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Light-Induced Spin State Switching in Copper(II)-Nitroxide Based Molecular Magnet at Room Temperature

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Abstract

Molecular magnets $\text{Cu}(\text{hfac})_2\text{L}^{\text{R}}$ exhibit an unusual type of photoinduced magnetostructural switching in exchange-coupled copper(II)-nitroxide clusters. Such photoswitching from strongly-coupled to weakly-coupled spin state (SS \rightarrow WS) was recently found to be ultrafast, thus enhancing the interest in these systems and the scope of their potential applications. However, up to date such SS \rightarrow WS photoswitching was demonstrated only at cryogenic temperatures, being limited by the absence of suitable SS states and short relaxation times at $T > 100$ K. In this work we selected model compound $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ residing in the mixed SS/WS state at room temperature and investigated it using femtosecond optical spectroscopy. Photoinduced spin dynamics was detected, and an ultrafast SS \rightarrow WS photoswitching was for the first time demonstrated at room temperature, constituting an important milestone in the development of copper(II)-nitroxide molecular magnets for practical purposes.

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2 Design of switchable molecular materials is a dynamically developing field within materials
3 science, where magneto-active compounds draw a particular attention due to the possibilities of
4 using them in data processing/storage and other molecular spin devices.¹⁻⁶ The latter type of
5 compounds is largely represented by spin-crossover (SCO) complexes, which were actively
6 investigated for a few past decades.⁵⁻²⁷ The change of the spin state could be driven by a number of
7 external stimuli including temperature, light, pressure, hard X-rays etc., among which light appears
8 to be most practically useful and therefore impressive. The discovery of photoswitching and Light-
9 Induced Excited Spin State Trapping (LIESST)⁷ phenomenon in SCO compounds stimulated a huge
10 amount of following studies, as well as intensive development of new types of photoswitchable
11 molecules.^{3,8,10-13,28}

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20 Relatively recently the manifestations similar to LIESST were observed for principally
21 different compounds.²⁹⁻³² Copper(II)-nitroxide molecular magnets $\text{Cu}(\text{hfac})_2\text{L}^{\text{R}}$ (where
22 hfac=hexafluoroacetylacetonate and L^{R} is a stable nitroxide radical) exhibit thermally-induced
23 magnetostructural transitions between the so-called weakly-coupled spin state (WS) and strongly-
24 coupled spin state (SS), which differ by geometry of metal coordination and the magnitude/sign of
25 exchange interaction between metal and its ligand (Fig.1). In most cases, the low-temperature state
26 is characterized by nitroxides occupying equatorial coordination positions of copper and strong
27 antiferromagnetic exchange coupling (SS state), whereas the high-temperature state corresponds to
28 the nitroxides in axial coordination positions (elongated Jahn-Teller axis of octahedron is flipped)
29 and weak ferromagnetic exchange coupling (WS state). Unit cell volume in these compounds
30 reversibly changes with temperature; therefore they are often called “breathing crystals”. Irradiation
31 with light allows switching of breathing crystals from SS to WS state, and such photoinduced WS
32 state remains metastable at cryogenic temperatures on the time scale of hours.³² Various aspects of
33 SS→WS photoswitching and reverse WS→SS relaxation were investigated using Electron
34 Paramagnetic Resonance (EPR) and steady-state Infrared Spectroscopy in a series of works.³³⁻³⁸
35 However, the information on kinetics and mechanism of the photoswitching could only be obtained
36 recently using femtosecond optical spectroscopy.³⁹ That study, first of all, clearly demonstrated that
37 photoswitching in breathing crystals is an ultrafast process occurring within available temporal
38 resolution and takes <50 fs. No signatures of the intermediate states upon intersystem crossing were
39 detected, therefore it was generally proposed that photoswitching in $\text{Cu}(\text{hfac})_2\text{L}^{\text{R}}$ family might be
40 faster compared to SCO compounds because it is partly spin-allowed,³⁹ even though a sub-50 fs
41 photoswitching in SCO compound was demonstrated slightly later.⁴⁰

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59 Along with the shortening of time required for photoswitching, another appreciable
60 improvement would be creating robust photoswitches functioning at room temperature. For SCO

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2 compounds such examples are already known,^{41-42,17} however, for the family of breathing crystals
3 $\text{Cu}(\text{hfac})_2\text{L}^{\text{R}}$ this task has not yet been achieved. One fundamental obstacle for this was the
4 occurrence of thermal $\text{SS} \leftrightarrow \text{WS}$ transitions only at low temperatures, resulting in ground WS state
5 for most of the compounds at room temperature, which excludes possibility of $\text{SS} \rightarrow \text{WS}$
6 photoswitching, while the reverse transition $\text{WS} \rightarrow \text{SS}$ was never reported in the breathing crystals.
7 However, recently we demonstrated that $\text{SS} \rightarrow \text{WS}$ photoswitching can also be induced for the
8 “thermally-unswitchable” breathing crystals, i.e. compounds exhibiting no or partial thermal
9 switching to WS state in the whole range of their thermal stability.³⁶ One of such compounds,
10 $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$, was illuminated at cryogenic (<20 K) temperatures, resulting in observation of EPR
11 spectra characteristic of photoinduced WS state. At $T > 20$ K the lifetime of this photoinduced state
12 was shorter than temporal resolution of EPR (a few minutes in steady-state mode), therefore the
13 photoswitching was not evidenced. However, it is fair to expect that for this and similar “thermally-
14 unswitchable” compounds the photoswitching occurs even at room temperatures. Therefore, in the
15 present work we have applied the femtosecond optical spectroscopy to this breathing crystal
16 $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ at room temperature, which allowed demonstrating for the first time the feasibility
17 of ultrafast photoswitching in these compounds at ambient conditions.

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31 The target compound $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ was synthesized according to the previously developed
32 procedures.⁴³ To embed microcrystals in poly(vinyl chloride) (PVC) films, 1–2 mg of each complex
33 was dissolved in 1 mL of solution prepared from 1 g of a low molecular weight PVC and 200 mL of
34 dichloroethane. This solution was equally sprayed onto a thin quartz plate to give a brown colored
35 transparent film. Upon drying in ambient atmosphere for 24 h, the film changed its color to blue
36 that indicates a formation of microcrystal particles (<30 μm) of target compound within PVC
37 films.^{35,39} Prior to the optical measurements, the identity of the compound embedded in PVC film
38 was verified by EPR. The details of the time-resolved optical experiments are provided in
39 Supporting Information (SI) and elsewhere.⁴⁴

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47 Figure 1 sketches the structure of polymer chains of $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ (Fig. 1a) and shows the
48 temperature dependence of its effective magnetic moment $\mu_{\text{eff}}(T)$ (Fig. 1b). The complete conversion
49 of $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ to WS state cannot be accomplished by temperature; however, as mentioned
50 above, this transition can be successfully driven by light at low temperatures.³⁶ UV-vis absorption
51 spectra are typical of breathing crystals containing spin triads nitroxide-copper(II)-nitroxide and are
52 dominated by intense bands of nitroxides (Fig. 1c). In addition, a weaker band appears at low
53 temperatures at ~ 500 nm (shaded green), assigned to the metal-ligand charge transfer (MLCT)
54 between nitroxide and copper, characteristic of SS state.^{34,35,39} In the case of $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ this
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band is still distinguishable even at room temperature, because $\sim 50\%$ of clusters still reside in the SS state (to be compared with pure WS state spectrum in Ref.39).

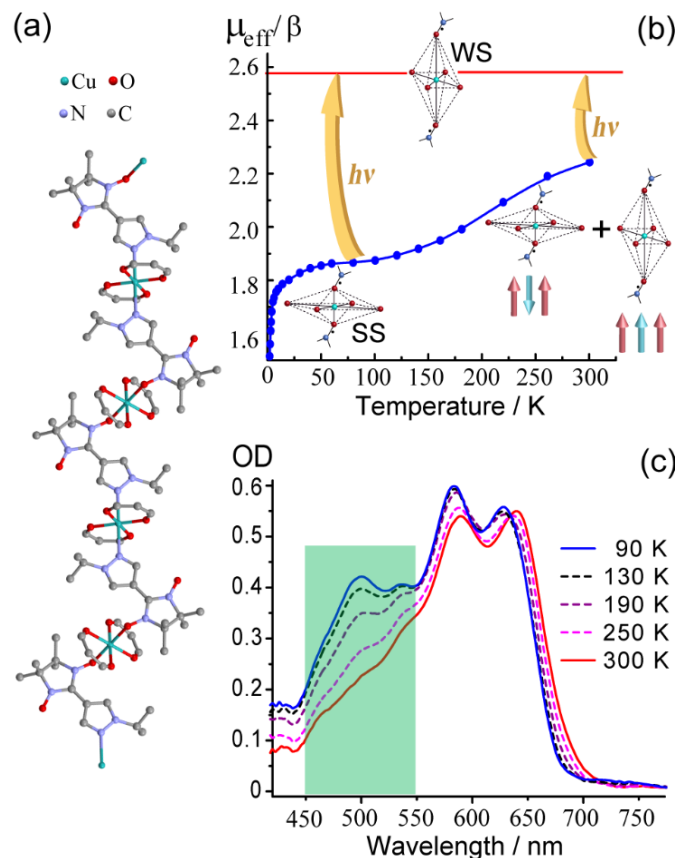


Figure 1. (a) Structure of polymer-chain complex $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$. (b) Temperature dependence of the effective magnetic moment (μ_{eff}) of $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$. Structures of SS and WS states are sketched, and corresponding arrangements of spins are shown using arrows. Yellow arrows illustrate the photoswitching to WS state, and horizontal red line – the magnetization level characteristic of WS state. (c) Temperature dependence of UV(vis) absorption spectrum of $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$. Spectral region of MLCT band is shaded with green.

In order to investigate the kinetics of photoswitching, we performed pump-probe time-resolved experiment by pumping the system at 675 nm and probing with white light (see SI for details). Remarkably, the signal at room temperature was readily detected (Fig. 2a), possibly evidencing the presence of photoinduced changes in the studied breathing crystal $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$. In general, the negative sign of the observed signal is consistent with previous observations of photoswitching at low temperatures (Ref. 39, and *vide infra*). However, the temporal evolution of the observed kinetics is more complicated compared to that obtained previously at low temperature for another breathing crystal.³⁹ The kinetic trace, representative of many measurements series and shown in Fig.2a, is contributed by several processes developing at different time scales and yielding effects of different signs (see more kinetic traces and detailed description of fitting procedures in SI). Abrupt jump-down at sub-picosecond time (at $t \sim 0$) is followed by a slower decrease of the optical density (OD) with characteristic time of a few picoseconds; then the trend is reversed and

OD increases on the time scale of hundred picoseconds. The observed kinetics can be satisfactorily simulated (Fig.2a, red) with two predominant contributions #1 and #2 (Fig.2b), which we are going to assign below.

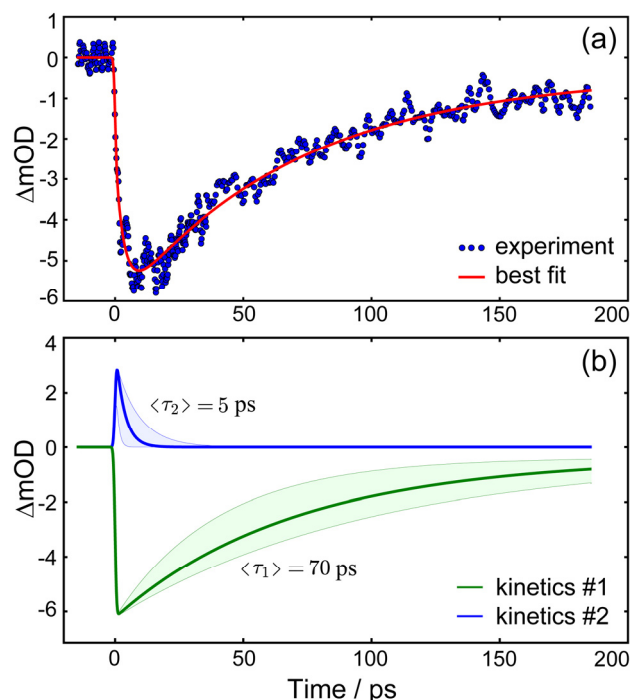


Figure 2. (a) Experimental kinetics $\Delta mOD(t)$ observed in pump-probe experiment for $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ at 300 K (black); $mOD=0.001 \cdot OD$. Pumping was performed at 675 nm, and probing with the white light integrated between 450-550 nm (see Fig.1c and SI for details). Red curve shows the simulation using the function $f(t) = \text{IRF} \cdot h(t)$, where $h(t) = A_1 \cdot \exp(-t/\tau_1) + A_2 \cdot \exp(-t/\tau_2) + C$, $\langle \tau_1 \rangle = 70 \text{ ps}$, $\langle \tau_2 \rangle = 5 \text{ ps}$, $C = -0.00037$; IRF= instrument response function. (b) Deconvoluted contributions #1 ($A_1 \cdot \exp(-t/\tau_1)$) and #2 ($A_2 \cdot \exp(-t/\tau_2)$). The shaded areas around deconvoluted kinetics illustrate the uncertainty due to S/N accumulated over different samples (SI).

The major contribution #1 consists of abrupt decrease of OD after the laser pulse and slow decay of this negative signal with the characteristic time $\langle \tau_1 \rangle = 70 \text{ ps}$ (determined with exponential fit). The smaller contribution #2 ($\approx 50\%$ of #1 by amplitude) exhibits an abrupt jump-up at sub-picosecond time followed by exponential relaxation with the characteristic time $\langle \tau_2 \rangle = 5 \text{ ps}$. In addition, there is a small negative plateau reached by the kinetics at $t > 180 \text{ ps}$ (see SI for details on fitting protocols). As an alternative to the model of two contributions (#1 and #2), the kinetic trace can also be reliably reproduced with two consecutive first-order reactions (ground state \rightarrow intermediate state \rightarrow final state), but the time-constant of 5 ps would imply an unusually long-lived intermediate in the photoexcitation scheme of breathing crystals.³⁹ Spectral wise, the supposed intermediate would not have exactly the same absorbance as the ground- or the final- state, at least to the extent where its spectral intensity if not spectral shape would differ. Neither of such changes was observed on the picosecond timescale for a sibling compound.³⁹ Furthermore, the relative

weights of kinetics #1 and kinetics #2 varied between different parts of the film sample (see SI). The above arguments strongly support two independent processes assignable to the two species (vide infra and SI), rather than two states of one species, in the kinetic description of our results. In order to assign and interpret the contributions #1 and #2, we have performed a series of auxiliary benchmark measurements.

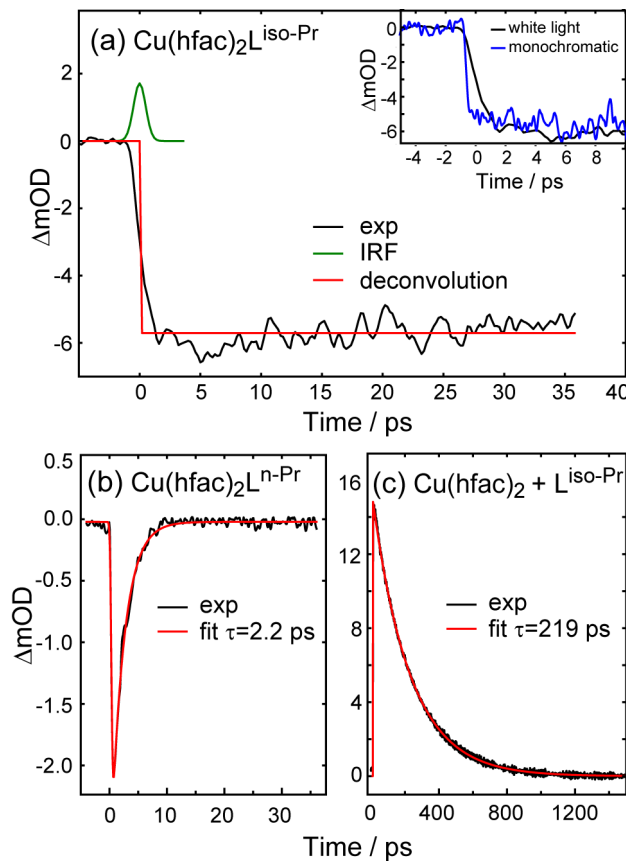


Figure 3. (a) Auxiliary pump-probe experiment on $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ at 100 K with monochromatic pump at 675 nm and white light probe; the experimental kinetics (black) is deconvoluted into a step function (red) using pump probe Gaussian cross-correlation function (green, IRF= 1.2 ps). The inset shows the same kinetics (black) compared with a monochromatic probe experiment (blue, IRF= 200 fs), pump unchanged; see SI for details. (b) Experimental kinetics of the sibling solid state compound $\text{Cu}(\text{hfac})_2\text{L}^{\text{n-Pr}}$ at ambient temperature. The fitted mono-exponential decay (red) has the time constant of 2.2 ps. (c) Experimental kinetics of the compound $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ dissolved in acetone and measured at ambient temperature. The fitted mono-exponential decay (red) has the time constant of 219 ps. In (b, c) the pump and probe wavelengths are 675 nm and 500 nm, respectively.

Figure 3a shows the experimental kinetics $\Delta\text{OD}(t)$ of $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ measured at low temperature 100 K, where the compound resides in pure ground SS state without any admixtures of WS state (see Fig.1b). The shape of the kinetics transforms into nearly pure stepwise jump-down, devoid of any additional processes (including fast relaxation), and closely resembles the previously observed photoswitching kinetics of the sibling compound $\text{Cu}(\text{hfac})_2\text{L}^{\text{n-Pr}}$ at 100 K.³⁹ Due to such remarkable similarity, this contribution, most clearly, corresponds to the SS \rightarrow WS switching, which

is (i) ultrafast (occurs within the temporal resolution of our setup) and (ii) has negative sign of $\Delta OD(t)$ due to the smaller OD in WS state compared to SS state (see Fig.1c). Since there is no reason to expect a different excitation pathways between SS and WS states at room temperature and at 100 K, we conclude that an abrupt decrease of the experimentally observed kinetics in Fig.2a is governed by the target SS \rightarrow WS photoswitching, evidenced at room temperature for the first time; its deconvoluted contribution is shown in Fig. 2b (kinetics #1). The time scale of the SS \rightarrow WS switching is shorter than the temporal resolution of the experiment, see the inset in Fig. 3a. Although we will not speculate on the nature of this initial process, being very likely the charge transfer from metal center to the nitroxide moieties with possible inter-system crossing, we can unambiguously conclude on its sub-200 fs timescale.

In addition to the SS \rightarrow WS photoswitching in $\sim 1/2$ of spin triads having the ground SS state at room temperature, one might anticipate some signals arising from photoexcitation of the other $\sim 1/2$ of triads having the ground WS state (Figure 4). These triads might participate in processes of the type WS \rightarrow WS* \rightarrow WS, where WS* is the excited state in WS coordination; such processes were previously studied using time-resolved EPR.³⁵ To learn how these processes manifest themselves at room-temperature, we investigated the sibling compound $\text{Cu}(\text{hfac})_2\text{L}^{\text{n-Pr}}$ which fully resides in WS state at $T > 250$ K (see SI for details). The observed pure WS \rightarrow WS* \rightarrow WS kinetics for this compound is shown in Figure 3b. This kinetic has a negative going ΔOD signal, and is characterized by an instantaneous buildup and a slower decay with time constant $\tau_3 \sim 2.2$ ps. Since the compounds $\text{Cu}(\text{hfac})_2\text{L}^{\text{n-Pr}}$ and $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ show very similar WS \rightarrow WS* \rightarrow WS kinetics at low temperature,³⁵ one would not anticipate large differences in corresponding manifestations at room temperature. Therefore, we expect that characteristic decay time for this WS* \rightarrow WS kinetics in $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ is also in the few ps range at 300 K. Thus, WS \rightarrow WS* \rightarrow WS process should contribute at short times (< 10 ps) to the observed kinetics in Fig. 2a. This contribution is not visible for the kinetics shown in Fig.2a, possibly due to the presence of competing positive-going kinetics (#2 in Fig. 2b) of different origin, not related to SS \rightarrow WS transition (see below and SI for details).

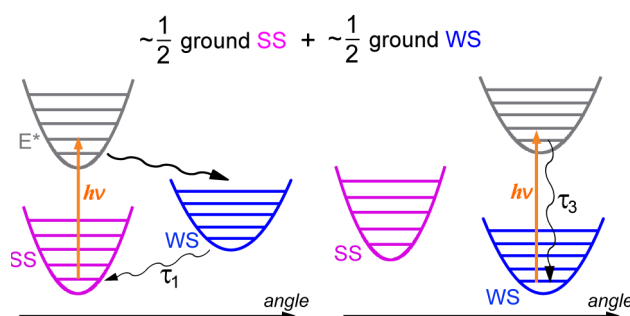


Figure 4. Energy schemes of $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ at 300 K residing in $\sim 50\%$ of ground SS states (left) and $\sim 50\%$ ground WS states (right). Corresponding photoexcitation, switching and relaxation processes are shown.

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3 Since the WS→WS*→WS contribution is expected at $t < 10$ ps, we can only assign the
4 observed slow (~ 70 ps) decay in $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ to a structural relaxation from photoinduced WS to
5 the ground SS states. The timescale of this relaxation curve corresponds well to the expectations.
6 Such WS→SS relaxation occurs on the scale of hours at $T < 20$ K, but speeds up to ~ 100 μs (~ 7
7 orders of magnitude) at $T = 90$ K,³⁹ therefore further acceleration by ~ 6 orders of magnitude at
8 $T = 300$ K up to ≈ 70 ps seems plausible. We have previously shown that structural WS→SS
9 relaxation at low temperatures is self-decelerating in both $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ and $\text{Cu}(\text{hfac})_2\text{L}^{\text{n-Pr}}$
10 compounds, being governed by a broad Gaussian distribution ($\sigma \sim 20\text{-}40$ cm^{-1}) of the energy barrier
11 between WS and SS state potential wells.^{33,36} The relaxation curves observed at low temperatures
12 begin with the rapid decay and then gradually slow down nearly to the plateaus. However, since the
13 dispersion of relaxation rates is governed by the Arrhenius law $\Delta k \sim \exp(-\sigma/kT)$, one would not
14 expect large effect for $\sigma \ll kT$, which is the case at 300 K. Therefore, the anticipated shape of the
15 relaxation kinetics at room temperature should not noticeably deviate from the mono-exponential.
16 Thus, the small plateau distinguishable at long time delays most likely owes to the local heating
17 effects. Indeed, upon photoexcitation the compound should reach thermal homogeneity on the sub-
18 nanosecond timescale,⁴⁵ while the SS/WS fractions settle to equilibrium at an increased temperature
19 on a time scale determined by the energy barrier between SS and WS states and this temperature.⁴⁶
20 At room temperature, such equilibration process will proceed significantly faster, on the
21 nanosecond time-scale, as was evidenced by a temperature jump study.⁴⁷ We can therefore argue
22 that the weak plateau following WS to SS relaxation is the signature of thermal population of WS
23 state. While further investigations at markedly different temperatures and on much longer time-
24 scales would be required to unequivocally pin down the origin of the plateau, they are well beyond
25 the foci of the current paper.

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44 Finally, the contribution #2 (Fig.2b) with positive ΔOD needs to be explained. Since both
45 SS→WS→SS and WS→WS*→WS have negative ΔOD , the small positive going contribution is
46 likely to arise from admixtures of uncomplexed species $\text{Cu}(\text{hfac})_2$ and $\text{L}^{\text{iso-Pr}}$, which are present in
47 the polymer film along with microcrystals in concentration $< 10\%$.³⁸ In order to determine the sign
48 of the corresponding kinetics, we have dissolved $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ in acetone, where it exists as a
49 mixture of $\text{Cu}(\text{hfac})_2$ and free radical ligand $\text{L}^{\text{iso-Pr}}$, and measured as such at room temperature
50 (Figure 3c and SI). Indeed, the obtained kinetics has a positive ΔOD ; however, the decay time of
51 this kinetics is ~ 2 orders of magnitude longer than positive contribution observed in the polymer
52 film. We tentatively attribute this difference to a greater number of channels for energy release,
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1 typically lattice phonon modes, available at solid state. At low temperatures, this contribution
2 becomes negligible by amplitude due to a largely predominant contribution of SS→WS→SS
3 process and, possibly, drastic prolongation of the relaxation times leading to saturation of these
4 signals. Thus, the small positive contribution #2 most likely arises from admixtures of uncomplexed
5 species, whose contribution does not enter into the kinetic scheme of SS→WS transition.
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11 In summary, the observed photoinduced kinetics of $\text{Cu}(\text{hfac})_2\text{L}^{\text{iso-Pr}}$ at room temperature is
12 dominated by SS→WS photoswitching in $\sim 1/2$ of spin triads occurring within the first ~ 200 fs after
13 the laser pulse, and following WS→SS relaxation with the time constant ~ 70 ps (#1). The effect of
14 photoexcitation of the other $\sim 1/2$ of spin triads in the ground WS state, i.e. the WS→WS*→WS
15 process, occurs within first ~ 10 ps; it has a smaller amplitude and temporally coincides with
16 contribution from impurity (uncomplexed species that gives opposite sign of ΔOD) in the film (#2).
17 The contribution of WS→WS*→WS process might be totally eliminated in the future by designing
18 breathing crystals whose thermal transition temperature is shifted to higher than room temperature.
19 Currently, a significant progress was achieved in this direction,⁴⁸ yet with still partial population of
20 unwanted WS states at room temperature. In order to exclude totally the presence of impurities
21 arising from uncomplexed species in the film, one needs to optimize synthetic protocols and obtain
22 more pure materials; before present studies, it was not realized that the presence of such impurities
23 contributes to the fs kinetics at room temperatures. We admit that the exhaustive interpretation of
24 contributions from uncomplexed species and WS→WS*→WS processes is not accomplished in the
25 present work; however, from practical point of view, these are unwanted contributions and should
26 be eliminated synthetically in the future.
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40 Most importantly, this study for the first time evidenced the feasibility of photoswitching in
41 nitroxide-copper(II)-nitroxide molecular magnets at room temperature. The switching occurs within
42 less than 200 fs, whereas the lifetime of the excited (WS) state is at least three orders of magnitude
43 longer (on the order of 100 ps), thus allowing numerous spin switching operations to be
44 accomplished. There is no real trapping of this excited spin state in contrast to low temperatures;
45 however, principal feasibility of such photoinduced magnetostructural transitions is highly
46 important for designing photomagnetic switches. As next steps of this research, we foresee a design
47 of breathing crystals whose ground state at room temperature is a pure SS state, and a design of
48 compounds exhibiting thermal hysteresis around room temperature. Both directions require
49 thorough synthetic efforts, which are now justified by the first observation of photoswitching in this
50 family at room temperatures.
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Supporting Information

Structure and magnetic properties of sibling $\text{Cu}(\text{hfac})_2\text{L}^{\text{n-Pr}}$ compound. Details of femtosecond optical spectroscopy experiments. This material is available free of charge via the Internet at <http://pubs.acs.org>.

Acknowledgements

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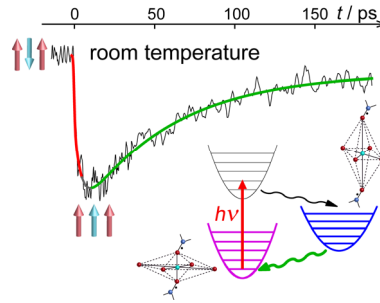
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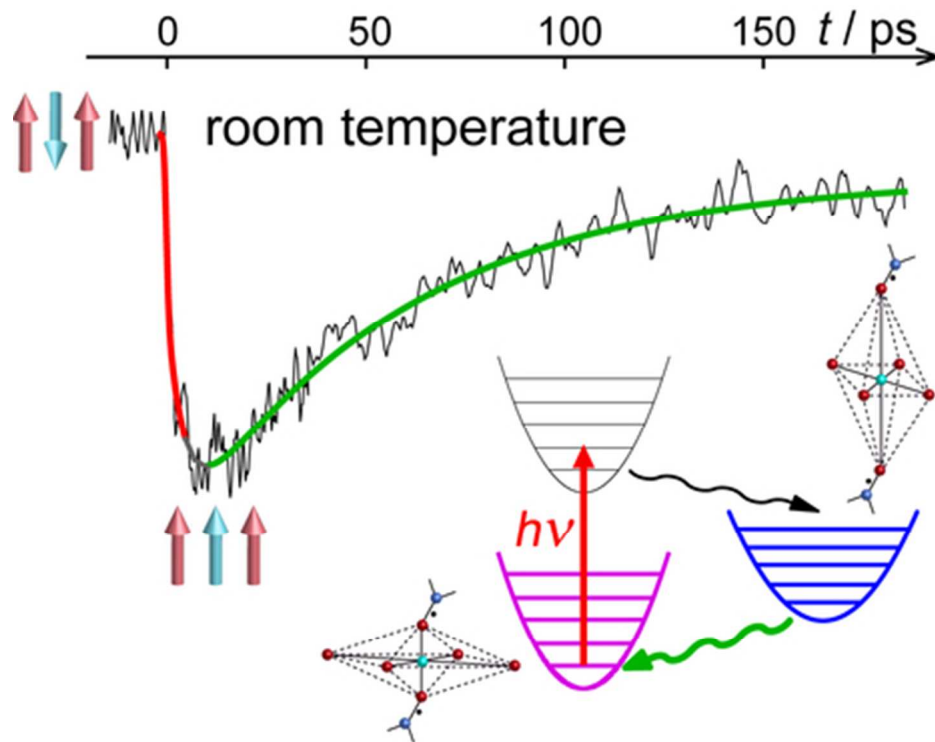
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Table of Contents Graphic





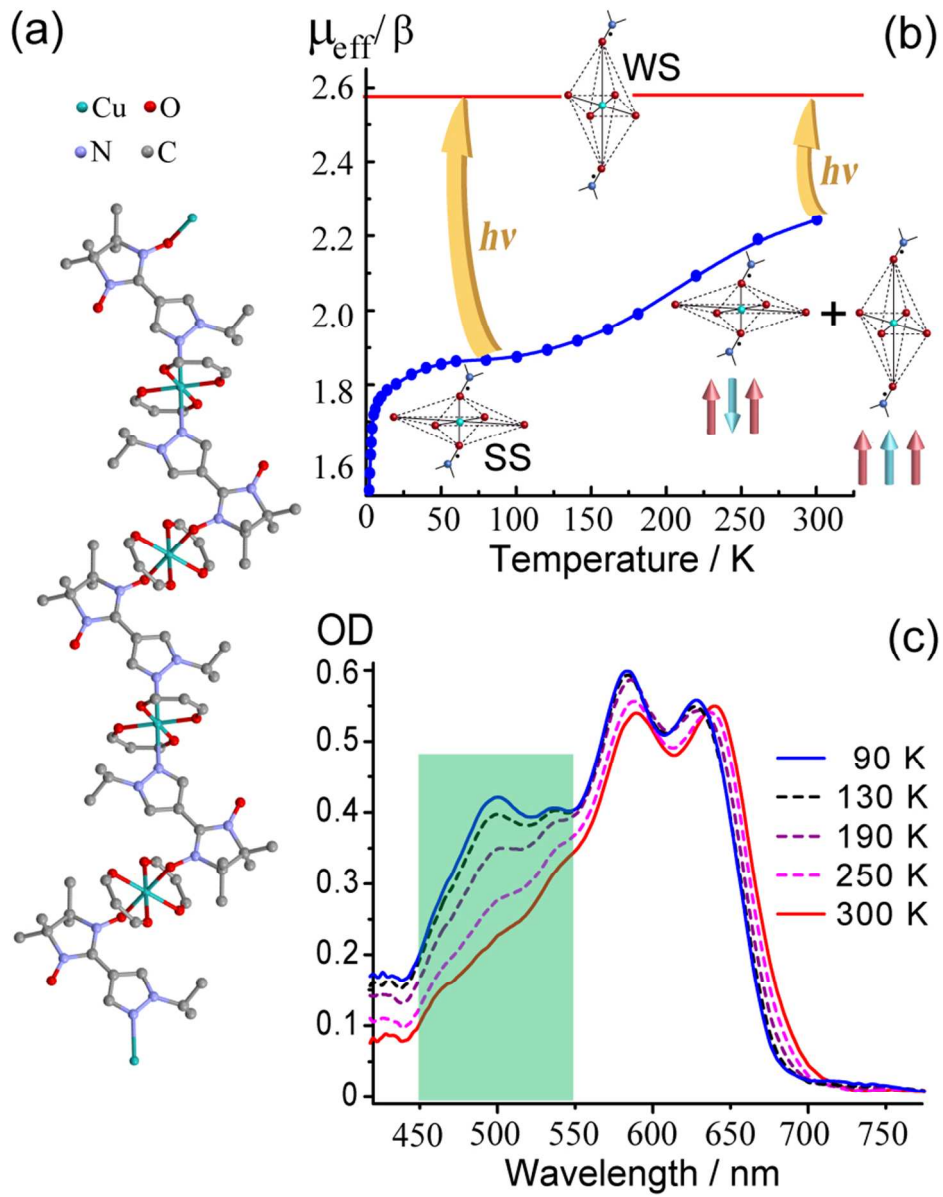


Figure 1

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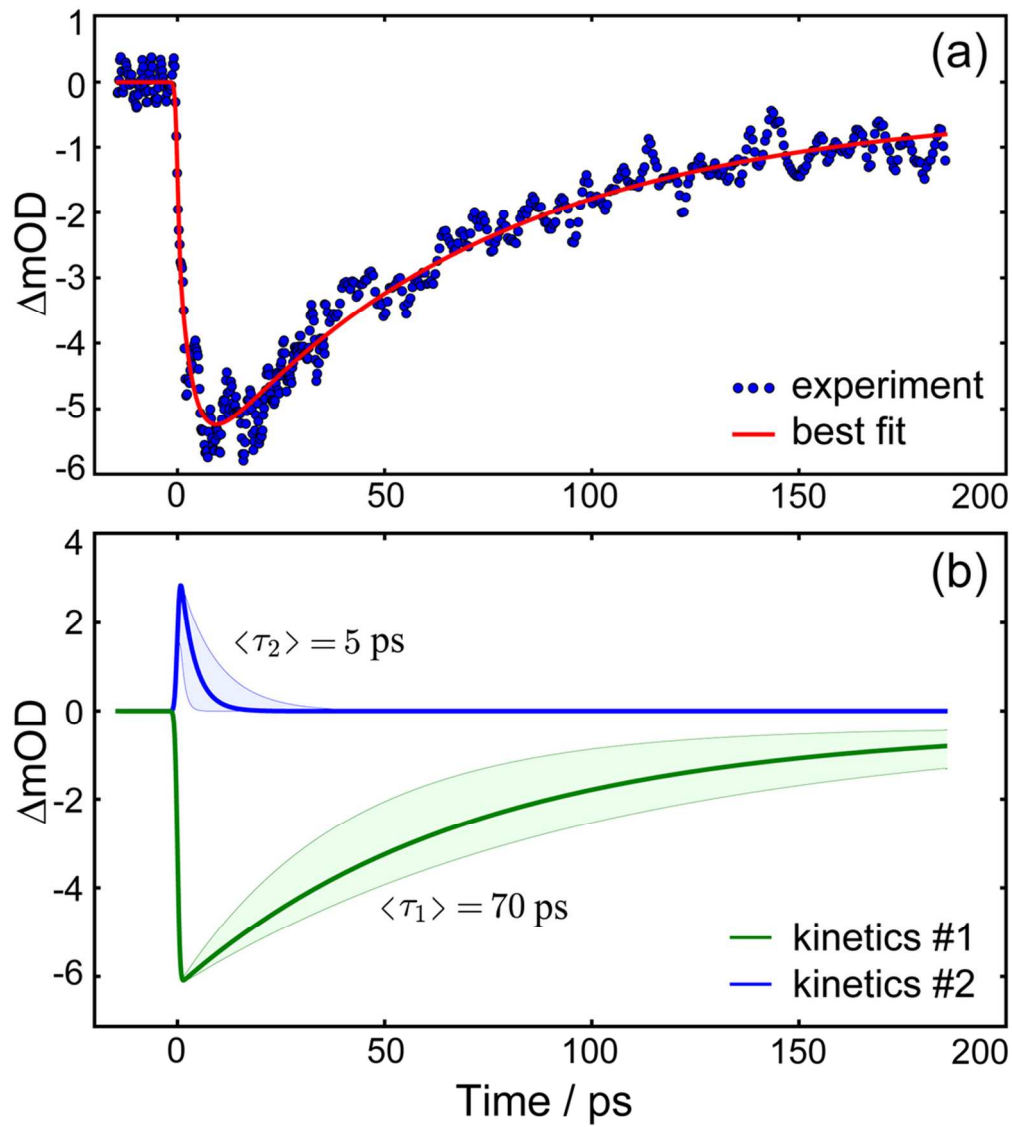


Figure 2

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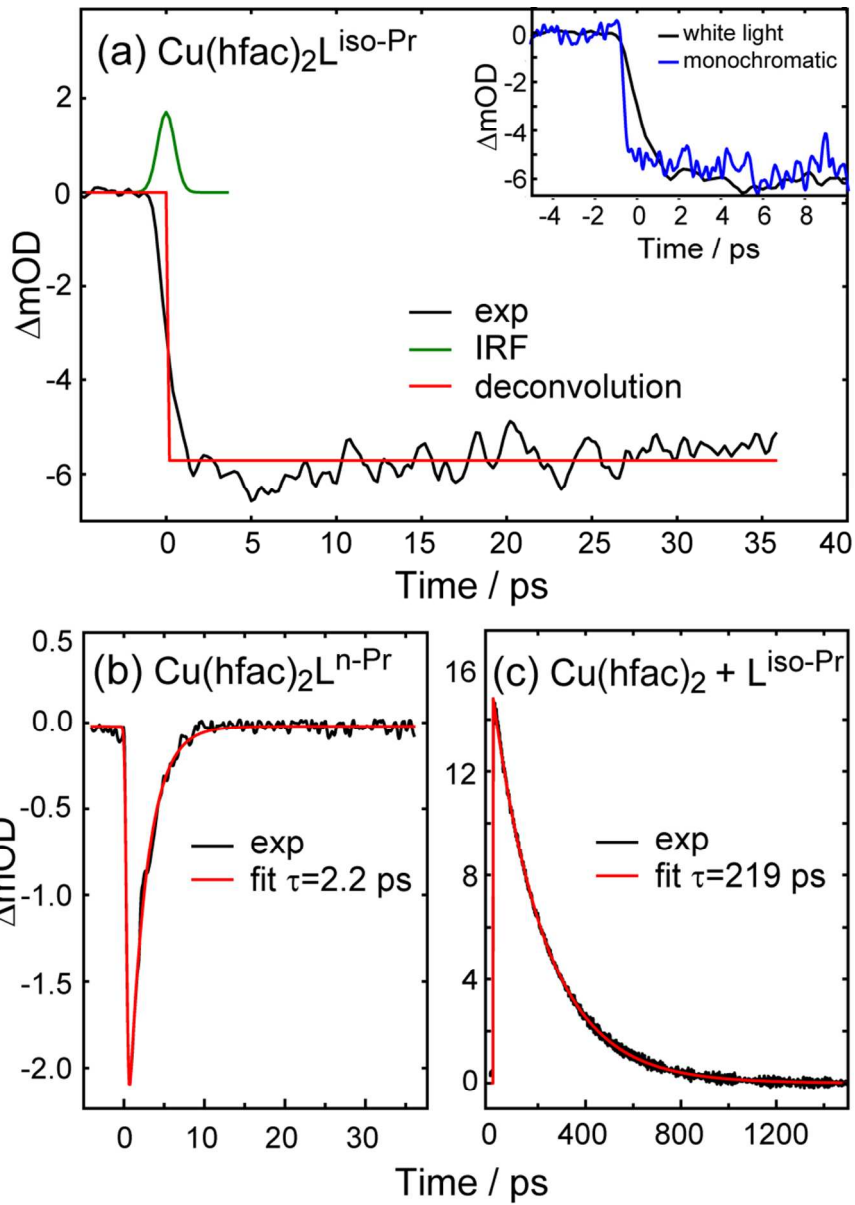


Figure 3

82x113mm (300 x 300 DPI)

