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PRESSURE BROADENING AND SHIFTING COEFFICIENTS AS TESTS OF H₂-He POTENTIAL ENERGY SURFACES



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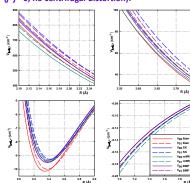
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We have used the Schaefer and Köhler¹ (SK), the modified Muchnick and Russek² (MR), the Boothroyd, Martin and Peterson³ (BMP), and the Bakr, Smith and Patkowski⁴ H₂-He potential energy surfaces in order to calculate, using the close coupling method, pressure broadening and shifting coefficients. The helium pressure broadening and shifting generalized cross sections of the isotropic raman Q(1) lines of the fundamental of D2 and H2 as well as the purely rotational Stokes S0(1) line of H2 were computed. We have decomposed the pressure broadening cross-sections in a purely inelastic and a purely dephasing, including a vibrational one for the Q lines, contributions. Such a decomposition allows us to better understand the main differences that exist between these potentials. The old Schaefer and Köhler PES and the most recent one, namely the Bakr et al PES, give close results in quite good agreement with available experimental data.

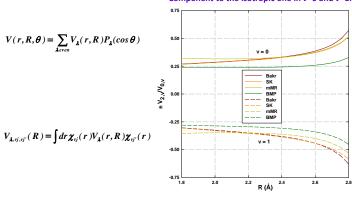
 $V(r,R,\theta) = \sum_{\lambda} V_{\lambda}(r,R) P_{\lambda}(\cos\theta)$

1. PESs used

Comparison of the isotropic components in v=0 and in v=1(j=j'=0, no centrifugal distortion):



Component to the isotropic one in v=0 and v=1:



Comparison of the ratio of the first anisotropic

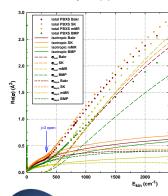
The modified MR PES shows the less differences in its isotropic parts in v=0 and v=1 while the SK PES shows the most differences. The modMR PES differs significantly from the others at long range. The BMP PES is the less anisotropic PES.

2. Pressure broadening and shifting cross-sections and coefficients

$$\sigma^{(n)}(v_ij_iv_fj_f;E_{kin}) = \sigma_{inelastic}(v_ij_i,v_fj_f;E_{kin}) + \sigma^{(n)}_{dephasing}(v_ij_i,v_fj_f;E_{kin}) \\ \sigma^{(n)}_{PB}(if;E_{kin}) = Re\Big(\sigma^{(n)}(if;E_{kin})\Big) \\ \sigma^{(n)}_{PS}(if;E_{kin}) = Im\Big(\sigma^{(n)}_{dephasing}(if;E_{kin})\Big) \\ \sigma^{(n)}_{PS}(if;E_{kin}) = Im\Big(\sigma^{(n)}_{dephasing}(if;E_{kin})\Big)$$

$$\sigma_{inelastic}(v_i,j_i,v_f,j_f;E_{kin}) = \frac{1}{2} \Biggl(\sum_{j_i \neq j_f} \sigma(v_i,j_i \rightarrow v_i,j_i';E_{kin}) + \sum_{j_i \neq j_f} \sigma(v_f,j_f \rightarrow v_f,j_f';E_{kin}) \Biggr) \\ \gamma + i\delta = n_b \overline{v} \Biggl(\frac{1}{k_B T} \Biggr)^2 \int dE_{kin} E_{kin} \exp(-E_{kin}/k_B T) \sigma^{(n)}(if;E_{kin}) + \sum_{j_i \neq j_f} \sigma(v_i,j_f \rightarrow v_f,j_f';E_{kin}) \Biggr) \\ \gamma + i\delta = n_b \overline{v} \Biggl(\frac{1}{k_B T} \Biggr)^2 \int dE_{kin} E_{kin} \exp(-E_{kin}/k_B T) \sigma^{(n)}(if;E_{kin}) \Biggr) \\ \gamma + i\delta = n_b \overline{v} \Biggl(\frac{1}{k_B T} \Biggr)^2 \int dE_{kin} E_{kin} \exp(-E_{kin}/k_B T) \sigma^{(n)}(if;E_{kin}) \Biggr) \\ \gamma + i\delta = n_b \overline{v} \Biggl(\frac{1}{k_B T} \Biggr)^2 \int dE_{kin} E_{kin} \exp(-E_{kin}/k_B T) \sigma^{(n)}(if;E_{kin}) \Biggr) \\ \gamma + i\delta = n_b \overline{v} \Biggl(\frac{1}{k_B T} \Biggr)^2 \int dE_{kin} E_{kin} \exp(-E_{kin}/k_B T) \sigma^{(n)}(if;E_{kin}) \Biggr) \\ \gamma + i\delta = n_b \overline{v} \Biggl(\frac{1}{k_B T} \Biggr)^2 \int dE_{kin} E_{kin} \exp(-E_{kin}/k_B T) \sigma^{(n)}(if;E_{kin}) \Biggr)$$

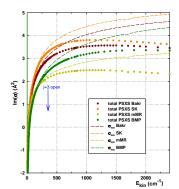
3. Raman isotropic Q(1) line of the fundamental band of D₂



Contributions to the cross-sections

widths: total inelastic elastic dephasing vibrational dephasing (ΔV_{iso})

> shifts: vibrational dephasing (ΔV_{iso})

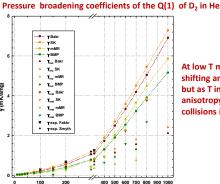




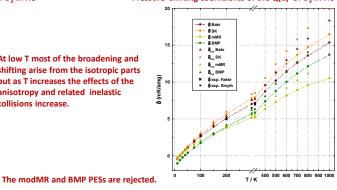
[2] P. Muchnick and A. Russek, J. Chem. Phys. 100, 4336 (1994)

[3] A.I. Boothroyd, P.G. Martin, M.R. Peterson, J. Chem. Phys. 119, 3187 (2003) [4] B.W. Bakr, D.G.A. Smith, K. Patkowski, J. Chem. Phys. 139, 144305 (2013)

Pressure shifting coefficients of the Q(1) of D₂ in He

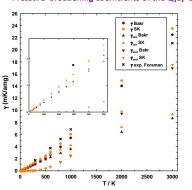


At low T most of the broadening and shifting arise from the isotropic parts but as T increases the effects of the anisotropy and related inelastic collisions increase.

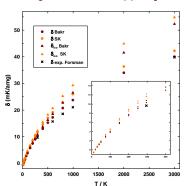


4. Raman isotropic Q(1) line of the fundamental band of H2 in He

Pressure broadening coefficients of the Q(1) of H2 in He

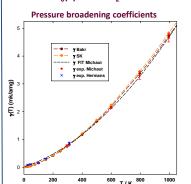


The SK PES gives better agreement with the experimental⁷ PB coefficients while the Bakr PES provides better agreement for the PS coefficients.



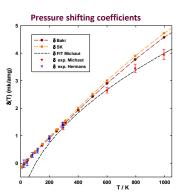
Pressure shifting coefficients of the Q(1) of H2 in He

5. Stokes S₀(1) line of H₂ in He



In this case it is important to take into account the centrifugal distortion especially for the shifts.

The SK and Bakr et al PESs lead to very close results in good agreement with the experimental values8,9 over the full range of T.



Work in progress proves that the relative differences between PB coefficients derived with the SK and Bakr PESs are smaller than relative differences between coefficients taking into account inhomogeneous broadening or not. Including the inhomogeneous mechanism leads to much better agreement with observed widths.



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[7] J.W. Forsman, J. Bonamy, D. Robert, J.-Ph. Berger, R. Saint-Loup, H. Berger, Phys. Rev. A. 52, 2652 (1995) [8] P.W. Hermans, A. Van Die, H.F.P. Knaap, J.J.M. Beenakker, Physica A 122, 233 (1985) [9] X. Michaut, R. Saint-Loup, H. Berger, M.-L. Dubernet, J. Bonamy, D. Robert, J. Chem. Phys. 109, 951 (1998)

