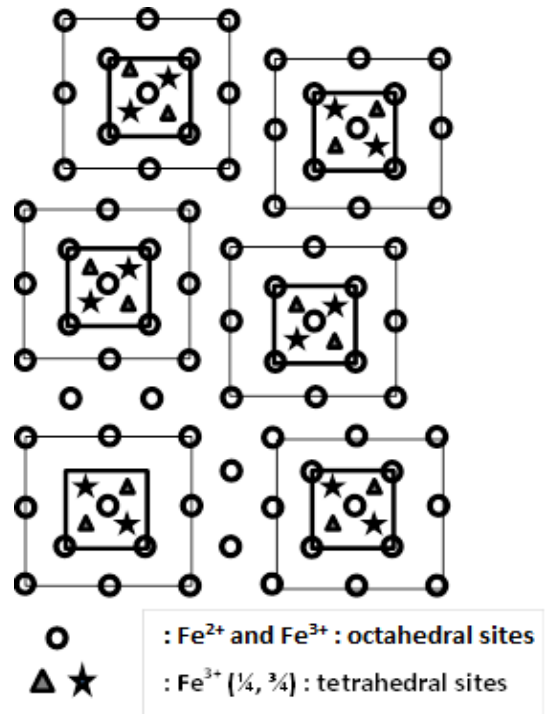


Fig.18 Various orderings of clusters (10:4) in the fcc lattice of wüstite. The basic cluster can be the (13:4), the (10:4) or other possibilities with 4 interstitial sites occupied by cations Fe^{3+} . The disorder can develop by statistical displacement in one, two or three directions. $2.55X$ could be the smaller intercluster distance (X = cell parameter of the “FeO” lattice).

Fig. 17. The cluster (10:4) ZnS blende <110> type.

★ = Fe^{3+} in tetrahedral site
 □ = Iron vacancy in octahedral site



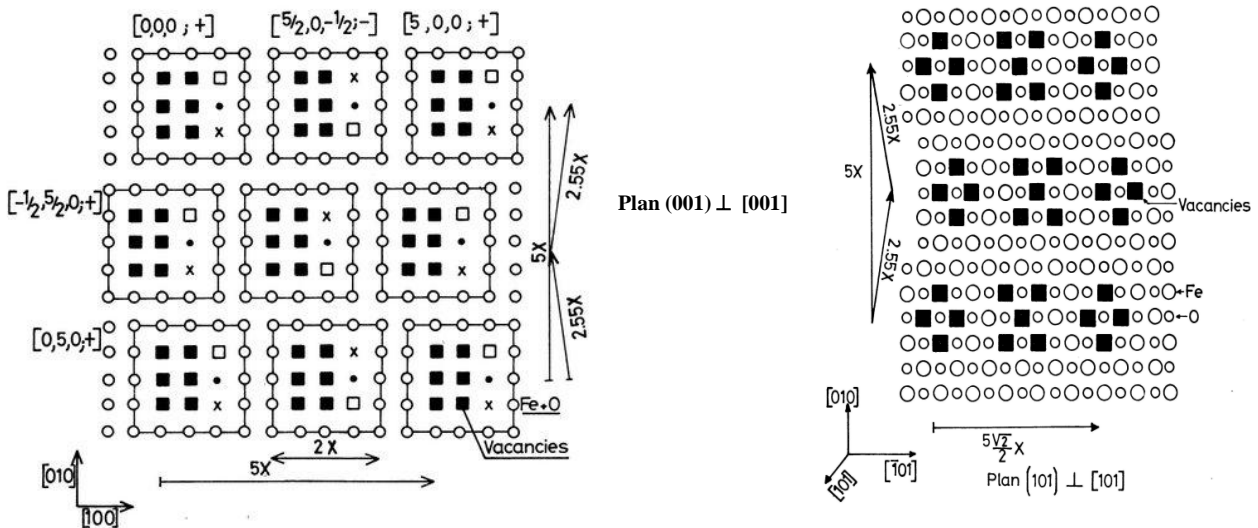
When re-interpreting the excellent images obtained from electron microscopy experiments by **Ishiguro** and **Nagakura** [151], Gavarri, Carel, Weigel [152], then **Nihoul**, Gavarri, Carel [153] show that this cluster corresponding to the ideal ratio $R = 2.5$ is compatible with the experimental data available (at least for quenched wüstites). This cluster depicted in Fig 17 is formed of a “FeO” cubic cell surround by an envelope of Fe^{3+} cations with total 3D dimensions ($2X \times 2X \times 2X$).

In the case of crystals quenched from 1 000 °C, with a composition $z = 0.10$ [152], the formation of local ordered domains of these clusters can be at the origin of all observations corresponding to a mean ratio R close to 2.5, and a quasi-invariant superstructure. The authors propose a progressive ordering of clusters as z increased: for low z values, repetition distances $3a$ linkable to pseudo phase P by Manenc [77] or W1 are observed. For intermediate z values, repetition distances $3a$ and $(2.55a \times 2.55a \times 2.55a)$ are formed in local zones linkable to W3 or P''.

A specific well-ordered superstructure can be based on a 3D supercell with parameters $(5a \times 5a \times 5a)$. This superlattice can be associated with the limit evolution at the boundary W3/Fe₃O₄. Such a superstructure was depicted at equilibrium by Shaïovitch, Zvinchuk and Vergazov [92]. Finally the distinction between « *allotropic* varieties of wüstite » or pseudo-phases can be due to both modifications of long-range order of clusters and the coexistence of different clusters that include free vacancies.

To define an isolated cluster correctly, particularly the ZnS blende defect cluster (10:4), it is necessary to take into account its envelope of Fe²⁺ and Fe³⁺ cations in octahedral sites. The charge of the cluster can be postulated as a first approximation by assuming the individual charges of vacancies and iron atoms. In Fig. 18 is illustrated the possible arrangement of these defect clusters (10:4), in the fcc lattice. The minimum distance between clusters is $[(2.5a)^2 + (0.5a)^2]^{1/2} = 2.55a$, and with allowed distances up to $3a$. This model is based on a cluster core (10:4) with double envelope of octahedral iron cations.

If we assume that a cluster (m:n) has a global charge $Q = 3n - mq$ (m vacancies with formal charges q ($=2, 1, 0$) and n interstitials Fe_i⁺⁺⁺ with charges $+3$), then the (10:4) cluster has a charge of $Q = -8, +2$, or $+12$ for $q = 2, 1$ or 0 , respectively. The electroneutrality is ensured by envelopes with charges $-Q$. Using this type of defect structure, different orderings of clusters such as the (10:4) one can be envisaged as shown in Fig. 18. A model of ordered clusters (10:4) is proposed which allows the simulation of HREM images in the two directions [001] and [101]. These hypothetical distributions are in agreement (Fig. 19) with the HREM images by Ishiguro and Nagakura [151] [152] [153].

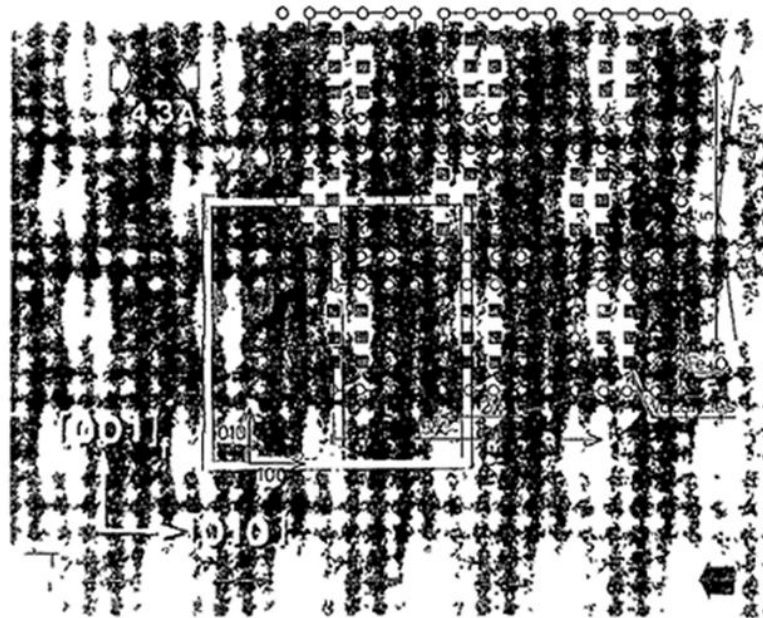


From crystallographic Ref. [152]

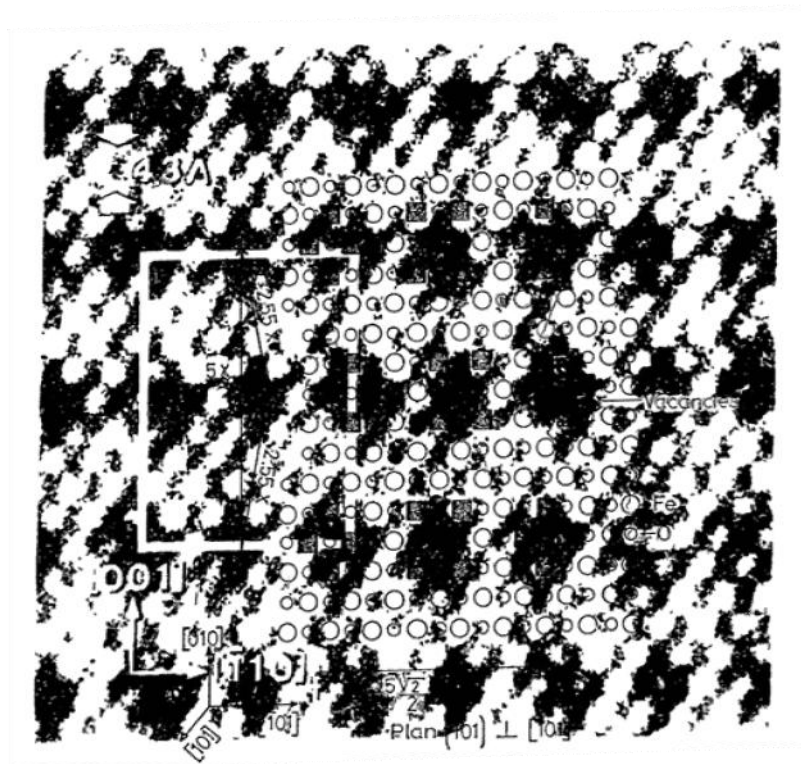
Fig. 19. Graphical modeling of local ordering of the (10:4) blende ZnS clusters in wüstite deduced by crystallographic considerations and drawn by identification with experimental HREM images by Ishiguro and Nagakura [151].

The graphical of HREM image (See ref. [151]) and corresponding graphical modeling [152] [153] are exactly superimposable:

A) Projection on (001) : Image HREM (Fig. 3 p. L724, Ref. [151]) and modeling (left part in Fig. 19)



B) Projection on (010) : Image HREM (Fig. 4 p. L724 in Ref. [151]) and modeling (right part in Fig. 19)



Symboles ■ for iron vacancies, ○ for Fe – O rows

8. MIXTURE OF DEFECT CLUSTERS AND/OR CLUSTER ZONES

8.1. Mixtures of defect clusters

Based on previous publications by Gavarri *et alii* [148] [154], a specific relationship can be established taking into account ordered zones of clusters and zones in which free vacancies are distributed. A mixture of different clusters $(m:n)_1$ and $(m:n)_2$ can coexist with additional free vacancies.

As a first approximation, the authors consider a lattice based on N cells of Fe_{1-z}O , with a proportion $V(z)$ of zones without defect. If N_i is the number of cells where clusters $(m:n)_i$ are formed, with the characteristics m_i and n_i of a specific cluster, and the distance $k_i a$ between two clusters, then the relation between cluster size and long-range order in the case of the fcc lattice of wüstite can be expressed as follows:

$$1-V(z) = 4z \cdot \sum_i \{ \pi_i (k_i)^3 / (m_i - n_i) \} \text{ with } \pi_i = N_i/N \quad (25)$$

For a unique cluster with continuous order (high z values), the relation (25) yields:

$$(m-n) / (k)^3 = 4z \rightarrow (m-n) = 4zk^3 \quad (26)$$

Considering the experimental values of k available in the literature, ranging between 2.5 and 3, this parameter k is characteristic of zones in which clusters are ordered partially. This means that, for low values of z , zones without clusters would exist and that they would be occupied by isolated defects *i.e.* free vacancies. The major part of the crystal lattice would be dominated by disordered point defects, and electrical conduction would be due to holes. For intermediate values of z , the cluster zones would be progressively joining together, forming percolation paths of ordered zones. A mixed conduction mode would be expected. For high values of z , the lattice would be occupied fully by ordered clusters, with likely small changes in the arrangement of clusters. Electrons in conduction band would mainly condition the electrical conduction.

This hypothetical description is in correspondence with the W1 disordered pseudo-phase and a conduction of p type, the intermediate pseudo-phase W2 as a mixture of cluster zones and free vacancy zones with mixed conduction, *i.e.* (holes +) and (electrons -) coexisting at short distance, and the ordered pseudo-phase W3 with conduction of n type from electrons in the conduction band.

8.2. Variation of parameter k

Numerous structural data can be found in the studies by **Bauer, Pianelli *et alii*** ([10] Part I and ref. [155] Part II). These authors develop a series of experiments using X-ray diffraction on powders of quenched wüstites. They observe linear variations of the cell parameter as a function of composition z and position of superstructure peaks. They characterize these peaks with parameter $\delta = 1/k$, where k designates the superstructure parameter discussed previously. Parameter δ varies between 0.37 and 0.39, corresponding to values of k between 2.7 and 2.56. This variation observed in quenched samples is in good agreement with the model presented in (section 7.6. *TEM analysis*) with a minimum k value of 2.55. Authors ([10] [155]) consider that their data did not permit to observe any phase transition. However, they suggest that wüstite can be described by an incommensurate lattice of clusters, with the presence of different clusters based on cluster (4:1).

8.3. Elementary statistical analysis of diffuse scattering

Gartstein and Mason [108] show that p conduction is observed in wüstite and that electrical conduction is based on a hopping mechanism between adjacent clusters. Later, Gartstein *et alii* publish a crystallographic analysis of a $\text{Fe}_{1-0.07}\text{O}$ single crystal using X-ray diffraction at 1 173 K ([134] [156]). They develop a detailed analysis of diffuse scattering and Bragg peaks, and conclude to the existence of a mixture of clusters (13:4) and asymmetric new clusters (5:2) constituted of 5 vacancies and 2 interstitials ($R = 2.5$). They propose an ordering with an average repetition distance of $2.5a$, close to the hypothetical value of 2.55 (see [153] [154]).

8.4. Summary

To conclude this section, it can be noted that there is an inherent difference between the concept of a mixture of disordered clusters with various forms, and the concept of zones with identical disordered clusters, coexisting with zones of free point defects.

The concept of incommensurability, applied to specific compositions z , was also proposed as a mathematical approach of disorder (see Yamamoto [157] and Weigel *et al.* [158]). Finally, the proposition of the irregular cluster (5:2), resulting from the analysis of diffuse scattering by Gartstein *et alii* [134] and Radler *et al.* [142] seemed to be compatible with the quasi-constant value $R=(z+t)/t = 2.4 \pm 0.4$ [148].

9. DEFECT CLUSTERING AND EQUILIBRIUM EQUATIONS

In this section, using basic equilibrium equations renewing at the same time by Stokłosa in Ref. [25] the link between thermodynamic and defect structure (see particularly [1] [11] [20] [23] and [48]), a new presentation of defect clusterings for the wüstites is proposed.

It is important to remind that, in $Fe_{1-z}O$ where only Fe^{2+} (also noted Fe_{Fe}) and Fe^{3+} cations (also noted Fe_{Fe}° or $Fe_{Fe}^{\circ\circ}$, for cations in octahedral or tetrahedral sites respectively) can coexist, cation vacancies can have one of three possible stable charges: doubly charged vacancies V_{Fe}'' , simply charged vacancies V_{Fe}' , and neutral vacancies V_{Fe}^{\times} (noted V_{Fe}). Singly charged iron vacancies V_{Fe}' were proposed previously in place of doubly charged vacancies V_{Fe}'' by Kofstad and Hed [120]. Vacancies of the oxygen sublattice are not considered because they occur in negligible numbers compared to the iron vacancies.

In section 2.2. (*Modifications or pseudo-phases*) the equilibrium constant (4b) of the hypothetical equilibria (4a), in which the activity of oxygen O_2 in the solid is taken equal to 1, can be developed with $q=0, 1$ or 2 following the general expression in which $s=2(q+1)$

$$\log z = (1/s) \cdot l' + (1/s) \log (K/C)^2 \quad (\text{with } C = (q)^{2q}) \quad (27)$$

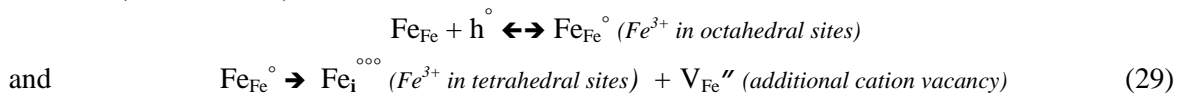
which is identical to equation (4c).

The three hypothetical values of s corresponding to $q = 2, 1$, and 0 are $s = 6, 4$, and 2 respectively.

The modifications of q can be described in terms of two equilibria involving holes h° in the valence band



Therefore, three domains corresponding to the three types of iron vacancies can exist in all systems presenting significant cation deficiency. In the case of $Fe_{1-z}O$, additional reactions occur since two valences (Fe^{2+} and Fe^{3+}) exist:



As a first approximation, assuming that the sole point defects are iron vacancies, associated with Fe^{3+} (Fe_{Fe}°) and coexisting with Fe^{2+} (Fe_{Fe}) in $Fe_{1-z}O$, the general relationship between composition z and oxygen partial pressure p' is as follows

$$q Fe_{Fe} + \frac{1}{2} O_2 \leftrightarrow q Fe_{Fe}^{\circ} + V^{q(\prime)} + O_2 \rightarrow K_q = [Fe_{Fe}^{\circ}]^q \cdot [V^{q(\prime)}] / ([Fe_{Fe}]^q \cdot p'^{1/2}) \quad (30)$$

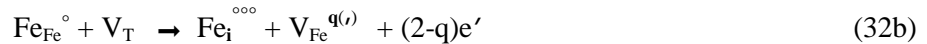
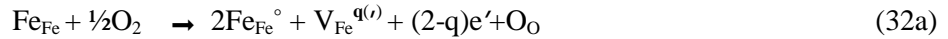
In expression (30) $[Fe_{Fe}] \approx 1-3z$ and $[Fe_{Fe}^{\circ}] \approx 2z$.

As all molar fractions are related to the composition z , the next relation can be derived

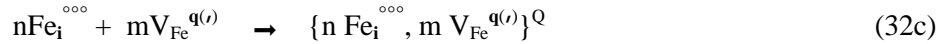
$$K_q = [\text{Fe}_{\text{Fe}}^\circ]^q \cdot [\text{V}_{\text{Fe}}^{q(\cdot)}] / ([\text{Fe}_{\text{Fe}}]^\circ)^q \cdot p'^{1/2/2} = 2^q \cdot z^{(q+1)} / (1-3z)^q \cdot p'^{1/2} \quad (31)$$

If this simple model is correct, then the term $(1-3z)$ should vary slowly with z in a reduced composition range, meaning that the activity of Fe_{Fe} can be considered as being constant. Under this assumption, z is roughly proportional to $(p')^{1/s}$, with $s = 2(q+1)$ being an integer characteristic of the nature of the defects. For $q = 2$ (doubly charged vacancies) we obtain $s \approx 6$. For $q = 1$ (singly charged vacancies), $s \approx 4$. For neutral vacancies, $s \approx 2$.

Now, let us consider complex clusters $(m:n)^Q$ with variable formal charge Q and successive reactions in which vacant tetrahedral sites V_T are introduced:



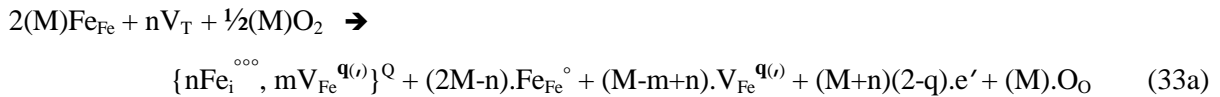
The formation of stable cluster can be written:



with $Q = 3n - qm$ (positive or negative charge). The cluster will be noted: $\text{CL} = \{n\text{Fe}_i^{\circ\circ\circ}, m\text{V}_{\text{Fe}}^{q(\cdot)}\}^Q$.

The formal charge Q of a cluster is assumed to result as a first approximation from individual charges of vacancies with the invariant charge of interstitials Fe^{3+} ($\text{Fe}_i^{\circ\circ\circ}$). However, the interstitial Fe^{3+} can be modified into Fe^{2+} interstitial cation due to electron jumping, and as a result a more complete model should take into account various possible Q values of cluster charges due to the modification of interstitial cation charges.

Considering the formation of clusters CL coexisting with point defects $\text{Fe}_{\text{Fe}}^\circ$ and $\text{V}_{\text{Fe}}^{q(\cdot)}$, a more elaborate equation (32a) can be expressed as follows:



with the hypothesis: $Q = 3n - qm$ (positive or negative).

In equation (33a), the integer $M \geq m$ is introduced arbitrarily to take into account the possible presence of additional free species $\text{Fe}_{\text{Fe}}^\circ$ and $\text{V}_{\text{Fe}}^{q(\cdot)}$, associated with free electrons in the conduction band. The additional vacant tetrahedral sites are the vacant sites V_T able to receive an interstitial cation $\text{Fe}_i^{\circ\circ\circ}$.

As a major fraction of these sites (8 sites for 4 FeO units in a fcc cell) is vacant, their activity can be considered as being quasi-constant, and does not play a significant role in the calculation of equilibrium constants. The resulting equilibrium constant is as follows, with the reduced notation $[\text{CL}] = [\{n\text{Fe}_i^{\circ\circ\circ}, m\text{V}_{\text{Fe}}^{q(\cdot)}\}^Q]$:

$$K(T) = [\text{CL}].[\text{Fe}_{\text{Fe}}^\circ]^{(2M-n)}.[\text{V}_{\text{Fe}}^{q(\cdot)}]^{(M-m+n)}.[e']^{(M+n)(2-q)} / ([\text{Fe}_{\text{Fe}}]^{2M}.\text{V}_T^n.(p')^{M/2}) \quad (33b)$$

In this expression, the cluster concentration $[\text{CL}]$ is considered as being proportional to z at least as a first approximation. Applying the same approximation as in previous equilibrium equations, the equilibrium constant can be expressed as follows:

$$K(T) = z.z^{(2M-n)}.z^n.z^{(M-m+n)} / ((1-3z)^{2M}.(p')^{M/2}) \approx z^{(2m+1+(2-q)(m+n))} / (1-3z)^{2m}.(p')^{m/2} \quad (33c)$$

No activity coefficient is taken into consideration. Consequently, in a restricted range of composi-

tion z , no significant change of interactions occurs between point defects or clusters. To illustrate the possible values of parameter s corresponding to the three hypothetical q values, l' has been simulated as a function of $\log z$ using the expressions of $K(T)$. The slopes of the curves (l' vs $\log z$) were determined. The two specific clusters (4:1) and (10:4) have been chosen to calculate the various coefficients s depending on the charges q of vacancies.

The results are:

- for a cluster (4:1), the three coefficients s are equal to 5.7, 8.2 and 10.7 for $q=2, 1$ and 0 , respectively;
- for a cluster (10:4), the three coefficients s are equal to 5.4, 8.2 and 11 for $q=2, 1$ and 0 , respectively.

It can be noted that, for either the cluster (4:1) or the cluster (10:4), the coefficients resulting from the calculations are quite similar for a given value of q . Given the approximations included in the model and expected experimental uncertainties, these simulated values (from $s = 5$ to 11) can be considered as being congruent with the experimental s values obtained in the literature and ranging between 4 and 10 . In other terms, the existence of pseudo-phases might be reasonably thought to be associated with the variation of the electrical charges (q -) of cation vacancies, and more generally with the variation of the electrical charges Q of the clusters.

These changes of electrical charges are strongly correlated with the increasing density of them. In other terms, two neighboring vacancies V'' with the same negative charge $2-$ can produce two singly charged V' or neutral V^x vacancies, which will decrease the coulomb repulsion between them, and require additional free electrons migrating in the conduction band.

Now, using this general approach delivering specific s values, it is possible to justify the various modifications observed by many authors including Vallet *et al* [30] [31] [32], Geiger, Levin and Wagner [67], Fender and Riley [69], Takayama and Kimizuka [72], Worral and Coley [75], Stokłosa *et alii* [25] [85]. In particular, let us recall the values $s(W1) = 4.6$, $s(W2) = 6.2$, $s(W3) = 9.2$, extracted from the data by Vallet *et al.*, as well as values $s(W1) = 4.8$, $s(W2) = 6$, $s(W3) = 9$ extracted from the data by Takayama and Kimizuka [72] (see section 3.7. Takayama and Kimizuka's *phase modifications*).

Considering the set of values $s = 4.6, 6, 9$ associated with $W1, W2$ and $W3$ respectively, for low values of z , the existence of doubly charged iron vacancies V_{Fe}'' can be associated with a major population of holes corresponding to mode 'p' of conduction. For high values of z , the existence of singly charged or neutral vacancies can be associated with mode 'n' of conduction, *i.e.*, to the injection of electrons in the conduction band. For intermediate values of z , an intermediate state would exist with a mix of holes in the valence band and electrons in the conduction band, following a rate varying from $z_{1/2}$ to $z_{2/3}$.

These considerations can be supported by the p-n transition [103] described many times in the literature (see Fig. 10). The authors provide a detailed interpretation for the existence of various defect configurations corresponding to pseudo-phases in the entire wüstite phase diagram.

For mean values of z , neutron diffraction data show that the clusters are assembled in limited ordered zones. For high values of z , the ordered cluster zones percolate (see [151] HREM images for $z=0.08$ and $z=0.10$, and **Addenda 1- Modeling** p. 40). The change in vacancy charge as z increases can be accompanied by injection of free electrons in the conduction band. As temperature T increases, the size of the cluster zones decreases, and the number of cluster zones increases. Correlatively, the population of free defects increases.

Taking into account all these considerations the $W1$ pseudo-phase can be associated with a small number of clusters assembled in cluster zones, and a large population of free defects with numerous free V_{Fe}'' vacancies. In this domain, holes can be the major electric charge carriers. The $W3$ pseudo-phase is

constituted of large cluster zones, covering the entire lattice adjacent to the magnetite zone. The neutralization of vacancies (or decreasing q charges) shall be associated with a large number of electrons occupying the conduction band. As the W2 pseudo-phase was previously considered as a mix system of pseudo-phases W1 and W3 by several authors, an intermediate structural configuration can exist, corresponding to the formation of percolation paths allowing a mix conduction of holes and electrons.

10. PERCOLATION APPROACH

A primitive « composite picture » of a possible evolution of the defect structure needing only two clusters and distinct populations through the subdomains was proposed earlier [126]. Nearly at the same time, a similar simple model of the defect structure was proposed in terms of the so-called theoretical Cluster Component Method (see Section 6.3. *The C. C. M.*, Fig. 12 p. 22).

In paper [21], p. 376 Gleitzer envisages a distribution of the charges and the clusters following the percolation model. That can be basically represented from Fe/W to W/Fe₃O₄ boundaries. as following:

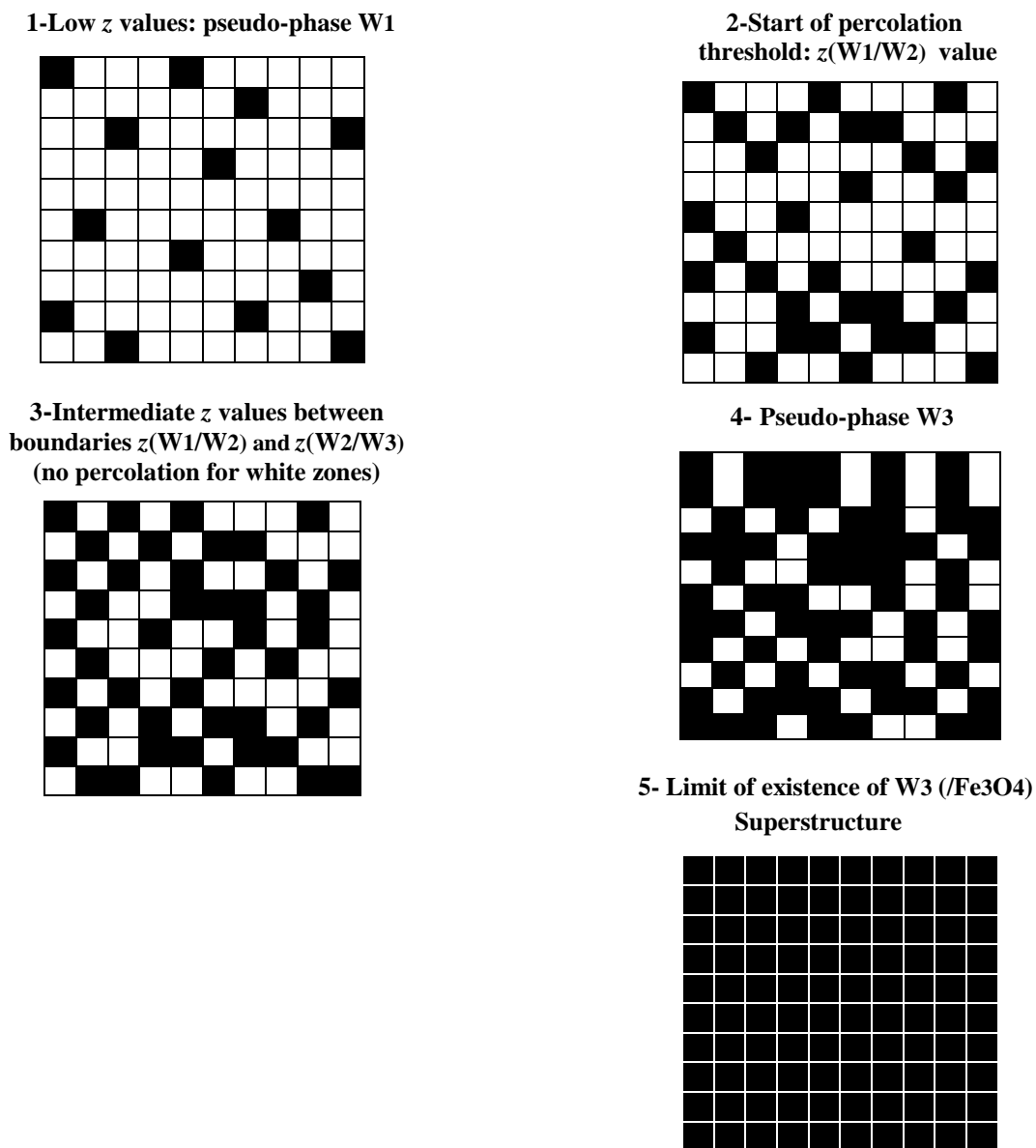


Fig. 20. Three configurations for CLZ (in black) disordered distributions, typically at 1 000 °C from Fe/W1 to W3/Fe₃O₄: W1 = majority of zones without defect clusters; W3 = majority of zones with defect clusters; W2 = intermediate configuration.

In this figure a representation of the distribution of the cluster zones (CLZ) inside the lattice is drawn as a schematic percolation process. The size of cluster zones (CLZ) (represented by black squares) is conditioned by the equilibrium temperature. White squares represent zones with only disordered point defects (free vacancies, individual Fe^{3+} in octahedral or tetrahedral sites).

- In the W1 domain, the CLZ's are isolated and immersed in a defect lattice. The transition line W1/W2 corresponds to a specific threshold $z_{1/2}$ value (for fixed T) where there is a starting percolation state.
- The W2 domain might be constituted of large black zones, with percolative paths through the lattice or both types of zones. The W2/W3 transition line might be reached when the residual white zones are no more percolative.
- The W3 state is reached when black zones envelop white zones up to a limit state where all the lattice is covered by a continuum of clusters, *i.e.*, at the boundary W3/Fe₃O₄.

The effect of temperature is to modify the sizes of clusters zones and the limits of composition for each boundary. As temperature increases, the cluster zones are destabilized, and the size of these zones decreases. Because of the complexity of the system, it is obvious that no unique specific model can represent the various evolutions. Nevertheless, this approach shows that the modifications W1, W2, W3 may be justified simply from changes in electrical charges of point defects, so of clusters, and are strongly correlated with composition z , and the degree of ordering. This scheme implies nearly continuous transformations from W_i to $W_{(i+1)}$ at the level of the slight deviations from the main diffraction process. However, the discontinuities observed in the partial molar properties [see *supra* Fig. 6 p. 11: $\bar{H}_{\text{O}} - \frac{1}{2}\bar{H}(\text{O}_2)$, $\bar{S}_{\text{O}} - \frac{1}{2}\bar{S}(\text{O}_2)$, and Fig. 7 p. 12: $\bar{H}_{\text{Fe}} - \frac{1}{2}\bar{H}(\text{Fe Cr})$] and in the kinetic reduction rates by [75] at the compositions x_{ij} of boundary of the subdomains correspond to abrupt and small-scale differences in internal energy and configurational entropy, due to the change of cluster type (or charge), and to the change in long range order of the clusters (see also [115] [116] in Annex 3-Fig. A3 p. 39).

Note: Metastabilities encountered in most of the experimental studies can modify the location of the transitions $i \rightleftharpoons j$, *i.e.* the characteristic z_{ij} values as determined by thermodynamic assessments can even suppressed them. As examples the thermogravimetric isotherms (W area) in Raccach's study [42] (see V.-ADDENDA 1- Metastability Fig. a) and b) p. 17 *hal* Rennes: 0341 75 62 (11-05-2021) and Cornell: arxiv.org/abs/2111.04346. English translation) are grouped together considering a sole modification in certain temperature intervals, without break all along from $\gamma\text{-Fe}/W_i$ to $W_i/\text{Fe}_3\text{O}_4$: W3 at 920, 950, 1 000 °C, W2 at 1 067, 1 122, 1 140 °C, W1 at 1 160, 1 167, 1 200, 1 250 °C.

11. CONCLUSIONS

Point defects: Up to now, the subdivision of the equilibrium phase diagram of wüstite appears to be complex with the existence of several domains corresponding to different thermodynamic behaviors and high order transitions (at least of second order). The departure from stoichiometry can be described by a complex system of cation vacancies noted V'' , V' or V^x , depending on their electrical charge or neutrality, giving rise to $\text{Fe}_{\text{Fe}}^\bullet$ (cations Fe^{3+}) coexisting with Fe_{Fe} (cations Fe^{2+}) in octahedral sites, and $\text{Fe}_i^{\bullet\bullet\bullet}$ (cations Fe^{3+}) occupying the tetrahedral sites of the NaCl-type lattice. Changes in electrical charges must be taken into consideration, with electrons in the conduction band and holes in the valence band.

Clusters: Works from several different sources demonstrate the existence of clusters of defects constituted of iron vacancies coupled with cations Fe^{3+} in tetrahedral interstitial sites, and mixture of $\text{Fe}^{2+}/\text{Fe}^{3+}$ cations in octahedral sites of the basic fcc lattice. These clusters can coexist with free vacancies. However, at present, the exact form of clusters noted (m:n) (m vacancies $V^{q(n)}$ linked to n interstitials $\text{Fe}_i^{\circ\circ\circ}$) is far from being determined definitively.

The ratio $R = (z+t)/t$ is obtained mainly from neutron diffraction experiments in the range 2 to 4 depending on the authors. Several authors observe quasi-constant R values in the range 2 to 3, while others observe a variable or constant value in the range 2.8 to 4. It can be noted that if the wüstite lattice is occupied solely by an $(m:n)$ isolated cluster, then the ratio R is equal to $(z+t)/t = m/n$.

The pseudo-phases: inside the initial phase diagram some experimental evidences from different authors support the existence of distinct areas corresponding more than likely to pseudo-phases W1, W2, W3.

The term « boundary » designates a threshold of percolation at the exact separation between two thermodynamic behaviors, corresponding to progressive evolutions of short- and long-range ordering. An additional first order transition $W \rightleftharpoons W'$ is directly linked to the first order transition $\alpha\text{-Fe} \rightleftharpoons \gamma\text{-Fe}$ at 912 °C.

Up to now, the lack of systematic studies relating to the isotherm at 912 °C limits the correct definition of this part of the Fe-O phase diagram, despite the availability of some advanced thermodynamic studies. The interpretation of them remains to be developed.

Changes in cluster form: with respect to the existence of clusters $(m:n)$ and the variation of the ratio R observed by some authors, a first description of the pseudo-phases W1, W2, W3 is proposed in terms of structure changes of these clusters, with an increasing cluster size as z increased.

The ZnS-blende type (10:4) cluster can agree with neutron diffraction experiments at equilibrium and in quenched samples, and with electron microscopy images for some compositions z of quenched samples. The (12:4) $\langle 110 \rangle$ specific cluster, proposed by Lebreton and Hobbs [129], can also agree with a superstructure of P'' type ($5a \times 5a \times 10a$) for a sufficiently high z composition.

The existence of separated ordered cluster zones can agree with the slowly variable k parameter. Let us note that significant changes in cluster form (*e.g.*, (13:4) or (10:4) transformed into (4:1)) as z or T vary are not demonstrated for the entire phase diagram. In addition, in studies of samples at equilibrium or after a quenching process, all observations of superstructure Bragg peaks or diffuse scattering due to disorder show that the k value (k close to 2.6-2.7) varies slowly with z , thus characterizing a quasi-constant repetition distance between clusters. If the form of clusters $(m:n)$ would change strongly, the repetition distance $k.a$ would change also strongly, which is not observed.

New possible descriptions of the pseudo-phases: the pseudo-phase W1 (low z values and high T values) would be characterized likely by free vacancies coexisting with a weak proportion of small zones of clusters. The pseudo phase W3 (high z values and low T values) would correspond to clusters covering the whole lattice with a minimum distance of $2.55a$ between clusters and to a $5X$ superstructure resulting from a zigzag ordering of clusters. The pseudo-phase W2 is identified as a solution of W1 and W3 of opposite energetic evolutions above its critical temperature near 300 °C. This pseudo phase is likely constituted of zones of ordered clusters and zones of free vacancies, with a progressive build up from boundary $z_{1/2}$ to boundary $z_{2/3}$ at which the junction arises between cluster zones that percolate.

The p to n transition might be associated with a starting of this percolation. This transition could be located in the W2 domain in a non-exclusive manner and is likely connected with a given change in electrical charges of vacancies and clusters.

Finally, an additional feature can explain the observed electrical modifications. As z increases, the electrical charges of vacancies ($V^{q(\cdot)} = V'', V'$ or V^x) and clusters could change with q varying from $q=2$ (low z) to $q=0$ (high z), due to interactions between charges. This modification of charges q is associated with an increasing number of electrons populating the conduction band.

Transition close to 912 °C: Up to now, the three varieties W'1, W'2 and W'3 associated to $\alpha\text{-Fe}$ below 912 °C were not strictly characterized from a structural point of view. According to Worrall and

Coley in their isotopic kinetic study, the defect structure identified on the surface of polycrystalline samples appears as being close to that identified for the pseudo-phases W1, W2, and W3 [24] [75].

However and most than likely the W_i and the W'_i being identified thermodynamically on both parts of the transition $\alpha\text{-Fe} \rightleftharpoons \gamma\text{-Fe}$ as soon as 1962 got specific defect structures in accordance with the pseudo phase diagram in a recent systematic description published in 2015 (see ADDENDUM 3, p. 41).

ANNEX 1 - Numerous sets from the literature of transitions located in diagram $\Theta(x)$

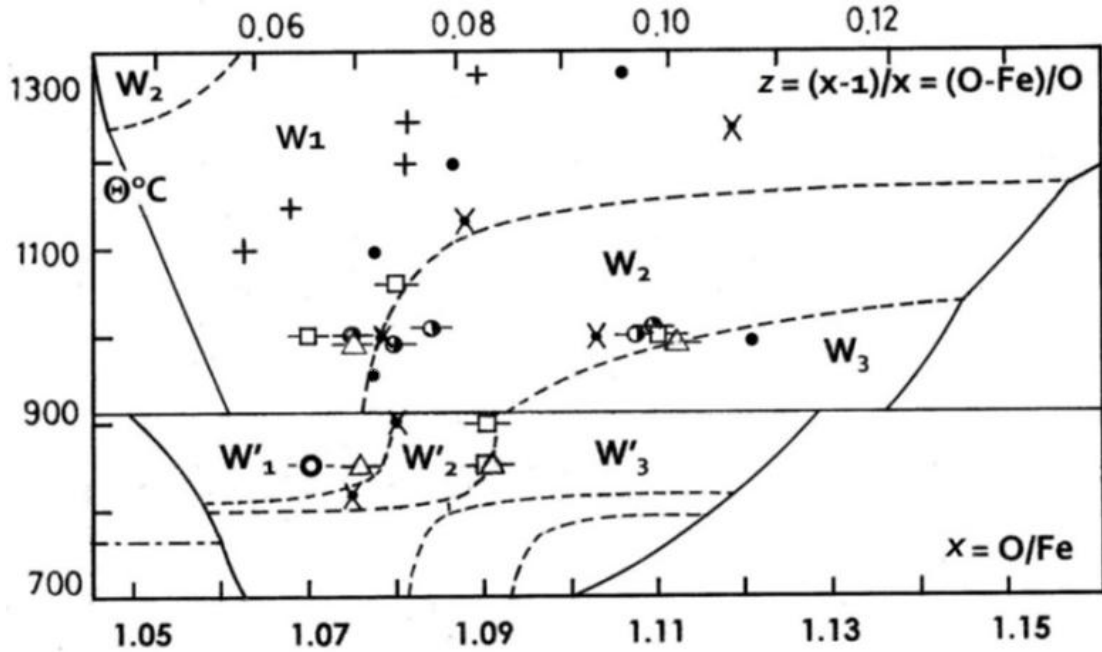


Fig. A1. Transitional effects located in the diagram $[\Theta^\circ\text{C}, x]$.

- a- Geiger, Levin, Wagner jr in ([67] p. 955 Table 4): “unexpected changes in the slopes of $\log \sigma$ vs $\log p(\text{O}_2)$ ” $\Rightarrow \square$
- b- Bransky, Hed [100]: breaking points at 1 000, 1 100, 1 200, 1 300 °C re-calculated for correlations of l' vs z $\Rightarrow \bullet$
- c- Riecke, Bohnenkamp in ([118] p.722 Fig 7): reaction rate vs $100y$ (Fe_{1-y}O) at 1 000 °C, $p_{\text{H}_2\text{O}}/p_{\text{H}_2}$ re-analyzed curves
(I), (III) oxidation, curve (II) reduction [(I): +10°shift, (III): -10°shift]: change in slope of the quasi-linear segments $\Rightarrow \bullet$
- d- Takayama, Kimizuka in [72]; see *supra* re-analyze of the isotherms in §3.7. « phase modifications » $\Rightarrow +$
- e- Molenda, Stokłosa, Znamirowski in ([85] p. 517; Fig. 2 and p. 519): experimental $\log(\sigma)$ vs l' (818-1 307 °C), successive segments defined from Raccach-Vallet’s data to model the curvature; Fig.4: isothermal $\sigma(y)$, resulting intersections of successive linear segments of $\sigma(y)$ $\Rightarrow \times$
- f- Worral, Coley in [75]: another interpretation of Fig. 9 at 850 °C (p. 821) give rise to an additional point close to $x = 1.078$ as a second order type transition, identified for $W'1 \leftrightarrow W'2$ $\Rightarrow \triangle$

Another transition for $x = 1.090$ appears as a discontinuity at the location of the transition $W'2 \leftrightarrow W'3$.

Point \bullet located close to $x = 1.070$ approximates the common intersection of the three lines drawn from re-interpreting the kinetic data at 850 °C (see Annex B2) $\Rightarrow \bullet$

See more numerous sets (W_i/W_j and p-n transitions) displayed in the diagram $T(\delta)$ ($\delta \equiv z$) by Stokłosa in Ref. [25] Fig.12.17. p. 339.

See also the tentative extension of the thermodynamic diagram to 5 varieties identified in MnOx :

ADDENDUM 2 (below p. 41): **On a likely drawing of the phase diagram of manganese monoxide.**
A very short review.

2 - Two first experimental evidences of the $\Theta(x)$ phase diagram with three subdomains (1964)

The first experimental confirmation of the three W_i out of Vallet's team is published by [67]. At the same time, the work of Landler and Komarek [115] interpreted in ref. [116] becomes available (see below **Section 3 -**).

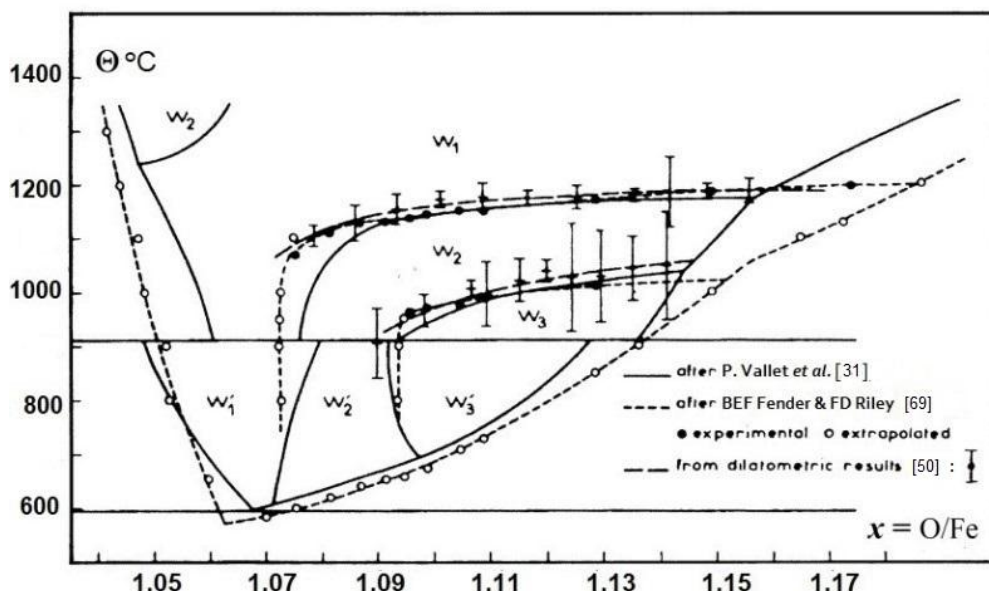


Fig. A2. Comparative plot of experimental phase diagram versions (see ref. [52], Fig. 3 p. 247) by Vallet and Raccach [30] and Fender and Riley [69]

3 - Thermogravimetric study of the reduction in H₂/H₂O mixtures. Separate sequences

A numerical analysis makes be revealed three diffusion sequences distributed across the analyzed domain by Landler and Komarek [115], [116]. Fig. A3 below is drawn as a summarized result of the determinations.

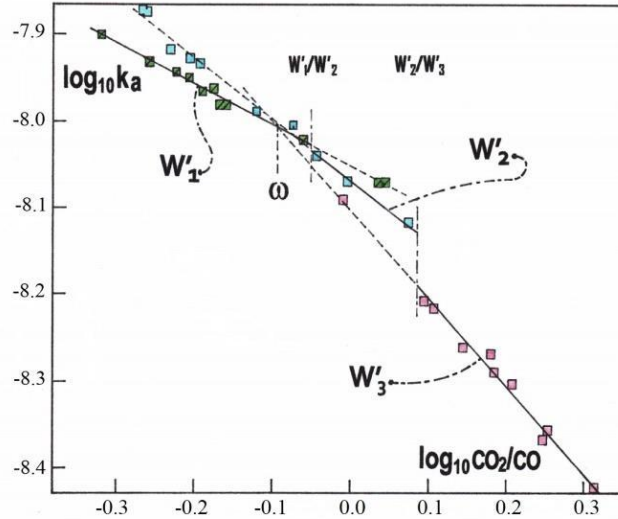
See Figure in Ref. [116])

Fig. A3. Typical runs of reduction in H₂/H₂O gas mixtures under 1 atm
(identified in Landler and Komarek Ref. [115])

39

4 - Re-analysis of kinetic data related to the isotherm at 850 °C by Worrall and Coley Ref. [75]

In Fig. 5 p. 819 and Fig. 9 p. 821 (Ref. [75]) it is possible to separate three linear variations of $\log k_a = -m \cdot \log(\text{CO}_2/\text{CO}) + \log k_0$ in place of only two variations. It results the figure below in which there are three linear segments possibly attributed to W'_1 , W'_2 and W'_3 , which converge at point ω for which $x = 1.0698$, a practically identical value to that of Chaudron 's point.



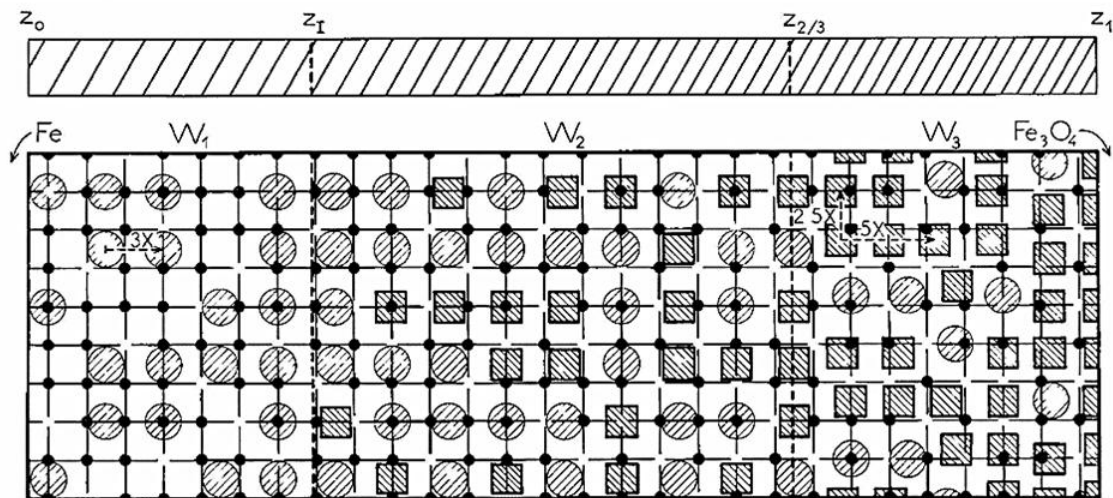
(identified in Worrall and Coley's paper, Ref. [75])

Fig. A4. Intersection W'_1/W'_2 for $\log \text{CO}_2/\text{CO} = -0.05052 \Rightarrow x=1.0776$, expected 1.0779 in the bulk;
 W'_2/W'_3 for $\log \text{CO}_2/\text{CO} = -0.08468 \Rightarrow x=1.0903$, expected $x=1.0901$;
point ω : $x = 1.0698$ where the three variations are convergent.

Let us observe that the variation related to W'_2 is placed at intermediary in the angle $[\log k_a(W'_1), \log k_a(W'_3)]$. The L.S. coefficients $m = 0.48 \pm 0.03$ (W'_1 ?), 0.71 ± 0.05 (W'_2 ?), 1.03 ± 0.09 (W'_3 ?) are statistically identical to those of equations ([36]), ([37]), ([38]) at 1268 K. This identity should be explained by surface phenomena not by bulk transitions following Ref. [75] p. 820.

ADDENDA (in 2023)

1 – Modeling: Composite Picture proposed in 1980 in a Round Table Meeting in Krakow-Poland.



See Figure p. 58 in Ref. [126]

An initial graphical representation of trends to ordering of defects with increasing z was presented at the First Round Table Meeting on Transport, Structure and Phase Transformations of Iron and Manganese Oxides, Krakow-Kozubnik, Pologne (8-9 Sept. 1980) in the hypothetical case of two kinds of clusters distributed in a planar lattice. This modeling was developed and published (see Ref. [154]). The Proceedings of the meeting are published later (see Ref. [126]). In 1982 Men' proposed an improved modeling based on the theory of CCM (see above p. 22 § 6.3 Fig. 12, and Ref. [122]) using the present structural suggestion.

2 – On a likely drawing of the phase diagram of manganese monoxide, C. Carel, C. R. Acad. Sc. Paris 295 (1982) Série II p. 853-56. A short review.

3 – Identification of the cluster types combined with the pseudo phases (A. Stokłosa, Ref. [25], Chap. 12, pp. 65)

As still referenced particularly with Smyth [22], Mrowec and Podgoreka [16], Stokłosa *et alii* [85], these previous authors make progress with the new thermodynamic and electronic approach making connections with the defect structure of parent oxide mixtures, then with pure wüstite. The pseudo phases appear progressively in their work.

Particularly in Ref. [85] by Molenda, Stokłosa, Znamirowski the three wüstites and their transitions are revealed (See Table 1 p. 520, and Fig. 4 p. 522, $\sigma(\Omega^{-1} \cdot \text{cm}^{-1} - y (\equiv z))$, after Wagner jr *et al.* [67], and Fender and Riley [69], and before Worral and Coley [75] in whose papers the 3×2 wüstites W_i and W'_i are strictly identified experimentally out of Vallet's group. Let us reference also to N. Takayama, E. Kimizuka [72].

Ref. [25] published in **2015** is added to the present version of the paper by Gavarri and Carel published in 2018. A. Stokłosa publishes a wide review with the current presentation (See PART II - DIAGRAMS OF THE CONCENTRATIONS OF POINT DEFECTS FOR THE OXIDES OF 3d TRANSITION METALS, Part 12. Iron oxide $\text{Fe}_{1-\delta}\text{O}$, p. 313-76, 168 References for period 1933-2008). It concerns specifically **the defect composition of wüstite** and the **concentration diagrams of point defects and clusters** combined with the mass action law and the electronic structure. That is an advanced systematics in the characterization of the “**three w_i** ”, in accordance with Vallet' - Carel - Gavarri's propositions. The complementary analyze of results concerning wüstite-manganosite alloys* led to improve the characterization of the pure pseudo phases and their transitions.

*As examples

- W_i doped with Mn [85] 1987: A. Stokłosa, H. M. Soliman, Point defect Structure of Mn-Doped Ferrous Oxide, Solid State Ionics 42 (1990) 75-84.

- A. Stokłosa, J. Zajecki, T. S. Wiltowski, A Correlation between the change of the chemical potential of ionic and electronic defects and the phase stability range as exemplified by solid solutions of manganese in iron oxide $(\text{Fe}_{1-x}\text{Mn}_x)_{1-y}\text{O}$, Solid State Ionic 116 (1999) 167-178.

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Presently is a second version with many specific corrections and few additional references.

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