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# **A new lateral method for characterizing the electrical conductivity of ion-exchange membranes**

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#### **Abstract**

A new method for determining the electrical conductivity of ion-exchange membranes was implemented with four commercial membranes (AMX, CMX, MK-40 and MA-41). It is based on lateral resistance measurements without direct contact between electrodes and membranes. The cell configuration made it possible to determine the membrane conductivity over a wide range of electrolyte concentrations (measurements were carried out in the range  $10^{-5}$  - 5x10<sup>-1</sup> M). The structural parameters of the different membranes were inferred from AC conductivities and the microheterogeneous model. They were found in good agreement with literature results obtained by normal measurements (i.e. with current lines oriented normally to the membrane surface), thus confirming the reliability of the proposed method. The main advantage of our method is the possibility to characterize ion-exchange membranes even at low salt concentration unlike usual non-contact methods based on normal measurements.

Counterion diffusion coefficients and transport numbers within the membrane gel phase were estimated. The inter-gel phase volume fraction of an anion-exchange membrane (AMX) was also determined for the first time from membrane DC conductivity inferred from streaming potential and streaming current measurements. An excellent agreement was found with structural parameters obtained from AC measurements.

**Keywords:** Ion-exchange membranes, Electrical conductivity, Lateral measurements, Microheterogeneous model, Electrokinetics

#### 1. **Introduction**

Ion-exchange membranes (IEMs) are used in various industrial processes and electrochemical devices such as electrodialysis or fuel cells, respectively. Electrical conductivity is one of the most important characteristics of IEMs that defines their practical suitability. For example, in electrodialysis the electrical potential gradient (driving force) applied through an IEM, and so the specific energy consumption required to operate a given separation, depend on the membrane electrical conductivity. This latter can be determined from electrical resistance measurements. It has been reported, however, that the experimental value of the electrical conductivity may depend on the measuring method [1-7].

The different methods that can be implemented to measure membrane resistance can be split into different categories according to three main criteria, namely (i) the current lines can be oriented normally or parallel to the membrane surface, (ii) electrodes can be in contact with the membrane or not, (iii) direct current (DC) or alternating current (AC) can be used.

Normal measurements performed in both AC and DC modes were reported in the literature. In AC mode, the most used techniques are (i) the difference method in which the membrane resistance is obtained from the difference between the cell resistance measured with and without the membrane [8-24], and (ii) the contact-mercury method [25-33]. The difference method does not allow characterizing IEMs in too dilute solutions since it becomes highly inaccurate as the solution resistance increases too much with respect to that of the membrane (in most reported works measurements were performed with electrolyte concentrations higher than  $\sim 0.01$  M). On the other hand the contact-mercury method permits, in principle, characterizing IEMs in solutions of any concentration. However, measurement can be impacted by the partial drying (de-swelling) of membrane samples when transferred from the equilibrating solution to the measuring cell and poorly reproducible results were reported with this technique [2].

The use of a highly toxic metal (mercury) is another drawback of this method. Measurements performed in DC mode were also reported from different cell configurations including the Guillou's cell [1], a six-compartment cell with a four-electrode arrangement [34-36] or a twocompartment cell with a four-electrode arrangement [37]. A major drawback of methods involving DC measurements is the formation of diffusion boundary layers at membrane/solution interfaces, which impacts resistance measurements, particularly at low electrolyte concentration [1].

Electrical resistance measurements performed with lateral configuration were also reported in AC mode [5, 7, 38-43]. Most often, the membrane ends are put in contact with metal foils or wire electrodes (direct contact configuration). It has been argued that in such a configuration conduction may be only superficial and then experimental data may not reflect the membrane bulk conductivity [1].

In the present work we propose a new method to determine IEM electric conductivity over a wide range of electrolyte concentrations. Notably, it enables accurate measurements at low electrolyte concentrations unlike standard methods based on normal measurements. The method is based on lateral measurements performed with a measuring cell without electrode/membrane contact. Measurements were carried out with a commercial electrokinetic analyzer (SurPASS instrument, Anton Paar GmbH) allowing direct resistance measurements in AC mode as well as indirect DC resistance determination from both streaming current and streaming potential measurements. Structural parameters of four commercial IEMs (AMX, CMX, MK-40 and MA-41) were obtained from the microheterogeneous model and further combined with the Donnan exclusion model in order to assess ion transport numbers through the membrane gel phase.

#### **2. Experimental**

#### **2.1. Membranes and chemicals**

Four commercial IEMs were used in this work, two cation-exchange membranes (CMX Neosepta®, Tokuyama Soda, Japan; MK-40, Shchekinoazot, Russia), and two anionexchange membranes (AMX Neosepta®, Tokuyama Soda, Japan; MA-41, Shchekinoazot, Russia).

AMX and CMX membranes are homogeneous membranes made of functionalized polystyrene cross-linked with divinylbenzene and mixed with finely powdered poly(vinylchloride) (PVC), which are coated on a PVC cloth used as a reinforcing material.

The membrane fixed-charge comes from the presence of sulfonate groups ( −¿  $-SO_3^{\lambda}$  ) in CMX

membrane and from the presence of quaternary ammonium groups  $+\dot{\iota}(C H_3)_3$  $-N^{\dot{b}}$ ¿ in AMX

membrane. Both charged groups are grafted directly to the basic polymer structure so that ionic charges are distributed over the whole membrane material.

On the other hand, MA-41 and MK-40 heterogeneous membranes are composite materials made of ion-exchange resins embedded in polyethylene. Ion-exchange resins consist in

$$
+ \lambda (CH_3)_3
$$
  
functionalized polystyrene (with  $-N^{\lambda}$  groups for MA-41 membrane and  $-SO_3^{\lambda}$ 

groups for MK-40 membrane) cross-linked with divinylbenzene.

The main properties of these IEMs are collected in Table 1.

Electrolyte solutions were prepared from KCl and NaCl (Fisher Scientific, analytical grade) and deionized water (resistivity: 18 M $\Omega$  cm). They were further used without pH adjustment (pH was found to be  $5.7\pm0.1$  for all solutions).

**Table 1.** Main properties of AMX, CMX, MA-41 and MK-40 membranes.



<sup>a</sup> Manufacturer data; Refs. [44-45]; <sup>b</sup> Ref. [33]; <sup>c</sup> Ref. [46]; <sup>d</sup> Ref. [47].

#### 2.2. **Membrane thickness measurement**

Membrane thickness was measured with a Dial Indicator (model 2046S, Mitutoyo, Japan) equipped with a 10 mm diameter flat contact point. Membrane samples were first soaked in a highly concentrated NaCl or KCl solution in order to exchange original membrane counterions with  $Na^+$  or  $K^+$  (CMX and MK-40 membranes) or Cl (AMX and MA-41 membranes). Samples were further equilibrated with the measuring electrolyte solution at room temperature (24  $\pm$  1 °C) for 24 h. Samples were then extracted from the solution and thickness was measured. For each sample, seven measurements were performed at different locations and the average value was considered as the membrane thickness  $(e_m)$ .

#### **2.3. Electrical resistance measurement**

Membrane electrical resistance were measured with a SurPASS electrokinetic analyzer (Anton Paar GmbH, Austria) after thickness measurements.

The cell we used required two samples of the same membrane materials. Each membrane sample was cut and adjusted to the sample holder dimensions (length (L): 2 cm and width (W): 1 cm) and fixed using double-sided adhesive tape. To prevent any leakage between the membranes and the sample holders, the membranes were firmly pressed against sample holders for 30–60 s.

Measurements were conducted with an adjustable-gap cell inside which both membrane samples face each other (see schematic of the measuring cell shown in Fig. 1). Thanks to micrometric screws the distance between the samples  $(h_{ch})$  could be set and varied without dismounting the cell [48]. The cell was surrounded by two cylindrical Ag/AgCl electrodes through which the measuring solution could flow (Fig. 1).

Since membrane swelling was likely to be modified when fixing the membranes onto the sample holders, circulating the measuring solution through the flow-cell enabled recovering equilibrium conditions prior to electrical resistance measurement. The solution flow was

created by a pair of syringe pumps and the resistance was measured in AC mode ( $r_{cell_{AC}}$ 

 by means of the conductivity meter incorporated in the SurPASS instrument and a pair of reversible Ag/AgCl electrodes with surface area of 10 cm<sup>2</sup> each.



**Fig. 1.** Schematic representation of the adjustable-gap cell used for electrical resistance measurements (adapted from [48]).

The SurPASS electrokinetic analyzer also allows determining the membrane electrical resistance from electrokinetic measurements, which is equivalent to measuring the electrical resistance in DC mode. In this case, the measuring solution is forced through the streaming channel formed by the two membrane samples facing each other (see Fig. 1). The fluid therefore drives ions within the mobile diffuse layer towards the low-pressure side.

This local ion shifting gives rise to a convective current, defined as the streaming current  $(I_s)$ that can be measured by the SurPASS instrument (streaming current mode). If a highimpedance voltmeter is put in the external loop, then the accumulation of ions at the channel end sets up an electric field which drives the counter-ions to move back in the opposite direction to the pressure-driven flow. This back-flow of counter-ions generates an electrical conduction current in the opposite direction to the streaming current (not only between the membranes but also through the membranes). When a steady state is reached, the conduction current and the streaming current balance each other. The resulting electrical potential difference that can be measured between the channel ends is defined as the streaming

potential  $\Delta \varphi$ <sub>s</sub>. The cell electrical resistance in DC mode *r*  $(i\textit{i}$  cell<sub>DC</sub>) ¿ can therefore be

inferred from both streaming potential and streaming current measurements by applying Ohm's law:

$$
r_{cell_{DC}} = \frac{\Delta \varphi_s}{I_s} \tag{1}
$$

Both streaming potential and streaming current were measured by applying pressure differences up to 300 mbar. The pair of syringe pumps used to circulate the measuring solution through the cell allowed reversing the flow direction periodically. This helped to limit electrode polarization during electrokinetic measurements [49].

All  $r_{cell_{AC}}$  and  $r_{cell_{DC}}$  experiments were performed at room temperature (24 ± 1 °C). The distance between the membrane samples  $(h_{ch})$  was set to  $100 \pm 2 \mu$ m (unless specified).

This latter was determined from volume flow rate  $(Q_v)$  measurements performed at various pressure difference  $(\Delta P)$  and the Hagen–Poiseuille equation, which reads as follows for a channel of rectangular cross-section (with  $W>>h_{ch}$ ) [50]:

$$
Q_V = \frac{2 W \left(\frac{h_{ch}}{2}\right)^3}{3 \eta L} \Delta P
$$

(2)

where  $\eta$  is the solution viscosity.

From the electrical point of view, the measuring cell consists in an electrical circuit with three parallel connected branches, i.e. membrane / solution-filled channel / membrane (Fig. 2). Consequently, it is suited for characterizing IEMs in diluted electrolyte solutions (at low concentrations, the channel resistance is much greater than that of the membrane) unlike most usual cells for which the membrane is surrounded by two (identical) electrolyte solutions thus forming a series circuit.



**Fig. 2.** (a) Different electrical current paths through the measuring cell; (b) equivalent electrical circuit; i<sub>m</sub> is the electrical current flowing through each membrane sample, i<sub>ch</sub> is the electrical current flowing between membranes,  $r_m$  is the membrane electrical resistance and  $r_{ch}$ solution-filled channel resistance.

The cell resistance that is accessible experimentally  $(r_{\text{cell}})$  is then related to both the membrane resistance  $(r_m)$  and the solution-filled channel resistance  $(r_{ch})$  as follows:

$$
\frac{1}{r_{Cell}} = \frac{1}{r_{ch}} + \frac{2}{r_m}
$$
 (3)

The membrane resistance can therefore be determined from:

$$
r_m = \frac{2 r_{ch} r_{Cell}}{r_{ch} - r_{Cell}}
$$

(4)

where the channel resistance  $r_{ch}$  is given by:

$$
r_{ch} = \dot{\zeta} \qquad \frac{L}{W h_{ch} \kappa_{ch}}
$$

(5)

with L, W and  $h_{ch}$  the channel length, width and thickness, respectively, and  $\kappa_{ch}$  the electrical conductivity inside the channel. This latter was considered equal to the measuring solution conductivity since the distance between the membranes samples (100  $\pm$  2  $\mu$ m unless specified) was always several orders of magnitude greater than the Debye length in the measuring solutions.

The membrane conductivity  $(\kappa_m)$  can be inferred from the membrane resistance as follows:

$$
\kappa_m = \frac{L}{W e_m r_m}
$$

(6)

where  $e_m$  is the wet membrane thickness.

In addition to the electrical conductivity, IEMs are commonly characterized by their specific resistance. The electrical resistance depends on the cell configuration (normal or lateral measurements; see Fig. 3), unlike the membrane conductivity. Although the apparent membrane thickness was L when lateral measurements were performed (see Fig. 3), the specific resistance  $(R_m)$  should be computed by considering the actual membrane thickness, i.e. em:

$$
R_m = \frac{e_m}{\kappa_m} \tag{7}
$$



**Fig. 3.** Effective membrane thickness for the different cell configurations. a) actual thickness  $e<sub>m</sub>$  in the case of normal measurements, b) apparent thickness L (corresponding actually to the sample length) in the case of lateral measurements (this configuration was used in the present work).

As mentioned above, electrical resistance measurements were performed in lateral configuration by means of an adjustable-gap cell.

It was then possible to perform a series of measurements by varying the distance  $h_{ch}$  between the two membrane samples facing each other in the measuring cell. From Eqs. (3) and (5) it follows that the reciprocal of the cell resistance is expected to vary linearly with  $h_{ch}$ :

$$
\frac{1}{r_{cell}} = \frac{W\kappa_{ch}}{L}h_{ch} + \frac{2}{r_m}
$$

(8)

Performing measurements for various h<sub>ch</sub> therefore allows determining the membrane resistance by plotting  $1/r_{\text{Cell}}$  as a function of h<sub>ch</sub> and extrapolating  $1/r_{\text{Cell}}$  for h<sub>ch</sub> = 0.

#### 3. **Results and discussion**

#### **3.1 Validation of lateral measurements**

We first demonstrated the reliability of our measurements performed in lateral configuration by comparing the specific resistances of four IEMs reported in the literature with our own measurements. These latter were performed in AC mode with a 0.5 M NaCl solution, which is a standard solution used by IEM membrane manufacturers to characterize their membranes.

Specific resistances obtained from our new lateral method were compared with both manufacturer's data and some literature results. All data are collected in Table 2. A very good agreement was obtained between our measurements and manufacturer's data. Moreover, our results lie in the range of experimental data reported in the literature with different methods based on normal measurements (membrane and measuring solutions in series). This is a strong indication that current lines did flow through the whole membrane matrix during our lateral measurements without contact between membranes and electrodes, and not only along the membrane surfaces. As mentioned above, measurements were carried out with a quite concentrate electrolyte solution (0.5 M NaCl). It should be stressed that lateral measurements are expected to be even more accurate when performed with more dilute solutions because of the parallel association of the membrane samples and the measuring solution (Fig. 2).





#### **3.2 Membrane electrical conductivity**

Membrane electrical conductivity was determined in KCl solutions of various concentrations ranging from  $10^{-5}$  to 0.5 M. Fig. 4 shows the variation of the membrane conductivity with KCl concentration. Overall, the trend of the different curves is similar to that described in literature [1-2, 8-9, 59-66] with a sharp increase in the membrane conductivity observed at low concentrations for all IEMs followed by a less pronounced rise for salt concentrations higher than *ca* 0.01 M (note that most data reported in the literature were obtained for concentrations higher than 0.01 M while our method enables accurate measurements at much lower concentrations).



**Fig. 4.** Logarithm of membrane conductivity vs. logarithm of electrolyte concentration (KCl solutions).

Additional measurements were performed by varying the distance between the membrane samples in the lateral measurement cell. Fig. 5 gives an illustration of results obtained with the AMX membrane in a 0.5 M KCl solution.

As expected from Eq. (8) the reciprocal of the cell resistance was found to vary linearly with the distance between the membrane samples. The membrane resistance was then inferred from the y-axis intercept without having to determine the solution conductivity explicitly. The AMX membrane resistance obtained from this extrapolation method was found in good agreement with the value obtained from single-gap measurements (distance between the membrane samples:  $100 \pm 2 \mu$ m) and Eq. (4) as shown in Table 3 (the deviation between the two different methods was about 5%).



Fig. 5. Reciprocal of the cell resistance vs. gap between membrane samples  $(h_{ch})$  in the lateral-measurement cell (AMX membrane in 0.5 M KCl).

**Table 3.** AMX membrane resistance measured in 0.5 M KCl by both single-gap and extrapolation methods.

Method	Membrane resistance	Membrane specific	
	$r_{m}$ (kΩ)	resistance	
		$R_m(\Omega.cm^2)$	
Single-gap measurement	20.12	2.58	
Extrapolation method	19.25	2.46	

According to Eq. (8) the solution conductivity inside the channel ( $\kappa_{ch}$ ) can be determined from the slope of the line shown in Fig. 5. It is worth mentioning that the deviation between  $\kappa_{ch}$ determined from Fig. 5 and the experimental solution conductivity (measured outside the cell) was less than  $0.3\%$  ( $\kappa_{ch}$  = 59.68 mS.cm<sup>-1</sup> from Fig. 5 and 59.84 mS.cm<sup>-1</sup> from direct measurement), which gives further evidence for the reliability of our measurements performed in lateral configuration.

#### **3.3 Membrane structure and transport properties**

The previous sections aimed at validating our new method for IEM characterization. The microheterogeneous model developed by Zabolotsky and Nikonenko [67] was further applied in the present work in order to analyze IEM structure. Hydration of ions and hydrophilic parts of the polymer matrix results in swelling of IEMs, which leads to a substantial reconstruction of their structure [68]. According to the microheterogeneous model, a hydrated IEM can be considered as a combination of two distinct microphases: the gel phase that contains a relatively uniform distribution of ionic groups and counter-ions compensating the fixed charge density, and the inter-gel phase filled with the electroneutral solution [67,69]. The conductivity in the gel phase  $(\kappa_1)$  is due to counterions compensating the membrane fixedcharge whereas the conductivity of the interstitial solution filling the inter-gel phase  $(\kappa_2)$ results from both coions and counterions from the measuring solution. In the microheterogeneous model the inter-gel phase conductivity is assumed to be the same as that of the measuring solution, i.e.  $\kappa_2 = \kappa_{ch}$ .

IEM structure is characterized by the gel and inter-gel phase volume fractions referred to as  $f_1$ and  $f_2$ , respectively  $(f_1 + f_2 = 1)$ . IEMs are usually classified according to the value of  $f_2$ , namely homogeneous membranes are characterized by  $0 \le f_2 \le 0.15$ , heterogeneous membranes are characterized by  $0.15 \le f_2 \le 0.25$  while membranes are considered as porous if  $0.3 \le f_2 \le 0.45$ [46,69].

Within the scope of the microheterogeneous model the membrane AC conductivity ( $\kappa_{m,AC}$ ) can be expressed as follows [67]:

$$
\kappa_{m,AC} = (f_1 \kappa_1^{\alpha} + f_2 \kappa_2^{\alpha})^{1/\alpha} \tag{9}
$$

where  $\alpha$  is a structural parameter reflecting the relative arrangement of the gel and inter-gel phases (  $-1 \le \alpha \le +1$  with  $\alpha = -1$  for a series phase connection while  $\alpha = +1$  for a parallel phase connection).

The isoconductance concentration  $(c_{iso})$  is defined as the electrolyte concentration for which the membrane conductivity, the gel phase conductivity and the measuring solution conductivity are equal  $(\kappa_m = \kappa_1 = \kappa_{ch} = \kappa_{iso})$ . Isoconductance concentrations  $(c_{iso})$  and conductivities at the isoconductance point  $(\kappa_{iso})$  obtained in the present work for the various IEMs are given in Table 4. The isoconductance concentration was found around 0.05 mol.L-1 for the different membranes.

**Table 4.** IEM parameters determined in the present work from AC resistance lateral measurements.

Membrane	<b>AMX</b>	$MA-41$	<b>CMX</b>	$MK-40$
Isoconductance concentration, $c_{iso}$ (mol. L <sup>-1</sup> )	0.028	0.030	0.065	0.042
Conductivity at isoconductance point, $\kappa_{iso}$ (mS.cm <sup>-1</sup> )	3.72	3.99	8.79	5.64
Volume fraction of inter-gel phase, $f_2$	0.11	0.24	0.10	0.25
Structural parameter, $\alpha$	0.20	0.15	$-0.21$	0.16
Counterion diffusion coefficient in the gel phase,	5.20	4.77	9.97	4.33
$\acute{D}_{\text{count}} \times 10^7 \text{ (cm}^2 \text{.} \text{s}^{-1})$				
Counterion diffusion coefficient in the gel phase,	0.989	0.986	0.992	0.991
$t_{\textit{count}}$				

Near the isoconductance point Eq. (9) shows that  $\kappa_m$  depends only slightly on  $\alpha$ . When

 $|\alpha| \ll 1$ , which corresponds to a quite random arrangement of gel and inter-gel phases, Eq. (9) can be approximately rewritten as follows [67]:

$$
\kappa_{m,AC} = \kappa_1^{f_1} \kappa_2^{f_2} \tag{10}
$$

and then,

$$
\ln \kappa_{m,AC} = f_1 \ln \left(\kappa_1\right) + f_2 \ln \left(\kappa_2\right) \tag{11}
$$

Zabolotsky and Nikonenko showed that Eq. (11), which predicts a linear variation of ln  $\kappa_{m,AC}$ vs. In  $\kappa_2$ , is a satisfying approximation in the concentration range  $\sim 0.1$  c<sub>iso</sub>  $\lt$   $\kappa_{\rm iso}$   $\lt$   $\sim 10$  c<sub>iso</sub> up

$$
_{\text{to}} \quad |\alpha| = 0.2
$$

.

Fig. 6 shows the experimental ln  $\kappa_{m,AC}$  vs. ln  $\kappa_{ch}$  dependence for AMX, CMX, MK-40 and MA-41 IEMs. Measurements were performed in KCl solutions of various concentrations ranging from  $10^{-5}$  to  $5x10^{-1}$  M. Results are in good agreement with the microheterogeneous model since a linear variation of ln  $\kappa_{m,AC}$  as a function of ln  $\kappa_{ch}$  (as mentioned above, this model assumes  $\kappa_2 = \kappa_{ch}$ ) was obtained with all IEMs for  $\sim$ 0.5  $c_{iso} < c_{iso} < \sim$ 10  $c_{iso}$  (Fig.6b).

The volume fraction of the inter-gel phase  $(f_2)$  of the different IEMs was determined from Eq. (11) and the slope of lines obtained in the concentration range  $\sim$ 0.5  $c_{\text{iso}}$  <  $c_{\text{iso}}$  <  $\sim$ 10  $c_{\text{iso}}$  (Fig. 6b). Results are collected in Table 4. For both AMX and CMX membranes  $f_2$  were found to be less than 0.15, which corresponds to homogeneous membranes. On the other hand,  $f_2$  was found to be 0.24 and 0.25 for the heterogeneous MA-41 and MK-40 membranes, respectively. Our results are in good agreement with  $f_2$  values reported in the literature for these membranes (Table 5), which gives additional evidence for the reliability of our lateral measurements of electrical membrane resistance.



**Fig. 6.** a) Logarithm of membrane conductivity vs. logarithm of KCl solution conductivity. b) Linear variations obtained for  $\sim 0.5$   $c_{\text{iso}} < c_{\text{iso}} < \sim 10$   $c_{\text{iso}}$ .

Membran	f <sub>2</sub>	$\alpha$	$D_{count} \times 10^7$	Electrolyte	Reference
e					
			$(cm2.s-1)$		
<b>AMX</b>	0.24		5.1	$Cl^-$ form	[8]
	0.06	0.35		NaCl	$[27]$
	0.099		4.1	NaCl	$[33]$
	$0.10 \pm 0.02$			NaCl	$[47]$
	$0.10 \pm 0.01$			NaCl	[51]
	0.06		4.5	<b>NaCl</b>	$[54]$
	0.11			<b>NaCl</b>	[69]
<b>CMX</b>	0.06		11.61	<b>NaCl</b>	$[46]$
	$0.11 \pm 0.02$			<b>NaCl</b>	$[47]$
	$0.11 \pm 0.01$			<b>NaCl</b>	[51]
	0.06			<b>NaCl</b>	[60]
	0.04			HC <sub>1</sub>	[60]
$MA-41$	0.18	0.31		<b>NaCl</b>	$[27]$
	$0.22 \pm 0.02$			<b>NaCl</b>	$[47]$
	$0.20 \pm 0.03$	0.24	4.7	<b>NaCl</b>	[67]
	0.17			<b>NaCl</b>	$[69]$
	0.18	0.31		<b>NaCl</b>	$[70]$
	0.22			<b>NaCl</b>	$[71]$
	0.23			<b>NaCl</b>	$[72]$
	0.10	0.50		LiC <sub>1</sub>	$[73]$
$MK-40$	0.18		7.35	<b>NaCl</b>	[46]
	$0.19 \pm 0.02$			<b>NaCl</b>	$[47]$
	0.18			<b>NaCl</b>	[60]
	0.20			<b>HCl</b>	[60]
	0.23			<b>NaCl</b>	$[69]$
	0.19	0.33		LiCl	$[73]$

**Table 5.** IEM parameters inferred from normal measurements (literature results).

Knowing  $f_2$  and  $\kappa_{\text{iso}}$ , the structural parameter  $\alpha$  was further determined for each membrane by

fitting experimental data shown in Fig. 6a with Eq. (9). We obtained  $0.15 \le \sqrt{\alpha} \le 0.21$  for all membranes (see Table 4), which indicates a rather random connection between gel and inter-gel phases whatever the kind of IEMs (homogeneous or heterogeneous). It should be stressed that data obtained for highly diluted solutions ( $c \le 10^{-4}$  mol. L<sup>-1</sup>) were disregarded in the fitting procedure since they were found to yield irrelevant  $\alpha$  values. The reason might be that the assumption made in the microheterogeneous model that the inter-gel phase is filled with an electroneutral solution having the same conductivity as the external solution is no longer valid at very low ionic strengths.

Indeed, electrical double layers expand as the electrolyte concentration decreases. At sufficiently low concentration they may overlap inside the inter-gel phase cavities which would then no longer be filled with an electroneutral solution but would contain an excess of counterions (and so  $\kappa_2$  would be greater than  $\kappa_{ch}$ ) [74].

Fig. 7 shows the concentration dependence of  $\ln \frac{K_m}{\kappa_{ch}}$  for the different IEMs. At low concentrations the membrane conductivity was found higher than that of the external solution because ion concentration inside the membrane was much higher than the external bulk concentration. Moreover, we observed in the low concentration range that the variation of

*κm* with the electrolyte concentration was more pronounced for MA-41 and MK-40 heterogeneous membranes than for AMX and CMX homogeneous membranes (Fig. 7b).





**Fig. 7.** a) Logarithm of membrane conductivity to solution conductivity ratio vs. logarithm of KCl solution concentration b) Enlargement for very low concentrations  $(10^{-5} - 5 \times 10^{-3} M)$ .

As discussed above, the conductivity at the isoconductance point  $(\kappa_{iso})$  is an essential quantity to get information on IEM structural properties. It can also be used to determine the counterion diffusion coefficient ( $\hat{D}_{\text{count}}$ ) in the gel phase [8, 67]:

$$
\acute{D}_{count} = \frac{RT \,\kappa_{iso}}{F^2 Q} \tag{12}
$$

where R is the ideal gas constant, T the temperature, F the Faraday's constant and Q is the ion-exchange capacity of the gel phase, which can be obtained from the membrane ionexchange capacity  $(Q_m)$  and the volume fraction of the gel phase  $(f_1)$ :

$$
Q = \frac{Q_m}{f_1} \tag{13}
$$

The counterion diffusion coefficient in the gel phase of the different IEMs was found to lie in the range  $4 \times 10^{-7} - 10^{-6}$  cm<sup>2</sup>.s<sup>-1</sup> (see Table 4), i.e. 20 to 50 times smaller than in the bulk phase, which results from the strong attractive interaction between counterions and the fixed-charges in the IEM gel phase. These values lie in the range of experimental results obtained from normal measurements (Table 5).

The gel phase conductivity, which can be obtained from Eq. (10), can be expressed as:

$$
\kappa_1 = \frac{F^2}{RT} \Big( \acute{D}_{\text{co}} \acute{C}_{\text{co}} z_{\text{co}}^2 + \acute{D}_{\text{count}} \acute{C}_{\text{count}} z_{\text{count}}^2 \Big)
$$
(14)

where  $\hat{C}_{\square}$  and  $\hat{D}$  denote ion concentrations and diffusion coefficients in the gel phase, respectively, *z* is the ion charge number, and subscripts "co" and "count" stand for coions and counterions, respectively.

The electroneutrality condition in the gel phase reads as follows:

$$
z_{co}^{\square} \hat{C}_{co} + z_{count}^{\square} \hat{C}_{count} + \omega Q = 0 \tag{15}
$$

with  $\omega = +1$  for an anion-exchange membrane and -1 for a cation-exchange membrane.

The interfacial Donnan equilibrium is given by,

$$
K_i = \frac{\dot{C}_i}{C_i} = K^{z_i} \tag{16}
$$

where  $K_i$  is the Donnan partitioning coefficient for ion *i*,  $C_i$  its bulk concentration, and  $K$  is the Donnan equilibrium constant.

For mono-monovalent electrolytes Eqs. (15) and (16) lead to:

$$
\dot{C}_{co} = \frac{-\omega Q}{2} + \sqrt{C_{\square}^2 + \left(\frac{\omega Q}{2}\right)^2} \tag{17}
$$

$$
\dot{C}_{\text{count}} = \frac{\omega Q}{2} + \sqrt{C_{\text{u}}^2 + \left(\frac{\omega Q}{2}\right)^2} \tag{18}
$$

where *C* is the external electrolyte concentration.

The counterion transport number in the gel phase  $(\hat{t}_{count})$  can be determined around the isoconductance point from the following equation with the help of Eqs. (12), (14), (17) and (18):

$$
\hat{t}_{count} = \frac{|z_{count}^{\Box}| \hat{D}_{count} \hat{C}_{count}}{|z_{count}^{\Box}| \hat{D}_{count} \hat{C}_{count} + |z_{co}^{\Box}| \hat{D}_{co} \hat{C}_{co}} \tag{19}
$$

The transport numbers of counterions in the gel phase were found around 0.99 for all IEMs (Table 4).

As mentioned in section 2.3 the SurPASS instrument also allows determining the cell electrical resistance from electrokinetic measurements (hereafter referred to as the DC method), namely streaming current and streaming potential measurements, which is

equivalent to the DC resistance *r*  $(i\textit{i}$  cell $_{\textit{DC}})$ ¿ . Fig. 8 shows the AMX membrane conductivity

determined from both AC and DC methods for KCl solutions of various concentrations. To the best of our knowledge it is the first time that streaming potential and streaming current measurements are combined to determine the electrical conductivity of IEMs. Both methods

led to close values with  $\kappa_{m,AC}$  only slightly higher than  $\kappa_{m,DC}$ .

The transport number of counterions (Cl⊤) inside the gel phase of AMX membrane was estimated to be 0.989 (Table 4). When  $\hat{t}_{count} \rightarrow 1$  and  $|\alpha| \ll 1$  the DC conductivity and the AC conductivity can be linked by the following approximate equation [67]:

$$
\kappa_{m,DC} = \kappa_{m,AC} t_{count}^{f_2}
$$
 (20)

where  $t_{count}$  is the transport number of counterions in the external solution, i.e. nearly 0.5 for Cl<sup>−</sup> in KCl solutions.

According to Eq. (20) the difference between the logarithm  $\kappa_{m,AC}$  and that of  $\kappa_{m,DC}$  is then approximately equal to -0.69 $f_2$ . From results shown in Fig. 8 we obtained  $f_2 = 0.11$ (average value), which is identical to the value determined from AC measurements (Table 4).



**Fig.8.** Logarithm of the AMX membrane conductivity determined from DC and AC methods vs. logarithm of the KCl solution conductivity.

#### **4. Conclusion**

The electrical conductivity of four commercial IEMs (AMX, CMX, MK-40 and MA-41) was determined by a new method based on lateral resistance measurements without direct contact between electrodes and membranes. It was first demonstrated that the cell configuration permitted probing the entire membrane volume and that conduction did not occur only along the membrane surface. Moreover, the parallel arrangement of the membrane samples and the electrolyte solution inside the measuring cell made it possible to determine membrane conductivity even at very low electrolyte concentration unlike usual non-contact cells. The microheterogeneous model was used to estimate the structural parameters of the different IEMs from electrical resistance measurements performed in AC mode. Transport parameters (counterion diffusion coefficients and transport numbers) within the gel phase of IEMs were further determined from the membrane AC conductivity and the combination of the microheterogeneous model with the Donnan-exclusion approach.

Finally, the membrane DC conductivity was determined for the first time from both streaming potential and streaming current measurements. The inter-gel phase volume fraction of AMX membrane inferred from these electrokinetic measurements was found in excellent agreement with AC measurements.

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