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Excited electronic structure of methylcyanoacetylene probed by VUV Fourier-transform absorption spectroscopy

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Abstract

High resolution photoabsorption spectrum of gas-phase methylcyanoacetylene ($\text{CH}_3\text{C}_3\text{N}$) has been recorded from 44 500 to 130 000 cm^{-1} at room temperature with a vacuum ultraviolet Fourier-transform spectrometer on the DESIRS synchrotron beamline (SOLEIL). The absolute photoabsorption cross section in this range is reported for the first time. Valence shell transitions and Rydberg series converging to the ground state $\tilde{X}^+ \ ^2\text{E}$ of the cation as well as series converging to electronically excited states ($\tilde{\text{A}}^+ \ ^2\text{A}_1$ and $\tilde{\text{C}}^+$) are observed and assigned. Time-dependent density-functional-theory calculations have been performed to support the assignment of the experimental spectrum in the low energy range. A tentative scaling of the previously measured $\text{CH}_3\text{C}_3\text{N}^+$ ion yield by Lamarre *et al.* (Lamarre *et al.* *J. Mol. Spectrosc.* **315** (2015)) is proposed, based on the comparison of the absorption data above the first ionization potential with the observed autoionization structures.

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1. Introduction

Methylcyanoacetylene (CH₃C₃N) is a molecule of astrophysical interest, detected in molecular clouds of the interstellar medium (TMC-1 [1], Sgr-B2 [2]), in circumstellar envelopes [3], and anticipated in the atmosphere of Titan (with
5 no firm detection up to now) [4, 5, 6, 7, 8]. It is the first member of the CH₃(CC)_nCN homologous series; the $n = 2$ species (methylcyanodiacetylene) has also been detected in Space [9].

Concerning the spectroscopic fingerprint of this species, the rotational structure of the ground electronic state \tilde{X}^1A_1 has been studied by microwave spectroscopy [10, 11, 12, 13]. Infrared spectroscopy studies in the gas phase have
10 been performed [14, 4, 6] to determine the 12 vibrational modes of this state (6 symmetric a₁ modes and 6 degenerate e modes in the C_{3v} symmetry group). The unique absorption spectrum involving electronic transitions reported to date ranges in the UV domain from 43 000 to 54 000 cm⁻¹ and has been measured
15 by Bruston *et al.* in 1989 [5]. It led to observation of 3 vibrational progressions of a valence-shell electronic transition. Bruston *et al.* compared this spectrum with the one of cyanoacetylene HC₃N in order to partially assign the observed combination bands. They derived vibrational frequency values for the ν_2 (C \equiv N and C \equiv C stretches) and ν_5 (C-C stretch) modes of the excited valence state [5].
20 To our knowledge, no other information is available on the excited electronic states of neutral CH₃C₃N, neither by experimental VUV spectroscopy nor by *ab initio* calculation. The only available data concern the cation: in the late seventies the He(I) and He(II) photoelectron spectra in the gas-phase have been measured at medium resolution (~ 160 cm⁻¹) [15, 16], and very recently the
25 threshold photoelectron spectrum of CH₃C₃N has been recorded by our group at better resolution (~ 80 cm⁻¹) [17]. In addition, the electronic absorption spec-

trum of methylcyanoacetylene cation has been obtained in 4.5 K neon matrix by Fulara *et al.* in 1985 [18]. These experiments allowed to determine the ionization potential and the vibronic structure of the 4 lowest electronic states of the $\text{CH}_3\text{C}_3\text{N}^+$ cation, namely ($\tilde{X}^+ \ ^2\text{E}$, $\tilde{A}^+ \ ^2\text{A}_1$, $\tilde{B}^+ \ ^2\text{E}$ and \tilde{C}^+). Although they relate to the cationic species, they are very useful to understand the excited electronic structure of the neutral molecule, in particular its Rydberg states which resemble the cationic state toward which they converge.

In the present work, we report the first VUV photoabsorption spectrum of methylcyanoacetylene in the gas phase, recorded from 44 500 to 130 000 cm^{-1} with a 4 cm^{-1} spectral resolution thanks to an original Fourier-transform spectrometer.

2. Experiment

2.1. Experimental details

The experiment was performed at SOLEIL synchrotron facility on the DESIRS beamline [19], using the full bandwidth ($\Delta\sigma/\sigma = 7\%$) of the undulator source to illuminate a unique VUV Fourier-Transform Spectrometer (FTS) installed in a high-vacuum chamber [20]. The beamline gas filter filled with xenon or argon was employed to remove high harmonics from the synchrotron light beam before it interacts with methylcyanoacetylene. High resolution absorption spectra were recorded at room temperature over consecutive spectral windows corresponding to 21 different undulator settings, and they were spliced together to build the overview spectrum spanning from 44 500 to 130 000 cm^{-1} . The typical resolution of the spectra recorded in this work was 4 cm^{-1} , which appeared to be sufficient as the narrowest features present in the spectrum have approximately a 15 cm^{-1} linewidth. Spectral calibration was achieved by recording absorption lines of impurities present in the gas filter. The sample was introduced through a Granville Phillips valve in a constant flow mode into a stainless-steel cylindrical cell ended at both extremities by open thin-capillaries (15 cm long each). The length of the windowless cell was 10 cm. The pres-

sure in the cell was monitored over the scans by a capacitive gauge located outside the main vacuum chamber. In most experiments the typical pressure inside the windowless cell was in the 0.1–1 mbar range. Calibration of the absolute photoabsorption cross section was performed using a MgF₂ windowed cell and different pressure values to assess the real column density distribution inside the windowless cell. The complete procedure is described in the following paragraph. Synthesis of methylcyanoacetylene was performed prior to the experimental campaign according to the procedure described in Ref. [17] and stored in a freezer. During the measurements, the sample tube was placed in a glass container in which a cooling fluid was circulating *via* a chiller to ensure a constant working temperature of +22°C and hence a constant vapor pressure of 35 mbar above the liquid.

2.2. Cross section calibration procedure

The FTS general method of absolute photoabsorption cross section measurement has already been described [21]. In principle, it requires the recording of a single spectral range using a MgF₂ windowed cell in order to infer the absolute absorption cross section for this particular range. The next step involves the windowless cell and allows recording the full spectrum by stitching the individual 7% bandwidth spectral windows over the desired VUV range. A general calibration curve is determined by recording various absorption spectra over a set of entrance pressures on this particular spectral window. Although in principle the calibration procedure requires only a single spectral window, a first set of scans has been recorded between 44 000 cm⁻¹ and 70 000 cm⁻¹ at low resolution and moderate signal-to-noise ratio in the windowed cell in order to set a proper absolute scale. Indeed, the range of photoabsorption cross sections for methylcyanoacetylene in the VUV spans more than three orders of magnitude, and such dynamics would have made difficult the determination of a correct calibration curve using a single spectral window. It appeared that the sample was particularly reactive and fragile when exposed to the high flux synchrotron beam in the windowed cell. The recording was tricky as the molecule suffered

photodestruction and the signal was dropping over a single low resolution scan of a few seconds inside the cell, due to a fast contamination of the cell windows. Nevertheless, it was possible to determine absolute values because the signal affected by the contamination decreases linearly over the acquisition time.

90 Therefore, a normalization on the beginning of the scan was done and a reduction of the data using the background recorded before the scan on an empty cell was performed. No variation of the set pressure was observed, therefore we assumed a constant column density within our error margin over the short time needed to record a spectrum (≈ 10 seconds). The windowless spectrum was

95 calibrated onto the windowed cell data up to $70\,000\text{ cm}^{-1}$. Above $70\,000\text{ cm}^{-1}$, the spectral ranges were normalized on the entrance pressure values and simply connected together as the cross section does not change significantly over this region. In the $44\,500\text{--}55\,000\text{ cm}^{-1}$ spectral range, corresponding to the low cross section region, the spectrum is exclusively based upon the data recorded

100 in the windowed cell. Indeed, the required pressure would have been too high to be easily supported by the windowless cell differential pumping system. Note that the experimental linewidth in this region is $\delta\sigma = 16\text{ cm}^{-1}$ for the reasons mentioned above. We believe that most of the absolute photoabsorption cross section uncertainty originates from the difficulty to obtain a constant signal

105 over the time needed to achieve a single scan in the windowed cell. The error estimation is based on the reproducibility of the individual spectral windows at various pressures and also on the non-overlap of two neighbored spectra during the stitching procedure of the different spectral windows. The maximum observed amplitude of the dispersion from scan to scan is $\pm 20\%$ below $58\,000\text{ cm}^{-1}$

110 and it decreases down to $\pm 10\%$ in the spectral region above $58\,000\text{ cm}^{-1}$. Note that between $55\,000\text{ cm}^{-1}$ and $57\,000\text{ cm}^{-1}$, due to a lack of time, only a single scan was recorded in the windowless cell with a limited absorption ($\approx 3\%$), for this reason the uncertainty may be higher in this particular region where the absorption cross section variation is quite important.

115 3. Quantum chemical calculation methodology

The present calculations employed the density functional theory (DFT), and were carried out with the Gaussian 09 Rev B.01 software [22]. The CAM-B3LYP functional [23] and aug-cc-pVQZ basis set [24, 25] were used. This Coulomb-attenuated exchange-correlation energy functional offers an enhanced
 120 description of long-range exchange interactions which improves the predictions of electronic excitation energies of low- n Rydberg states (n being the principal quantum number). In the first step, geometric parameters of the ground electronic state have been optimized assuming its C_{3v} symmetry. The electronic configuration of the ground electronic state is [core]($1a_1$)² ($2a_1$)² ($3a_1$)² ($4a_1$)²
 125 ($5a_1$)² ($1e$)⁴ ($2e$)⁴ ($6a_1$)² ($3e$)⁴. Subsequently, vertical excitation energies were obtained using the time-dependent DFT approach [26, 27, 28, 29, 30, 31]. All integrals have been calculated with an *UltraFine* grid option. Standard convergence criteria were applied.

Note that, optimizing the excited electronic state geometries as well as calculating the respective vibrational frequencies, fell beyond the scope of this
 130 work.

4. Results and discussion

4.1. Absolute absorption cross section

The absorption spectrum of methylcyanoacetylene from 44 500 to 130 000 cm^{-1}
 135 is displayed in Fig. 1 on a logarithmic scale. The vertical scale has been calibrated in absolute absorption cross section using the procedure described in section 2. One can distinguish 4 regions in the spectrum:

- up to 56 000 cm^{-1} , the spectrum features weak bands of the valence shell system $\tilde{B} \leftarrow \tilde{X}$ only (see below),
- between 56 000 cm^{-1} and 86 872 cm^{-1} (ionization potential towards the
 140 \tilde{X}^+ state of the cation), apart from strong valence shell transitions, the

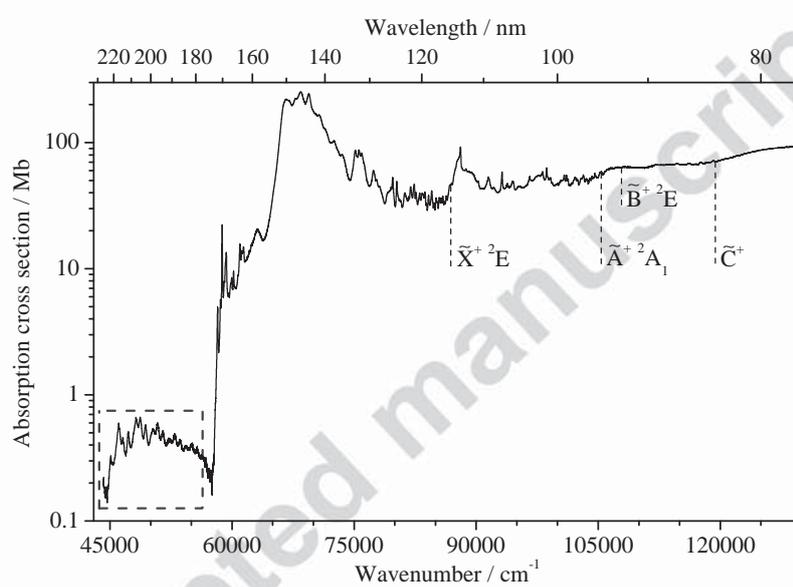


Figure 1: Absolute photoabsorption cross section of methylcyanoacetylene (1 Mb = 10^{-18} cm²). The vertical dashed lines represent the energetic positions of the lowest electronic states of the cation [17]. The dashed rectangle corresponds to the spectral region which has been previously investigated by Bruston *et al.* [5] (see text and Fig. 2).

spectrum is dominated by the Rydberg series converging to the \tilde{X}^+ state of the cation,

- between 86 872 cm^{-1} and approximately 108 000 cm^{-1} , the spectrum consists mostly in Rydberg series converging *a priori* towards the \tilde{A}^+ and/or \tilde{B}^+ states of the cation,
- above 108 000 cm^{-1} , the structures are probably due to Rydberg series converging towards the \tilde{C}^+ state of the cation.

The assignment of the observed features is explained in the two following sections.

The only published absorption measurement prior to the present work was recorded by Bruston *et al.* in the UV region (40 800 – 54 000 cm^{-1}) [5] materialized by a dotted rectangle in Fig. 1, thus allowing a comparison with our curve in the overlapping region displayed in Fig. 2. As shown in Fig. 2, the positions of the vibrational bands in the region 45 000–53 000 cm^{-1} are similar between the two sets of data, although the present study extends the progression beyond 54 000 cm^{-1} which was the monochromator transmission limit of Bruston *et al.* experiment. Between 44 500 and 48 000 cm^{-1} , the cross section of the two sets of data are almost identical, but in the high energy range the two traces deviate with a difference beyond our claimed uncertainty (20%). A possible contamination in Bruston *et al.* experiment (mentioned by the authors) could explain this extra absorption that would have been caused by the O_2 Schumann-Runge structured absorption bands due to the residual oxygen within their monochromator.

4.2. Valence state vibrational assignment

Fig. 3 displays a comparison of the measured absorption spectrum with theoretically (CAM-B3LYP) predicted vertical electronic excitations departing from the $\tilde{X}^1\text{A}_1$ ground electronic state. Noteworthy, the spectrum qualitatively resembles that of HC_3N [32, 33]. As can be seen, computations satisfactorily reproduce the observed valence shell transitions (displayed as purple sticks in

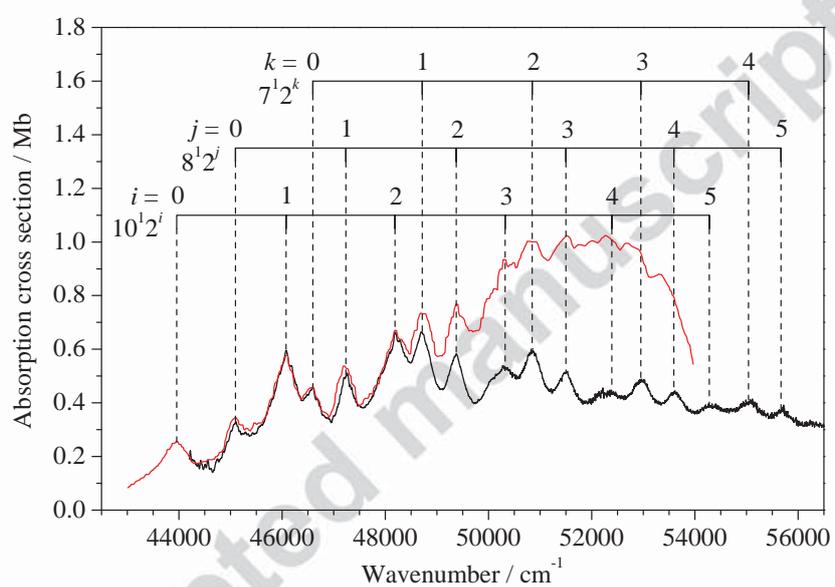


Figure 2: Comparison of the absolute photoabsorption cross section of methylcyanoacetylene between 44 500 and 56 000 cm^{-1} measured in this work (black line) and by Bruston *et al.* [5] (red line). See text for assignment description corresponding to vibrational progressions of the $\tilde{B}^1E \leftarrow \tilde{X}^1A_1$ valence shell transition.

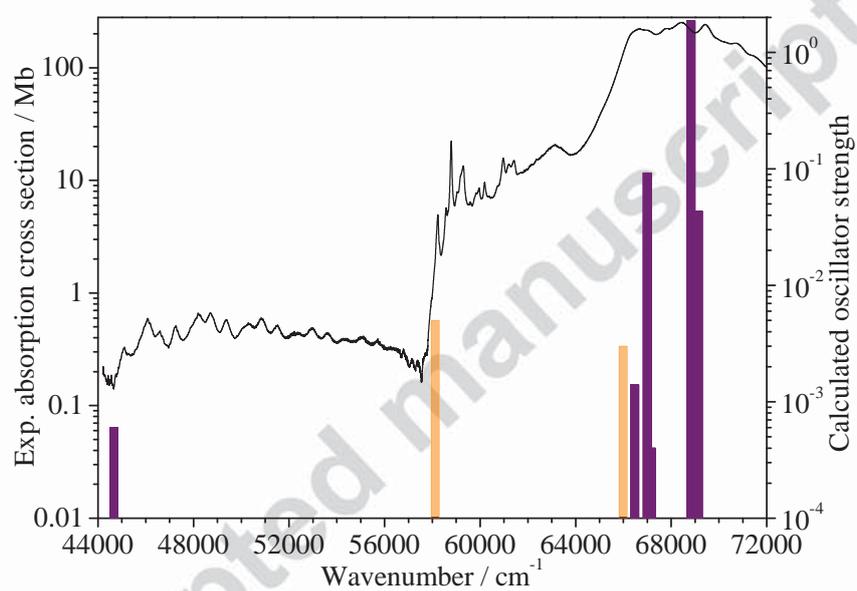


Figure 3: Comparison of the measured methylocyanoacetylene absorption spectrum with theoretically (CAM-B3LYP) predicted vertical electronic excitation energies from the \tilde{X}^1A_1 ground electronic state towards Rydberg states (orange sticks) and valence states (purple sticks).

Fig. 3) both in energetic positions and in relative intensities. This applies in particular to the strongest band, around $68\,800\text{ cm}^{-1}$, recognized as due to $3\ ^1A_1 \leftarrow \tilde{X}\ ^1A_1$ transition (see Table 1) and to the lowest energy $\tilde{B}\ ^1E \leftarrow \tilde{X}\ ^1A_1$ transition around $44\,700\text{ cm}^{-1}$.

Table 1: CAM-B3LYP-derived vertical electronic excitations originating from the $\tilde{X}\ ^1A_1$ ground electronic state of methylcyanoacetylene. Note that the labelling of the states refers to the state order at the neutral ground state geometry. The $3e$, $1e^*$ and $6a_1$ orbitals correspond to the highest occupied molecular orbital, the lowest unoccupied molecular orbital, and the orbital of the lone electron pair of the nitrogen atom, respectively.

Upper state	Orbital excitation	Vertical transition energies [cm^{-1}]	Predicted oscillator strengths
$1\ ^1A_2$	$3e \rightarrow 1e^*$	43.0×10^3	0
$1^1E (\tilde{B})$	$3e \rightarrow 1e^*$	44.7×10^3	6×10^{-4}
$2\ ^1E$	$3e \rightarrow 3s$	58.1×10^3	5×10^{-3}
$2\ ^1A_2$	$2e \rightarrow 1e^*$	64.0×10^3	0
$3\ ^1E$	$3e \rightarrow 3p$	66.0×10^3	3×10^{-3}
$4\ ^1E$	$2e \rightarrow 1e^*$	66.5×10^3	1×10^{-3}
$2\ ^1A_1$	$3e \rightarrow 2e^*$	67.0×10^3	9×10^{-2}
$5\ ^1E$	$3e \rightarrow 2e^*$	67.2×10^3	4×10^{-4}
$3\ ^1A_2$	$3e \rightarrow 2e^*$	67.3×10^3	0
$3\ ^1A_1$	$3e \rightarrow 1e^*$	68.8×10^3	1.9
$6\ ^1E$	$6a_1 \rightarrow 1e^*$	69.1×10^3	4×10^{-2}

175 The latter one, even though formally symmetry-allowed, is in fact closely similar to the forbidden $B\ ^1\Delta \leftarrow X\ ^1\Sigma^+$ system of HC_3N , since the considered chromophores of both molecules consist of essentially the same C_3N chains. This explains its weak intensity.

180 For the vibronic structure of this first observable electronic absorption system of $\text{CH}_3\text{C}_3\text{N}$ ($\tilde{B} \leftarrow \tilde{X}$), Bruston *et al.* proposed a partial assignment involving the ν_2 ($\text{C}\equiv\text{N}$, $\text{C}\equiv\text{C}$ stretches in phase) vibrational progression coupled

with either one quantum of ν_{11} (CCCCN skeleton bend) or a combination of one quantum of ν_{11} and one quantum of ν_5 (C–C stretches, opposite phases). Indeed, isotopic shifts measured in the UV spectrum of cyanoacetylene are consistent with those expected for ν_2 [34]. The combination of the ν_2 and ν_{11} vibrational modes was chosen by the authors by analogy with cyanoacetylene [35]. The combination of ν_{11} with ν_5 , proposed by Bruston *et al.*, was suggested by the values of vibrational frequencies found for the ground electronic state of $\text{CH}_3\text{C}_3\text{N}$. A third progression detected by Bruston *et al.* remained unassigned.

Here we propose a different but complete assignment for the three ν_2 -based progressions observed in the absorption spectrum. First, we suggest the mode ν_{10} (499 cm^{-1} in the ground state [17]) as a better candidate than ν_{11} (339 cm^{-1} in the ground state [17]) for the lowest energy "false origin". Indeed, that latter mode involves the CH_3CC part of the molecule, while the CCCCN skeleton bending ν_{10} is similar to the vibration around 500 cm^{-1} observed for the ground electronic state of HC_3N [35], of the cyanopolyynic chains such as HC_5N [36], C_2N_2 [37], C_4N_2 [38], and is always active in the respective $\text{B} \leftarrow \text{X}$ transitions. Taking into account the measured energy differences between the origins of the three progressions, as well as the fact that the discussed ground-state bending frequencies are reduced by approximately 100 cm^{-1} in the second excited state of the above mentioned molecules, the origins of the two remaining progressions should involve modes of frequencies around 1550 cm^{-1} and around 3050 cm^{-1} . These we propose to identify as, respectively, ν_8 (CH_3 antisymmetric deformation) and ν_7 (CH_3 antisymmetric stretch) [17], both of e symmetry. Analogous chain-stretching mode progressions, gaining strength *via* coupling to e-symmetry modes, have been recently observed for the $\tilde{\text{B}} - \tilde{\text{X}}$ system of $\text{CH}_3\text{C}_5\text{N}$ [39]. This reasoning allows for expecting the vibrationless origin of $\tilde{\text{B}} - \tilde{\text{X}}$ within the $43\,400 - 43\,600\text{ cm}^{-1}$ range, in reasonable agreement with our CAM B3LYP prediction of $44\,700\text{ cm}^{-1}$ (see Table 1).

The relative energies with respect to their origins of the three observed ν_2 progression bands of the $\tilde{\text{B}}^1\text{E} \leftarrow \tilde{\text{X}}^1\text{A}_1$ transition are plotted in Fig. 4 against the ν_2 quantum number. The data points exhibit a linear behaviour which

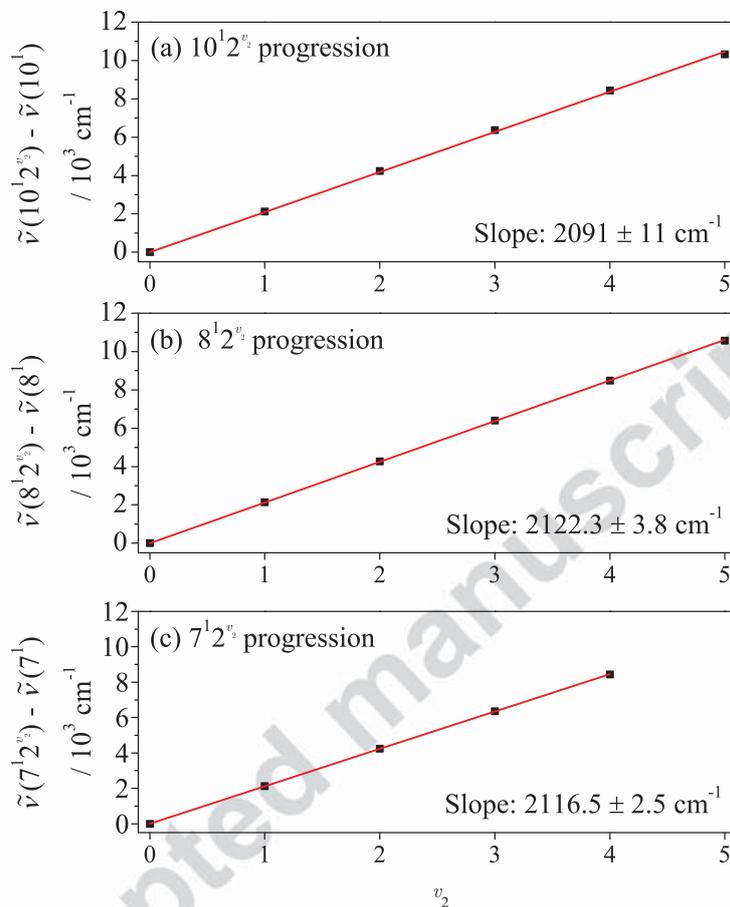


Figure 4: Plots of the 2^v level wavenumbers against the v_2 vibrational quantum number for the three observed v_2 progressions in combination with 10^1 (panel (a)), 8^1 (panel (b)) and 7^1 (panel (c)) with respect to their origin. The black squares are the experimental data and the red lines the linear fits.

implies that the anharmonic parameter of the ν_2 mode (anharmonicity constant $x_{2,2}$) can be neglected in a first approximation. From a linear fit of the three
 215 sets of data, we derived the following slopes: 2091 cm^{-1} , 2122.3 cm^{-1} , and 2116.5 cm^{-1} , for the ν_2 progressions involving, respectively, a quantum of either ν_{10} , ν_8 or ν_7 . Anharmonicity constants $x_{2,8}$ and $x_{2,7}$ are therefore smaller than

$x_{2,10}$, which is quite an intuitive result. Indeed, the vibrational modes localised mainly within the methyl group (ν_8 and ν_7) are not expected to interact strongly with ν_2 , this latter mode being responsible mostly for the CN stretching.

4.3. Rydberg series assignments

Energies of the lowest $\text{CH}_3\text{C}_3\text{N}^+$ electronic states, to which the observed bands of Rydberg series converge, have been recently measured and published by our group [17]. Pertinent information for the present work is reported in Table 2.

Table 2: Ionization potentials of $\text{CH}_3\text{C}_3\text{N}$ towards the four lowest electronic states of the cation, and cationic harmonic vibrational frequencies of the main active modes upon ionization from Ref. [17]. All values are in cm^{-1} and the uncertainties are given in parentheses. Note that the notation of the neutral species (in the C_{3v} point group) is used for the vibrational mode labelling of the cationic states (see Ref. [17] for details).

Cationic state	E_1/hc	Active mode and description	Harmonic frequency
$\tilde{\text{X}}^+ \ ^2\text{E}$	86 872(20)	ν_2 ($\text{C}\equiv\text{N}$, $\text{C}\equiv\text{C}$ stretches in phase)	2178(20)
$\tilde{\text{A}}^+ \ ^2\text{A}_1$	105 354(20)	ν_2 ($\text{C}\equiv\text{N}$, $\text{C}\equiv\text{C}$ stretches in phase)	2246(30)
		ν_6 ($\text{C}-\text{C}$ stretches in phase)	655(30)
$\tilde{\text{B}}^+ \ ^2\text{E}$	107 848(20)	ν_5 ($\text{C}-\text{C}$ stretches opposite phase)	1132(20)
$\tilde{\text{C}}^+$	119 390(50)	ν_8 (CH_3 antisym. deformation)	1586(40)

The difficulty in assigning the absorption spectrum comes from the high density of states and the fact that these states are not fully resolved. Here, we based our assignment on our recent analysis of the TPES spectrum of $\text{CH}_3\text{C}_3\text{N}$ [17]. In this study we determined which vibrational modes are active upon ionization towards each cationic electronic state. These results give us a hint about which vibrational modes can be observed in the transitions towards the Rydberg states. Indeed the well-known property of Rydberg states is to have a similar geometry to the cationic state towards which the Rydberg series converge.

From this consideration, one can assume that the Rydberg series converging to \tilde{X}^+ , \tilde{A}^+ , \tilde{B}^+ and \tilde{C}^+ will exhibit vibrational structures involving mainly the vibrational modes reported in Table 2. Note that we only take into account the vibrational modes which are the most active upon ionization [17].

Finally, to assign the Rydberg series we use plots of the absorption cross section against the effective principal quantum number, n^* ($n^* = n - \delta$, where n is the principal quantum number and δ the quantum defect), following the well-known Rydberg formula ($\tilde{\nu} = E_i/hc - R/(n - \delta)^2$, where E_i/hc is the ionization potential in cm^{-1} and R the Rydberg constant for $\text{CH}_3\text{C}_3\text{N}$ in cm^{-1}). In Appendix A, the plot corresponding to the Rydberg series converging to the ground vibronic state of the cation ($\tilde{X}^+ \ ^2\text{E } v^+ = 0$) is displayed as an example (the other plots are not presented in this paper). The aim of this kind of plot is to recognize the peaks which are present approximately at the same quantum defect value ($\delta = n^* - n$, thus the decimal part of n^*), provided no coupling between valence and Rydberg states exists.

Following this procedure for all the ionization limits ($\tilde{X}^+ \ ^2\text{E } v^+ = 0$, $\tilde{X}^+ \ ^2\text{E } v_2^+ = 1$, $\tilde{A}^+ \ ^2\text{A}_1 \ v^+ = 0, \dots$), we assign most of the observed structures. The corresponding assignment is summarized in Table 3 and the origins of each Rydberg state ($v = 0$) are labeled in Figs. 5, 6 and 7. The quantum defects obtained for the Rydberg series allow us to deduce the s, p or d character of the series by assuming that ns , np and nd series usually have quantum defects of about 1, 0.5 and < 0.2 , respectively.

In Fig. 5, the absorption spectrum of $\text{CH}_3\text{C}_3\text{N}$ between 60 000 and 88 000 cm^{-1} is displayed with 3 assigned Rydberg series converging to the adiabatic ionization potential of the $\tilde{X}^+ \ ^2\text{E}$ electronic ground state (86872 cm^{-1} [17]). The ns , np and nd Rydberg series have averaged quantum defects of 0.94, 0.63, and 0.13, respectively. The Rydberg states are indicated in this spectrum by dashed lines and labeled $ns \ 0^0$ [\tilde{X}^+], $np \ 0^0$ [X^+] and $nd \ 0^0$ [X^+], respectively. The vibrationally excited Rydberg states are not displayed in Fig. 5 for clarity (see Table 3 for detailed assignment and individual quantum defect values and figures provided as supplementary material). Note that the quantum defect of the

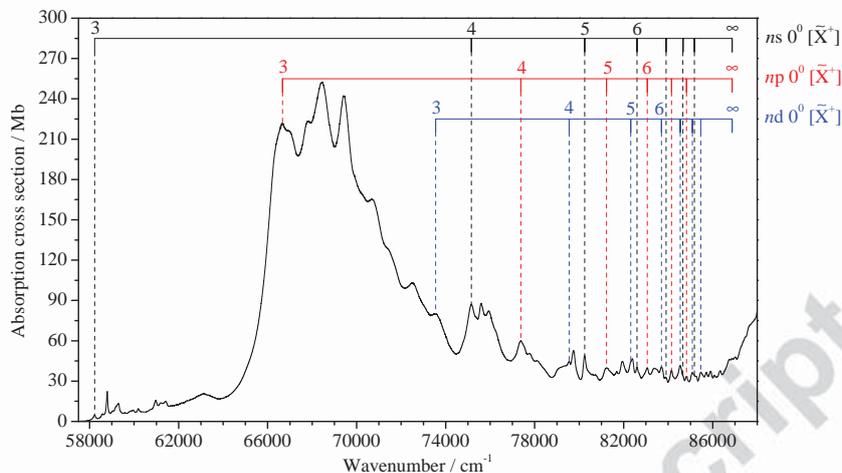


Figure 5: Absorption spectrum of methylcyanoacetylene between 60000 and 88000 cm^{-1} with the assignment corresponding to the Rydberg series converging to the $\tilde{X}^+ \ ^2E (v^+ = 0)$ cationic state. The ns , np and nd series are reported in black, red and blue, respectively.

265 $3s$ state is slightly higher (1.04) than the other members of this series (0.93) (see Table 3). With our modest calculations, the transition denoted $2 \ ^1E \leftarrow \tilde{X} \ ^1A_1$ in Table 1 is securely interpreted as the one occurring between the highest occupied molecular orbital and the $3s$ Rydberg orbital, and it unambiguously corresponds to the system of vibrationally resolved bands observed in the 58000 - 64000 cm^{-1} range. Its relatively high strength may stem from the Herzberg-Teller effect, namely an "intensity borrowing" from $3 \ ^1A_1 \leftarrow \tilde{X} \ ^1A_1$ transition that was not accounted for in the present calculations. The $3p$ [\tilde{X}^+] and $3d$ [\tilde{X}^+] states are also in the range of this valence transition. Thus, it is very difficult to unambiguously conclude on their location without taking into account a possible Rydberg-valence interaction. Here, we choose to assign the structures which correspond to the closest quantum defect value of the corresponding series. Note that the calculated wavenumber for the transition towards the $3p$ states only differs by 670 cm^{-1} (see Tables 1 and 2).

The Rydberg series converging to the first excited electronic state of $\text{CH}_3\text{C}_3\text{N}^+$ ($\tilde{A}^+ \ ^2A_1 (v^+ = 0)$, $IP = 105354 \text{ cm}^{-1}$ [17]) are depicted in Fig. 6 (see Table 3

for the assignment of the vibrationally excited Rydberg states and figures provided as supplementary material). The average values of the quantum defects

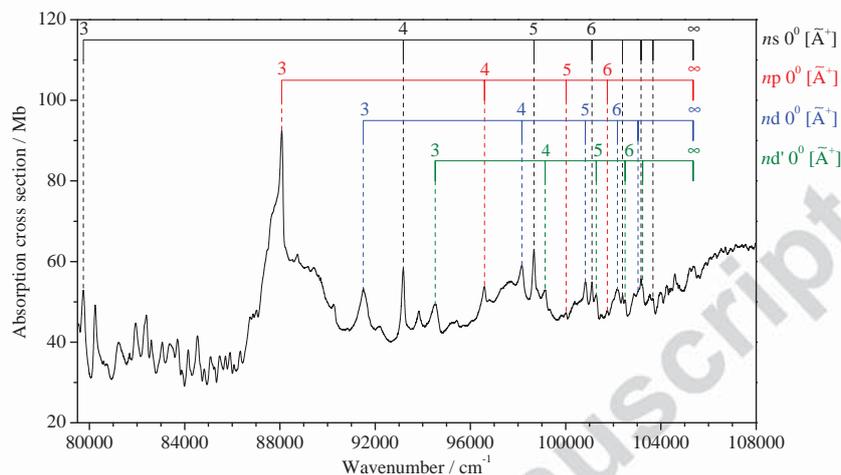


Figure 6: Absorption spectrum of methylcyanoacetylene between 79500 and 108000 cm^{-1} with the assignment corresponding to the Rydberg series converging to the $\tilde{A}^+ \ ^2A_1 (v^+ = 0)$ cationic state. The ns , np , nd and nd' series are reported in black, red, blue and green, respectively.

of the ns , np , nd and nd' series are 0.94, 0.47, 0.11 and -0.19, respectively. Individual quantum defects are reported in Table 3. Note that in this region, no
 285 Rydberg series converging to the $\tilde{B}^+ \ ^2E$ electronic state of the cation has been identified.

The absorption spectrum in the 108 000–120 000 cm^{-1} range is displayed in Fig. 7. Only one nd Rydberg series converging to the \tilde{C}^+ electronic state of $\text{CH}_3\text{C}_3\text{N}^+$ is identified. In this region, the Rydberg states are quite broad
 290 indicating a short lifetime of these states. This feature is very similar to that encountered in transitions towards the \tilde{C}^+ state of the cation in the threshold-photoelectron spectrum (TPES) [17]. The observed nd series has an average quantum defect of 0.15. Origin terms with a principal quantum number from 4 to 8 are presented in Fig. 7. The $n = 3$ term can be associated with four peaks
 295 between 105 500 and 106 500 cm^{-1} : 105 570, 105 865, 106 065 and 106 265 cm^{-1} .

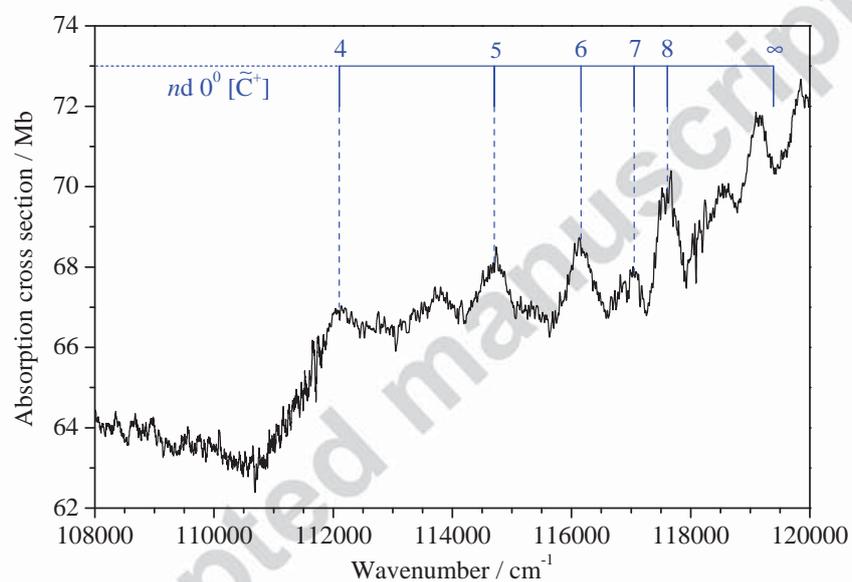


Figure 7: Absorption spectrum of methylcyanoacetylene between 108000 and 120000 cm⁻¹ with the assignment corresponding to the Rydberg series converging to the \tilde{C}^+ ($v^+ = 0$) cationic state. The nd series is reported in blue.

If we assign this term with one of these peaks, we find quantum defects of 0.18, 0.15, 0.13 and 0.11, respectively. No conclusion can be drawn from the quantum defect value only for the $n = 3$ Rydberg state hence it is not assigned here. The most active mode upon ionization towards the \tilde{C}^+ state is the ν_8^+ with a harmonic frequency of 1586 cm^{-1} [17]. In the TPES, a long progression along this mode was observed. Nevertheless, the absorption spectrum seems to possess only one series converging to the ground vibrational state of the \tilde{C}^+ electronic state. Only one of the assigned Rydberg states ($4d\ 8_0^1[\tilde{C}^+]$) seems to exhibit one vibrational state which might be assigned to the ν_8 mode. However, since the bands are broader in this region, it is difficult to conclude.

Our tentative assignment of the different Rydberg series described above is summarized in Table 3. Only 74 over more than 130 observed bands have been assigned in this work and high level *ab initio* calculations are needed to support a complete spectral analysis and to investigate the possible Rydberg-valence interactions.

4.4. Photoionization cross section

Based on the quasi perfect match observed between the absorption band intensities and the autoionization structure intensities observed in the $\text{CH}_3\text{C}_3\text{N}^+$ ion yield in Ref. [17] in the $100\ 000\text{--}105\ 000\text{ cm}^{-1}$ region, it is reasonable to assume that the photoionization quantum yield is close to unity in this spectral range. Indeed, it is quite unlikely that other processes (photodissociation for instance) would present exactly the same structures with the same relative intensities in this range. Within this assumption, we have scaled the photoionization yield of the $\text{CH}_3\text{C}_3\text{N}^+$ parent ion [17] on an absolute cross section scale and the resulting spectrum is displayed in Fig. 8. Note that this scaling only gives an upper limit estimate of the ionization cross section.

Between the first ionization potential and *c.a.* $100\ 000\text{ cm}^{-1}$, the difference between the two curves is due to neutral dissociation processes which compete with ionization. Above $105\ 000\text{ cm}^{-1}$, the photoionization cross section decreases whereas the photoabsorption cross section increases. This can eas-

Table 3: Position of the observed peaks in the absorption spectrum with the corresponding assignment. For Rydberg series converging to the electronic states of $\text{CH}_3\text{C}_3\text{N}^+$, the quantum defect δ is given in parentheses. All values are in cm^{-1} .

$\tilde{\nu}$	Assignment (δ)	$\tilde{\nu}$	Assignment (δ)	$\tilde{\nu}$	Assignment (δ)
58230	3s 0 ⁰ [X ⁺] (1.04)	85260	6p 2 ¹ [X ⁺] (0.62)	99810	4d' 6 ¹ [A ⁺] (-0.21)
60190	3s 2 ¹ [X ⁺] (1.05)	85471	9d 0 ⁰ [X ⁺] (0.15)	100020	5p 0 ⁰ [A ⁺] (0.46)
66670	3p 0 ⁰ [X ⁺] (0.67)	85900	6d 2 ¹ [X ⁺] (0.10)	100383	4d 2 ¹ [A ⁺] (0.10)
73568	3d 0 ⁰ [X ⁺] (0.13)	86072	7s 2 ¹ [X ⁺] (0.93)	100825	5d 0 ⁰ [A ⁺] (0.08)
75155	4s 0 ⁰ [X ⁺] (0.94)	86340	7p 2 ¹ [X ⁺] (0.64)	101104	6s 0 ⁰ [A ⁺] (0.92)
77380 ^a {	4s 2 ¹ [X ⁺] (0.93)	86755	7d 2 ¹ [X ⁺] (0.09)	101280	5d' 0 ⁰ [A ⁺] (-0.19)
	4p 0 ⁰ [X ⁺] (0.60)	86873	8s 2 ¹ [X ⁺] (0.90)	101458	5d 6 ¹ [A ⁺] (0.09)
79552 ^a {	4d 0 ⁰ [X ⁺] (0.13)	87022	8p 2 ¹ [X ⁺] (0.64)	101753 ^a {	6p 0 ⁰ [A ⁺] (0.48)
	4p 2 ¹ [X ⁺] (0.60)	88076	3p 0 ⁰ [A ⁺] (0.48)		6s 6 ¹ [A ⁺] (0.92)
79745	3s 0 ⁰ [A ⁺] (0.93)	88740	3p 6 ¹ [A ⁺] (0.48)	101970	5d' 6 ¹ [A ⁺] (-0.21)
80244	5s 0 ⁰ [X ⁺] (0.93)	90250	3p 2 ¹ [A ⁺] (0.49)	102175	6d 0 ⁰ [A ⁺] (0.12)
81230	5p 0 ⁰ [X ⁺] (0.59)	91503	3d 0 ⁰ [A ⁺] (0.19)	102385	7s 0 ⁰ [A ⁺] (0.92)
81688	4d 2 ¹ [X ⁺] (0.14)	92161	3d 6 ¹ [A ⁺] (0.18)	102495	6d' 0 ⁰ [A ⁺] (-0.20)
82320	5d 0 ⁰ [X ⁺] (0.09)	93180	4s 0 ⁰ [A ⁺] (1.00)	102860	6d 6 ¹ [A ⁺] (0.10)
82395	5s 2 ¹ [X ⁺] (0.94)	93710	3d 2 ¹ [A ⁺] (0.19)	103040 ^a {	7d 0 ⁰ [A ⁺] (0.11)
82600	6s 0 ⁰ [X ⁺] (0.93)	93830	4s 6 ¹ [A ⁺] (1.00)		7s 6 ¹ [A ⁺] (0.92)
83063	6p 0 ⁰ [X ⁺] (0.63)	94520	3d' 0 ⁰ [A ⁺] (-0.18)	103165 ^a {	8s 0 ⁰ [A ⁺] (0.92)
83374	5p 2 ¹ [X ⁺] (0.60)	95205	3d' 6 ¹ [A ⁺] (-0.19)		6d' 6 ¹ [A ⁺] (-0.21)
83700	6d 0 ⁰ [X ⁺] (0.12)	95430	4s 2 ¹ [A ⁺] (1.00)	103230	7d' 0 ⁰ [A ⁺] (-0.19)
83873	7s 0 ⁰ [X ⁺] (0.95)	96590	4p 0 ⁰ [A ⁺] (0.46)	103665	9s 0 ⁰ [A ⁺] (0.94)
84152	7p 0 ⁰ [X ⁺] (0.65)	96820	3d' 2 ¹ [A ⁺] (-0.19)	112100	4d 0 ⁰ [C ⁺] (0.12)
84539 ^a {	7d 0 ⁰ [X ⁺] (0.14)	97295	4p 6 ¹ [A ⁺] (0.45)	113760	4d 8 ¹ [C ⁺] (0.10)
	5d 2 ¹ [X ⁺] (0.07)	98160	4d 0 ⁰ [A ⁺] (0.09)	114705	5d 0 ⁰ [C ⁺] (0.16)
84665	8s 0 ⁰ [X ⁺] (0.95)	98670	5s 0 ⁰ [A ⁺] (0.95)	116160	6d 0 ⁰ [C ⁺] (0.17)
84823 ^a {	8p 0 ⁰ [X ⁺] (0.68)	98810 ^a {	4d 6 ¹ [A ⁺] (0.10)	117050	7d 0 ⁰ [C ⁺] (0.15)
	6s 2 ¹ [X ⁺] (0.90)		4p 2 ¹ [A ⁺] (0.47)	117610	8d 0 ⁰ [C ⁺] (0.15)
85088	8d 0 ⁰ [X ⁺] (0.16)	99135	4d' 0 ⁰ [A ⁺] (-0.20)		
85180	9s 0 ⁰ [X ⁺] (0.95)	99335	5s 6 ¹ [A ⁺] (0.95)		

^a two possible assignments.

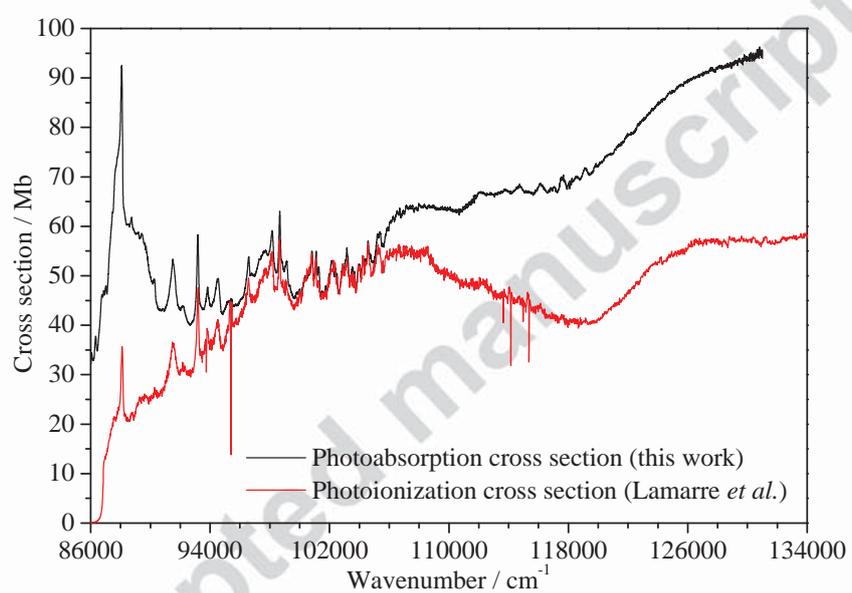


Figure 8: Tentative scaling of the CH₃C₃N⁺ ion yield measured by Lamarre *et al.* [17] (red line) on an absolute cross section scale using its photoabsorption cross section measured in the present work (black line) (see text for details).

ily be explained by invoking the opening of dissociative ionization channels of $\text{CH}_3\text{C}_3\text{N}$. Indeed, at this wavenumber the dissociative ionization process starts to be non-negligible (see Ref. [17]). Note that the ion yield reported in Ref. [17] only corresponds to the $\text{CH}_3\text{C}_3\text{N}^+$ parent cation signal and does not represent the total ion yield including the ionic fragments produced by dissociative ionization.

5. Conclusion

Absolute photoabsorption cross section of methylcyanoacetylene has been recorded for the first time over the 44 500–130 000 cm^{-1} range using the Fourier-transform spectrometer of the DESIRS beamline at SOLEIL synchrotron facility. Only the low energy range of our measurement in the region of a valence state transition has been published previously [5]. In this work, we propose a new and complete assignment of the vibrational progressions observed for this transition. A first assignment of the Rydberg series converging to different electronic states of the cation is also proposed. Preliminary calculations have been performed by TD-DFT method to support the analysis but more sophisticated calculations are needed to provide a deeper insight into the highly-excited electronic structure of methylcyanoacetylene.

From the structure intensities observed above the first ionization potential in our absorption spectrum and in the $\text{CH}_3\text{C}_3\text{N}^+$ photoion yield previously measured in our group [17], we conclude that the ionization yield in the 100 000–105 000 cm^{-1} range is close to unity which allows a scaling of the parent ion yield on an absolute cross section scale.

The high-quality of the absorption spectrum reported in the present work will be of great help to validate future high-level *ab initio* calculations of the complex excited electronic structure of this molecule.

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Appendix A. Plot of the absorption cross section versus n^*

A plot example of the absorption cross section against the effective quantum number (n^*) is presented here for the Rydberg series converging to the
365 $\tilde{X}^+ \ ^2E (v^+ = 0)$ ground state of $\text{CH}_3\text{C}_3\text{N}^+$. This plot allows to vertically visualize the different members of a pure Rydberg series (without any interaction). Note that the panel corresponding to the $n^* \in [2; 3]$ range is displayed to show that the 3s Rydberg state (identified by TD-DFT calculations) is probably perturbed by a Herzberg-Teller interaction with the close-lying valence state.

370 Appendix B. Supplementary material

Supplementary figures associated with this paper and displaying the complete Rydberg series assignment can be found in the online version at XXXX.

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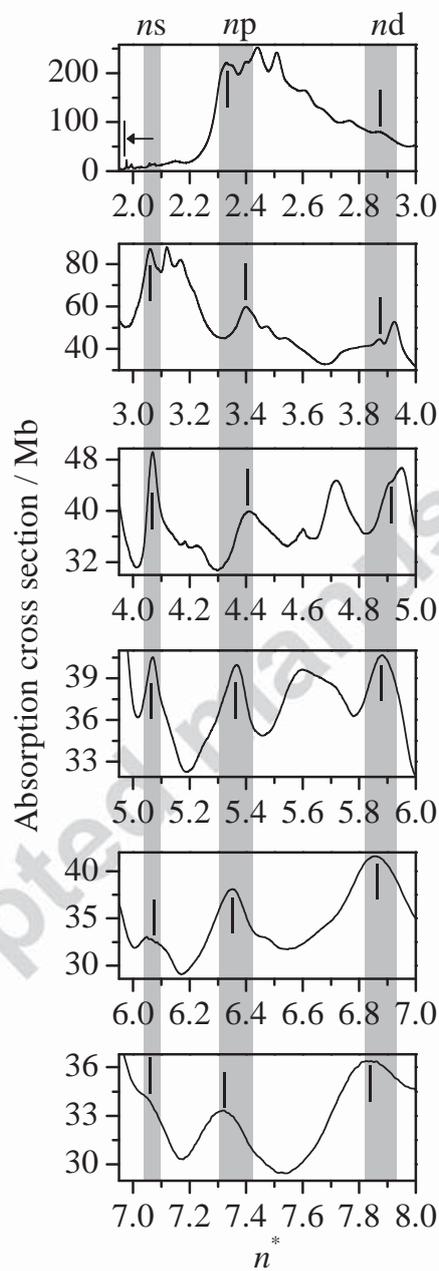


Figure A.1: Plot of the absorption cross section of methylcyanoacetylene against the effective quantum number n^* ($n^* = n - \delta$) for the Rydberg series converging to the $\tilde{X}^+ \ ^2E$ ($v^+ = 0$) state of $\text{CH}_3\text{C}_3\text{N}^+$. The gray areas link vertically bands with approximately the same quantum defect (δ) for each value of the principal quantum number (n). Each selected band is located by a vertical stick. The arrow in the upper panel shows that the 3s state is located at a slightly lower energy than expected, probably due to an interaction with the 3^1A_1 valence state (see section 4).

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