



HAL
open science

Tailoring the morphology of photowritten buried waveguides by helical trajectory in As₂S₃ glass

Ophélie Caulier, Pascal Masselin, Eugène Bychkov, David Le Coq

► **To cite this version:**

Ophélie Caulier, Pascal Masselin, Eugène Bychkov, David Le Coq. Tailoring the morphology of photowritten buried waveguides by helical trajectory in As₂S₃ glass. *Optics and Laser Technology*, 2016, 78, pp.56-61. 10.1016/j.optlastec.2015.09.032 . hal-01220883

HAL Id: hal-01220883

<https://univ-rennes.hal.science/hal-01220883>

Submitted on 26 Nov 2015

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Tailoring the morphology of photowritten buried waveguides by helical trajectory in As₂S₃ glass

Ophélie Caulier¹, Pascal Masselin¹, Eugène Bychkov¹, David Le Coq^{2,*}

¹Laboratoire de Physico Chimie de l'Atmosphère, EA 4493 Université du Littoral Côte d'Opale, 59140 Dunkerque, France

²Institut des Sciences Chimiques de Rennes, Eq. Verres et Céramiques, UMR CNRS 6226, Université de Rennes 1, 35042 Rennes Cedex, France

* Corresponding author: david.lecoq@univ-rennes1.fr

Phone/Fax : +33 (0)223235423/5611

ABSTRACT

This paper deals with the control of the refractive index variation (Δn) profile induced by femtosecond laser irradiation in the bulk of As₂S₃ glass. The writing technique consists of an original laser inscription based on a helical translation of the sample parallel to the laser beam. Instead of inscribing the core of the waveguide as usually performed, the laser beam is used to induce a negative refractive index variation and consequently to write the cladding of the waveguide. However, it should be noticed that the matter displacement resulting from local heating leads to a densification at the center of the helix, which induces a positive Δn relatively to the matrix. Therefore, the structure of Δn is complex, being composed of a positive core surrounded by a negative cladding. The influence of different parameters on the Δn profiles such as the sample translation velocity, the pitch and the radius of the helical displacement, and the pulse energy are analyzed. This study demonstrates that both Δn and its diameter can be varied in a wide range of values and picked independently, allowing the design of single or multimode buried infrared waveguide.

Keywords : Chalcogenide glasses - Femtosecond laser - Laser materials processing – Photoinscription – Infrared waveguide

I. Introduction

Focused ultrashort-pulsed laser pulses induce refractive index modifications in glass. This constitutes a real opportunity to generate three-dimensional optical components [1,2,3]. The advantages of this technique are numerous. Among them, the possibility to inscribe a waveguiding structure in a single step, without chemical process, is of particular interest because of its rapidity. Moreover, the laser writing allows the inscription of complex structures. For example, the realization of Mach-Zender interferometers [4] or of double waveguide couplers [5] have already been demonstrated. Due to their specific optical [6,7] or electrical properties [8,9] chalcogenide glasses (ChG) are an integral part of a scientific research domain [10], especially if the spectral region of interest is the infrared range.

Complementary common characteristics of ChG are their high optical nonlinearities [11] and their high photo-sensitivity [12] which make the inscription of refractive index variation (Δn) easy. Consequently, they have been the subject of many investigations in the laser photoinscription field [13,14,15,16,17,18].

The understanding of the fundamental and physical processes induced by the laser pulse is the basis of controlling the refractive index changes and consequently of realizing new integrated

photonic devices. According to the composition of the chalcogenide glasses, many phenomena can occur during the irradiation. In the case of As_2S_3 , at high repetition rate of femtosecond pulsed laser, the involved thermal effects have been shown to produce either low positive or negative Δn amplitude depending of the pulse energy when longitudinal writing geometry is used [19]. However, an original writing technique based on a helical translation of the sample throughout the focal point of the laser beam has been recently successfully demonstrated to obtain large positive Δn allowing infrared waveguiding [20]. This very promising technique opens a new route to the fabrication of buried infrared waveguides. The homogeneity of the inscriptions has also been achieved over more than 25 mm by solving the drawbacks due to the longitudinal trajectory [20].

In this paper, an overview of different parameters involved in the helical writing process, which directly influence the morphology of Δn in the glass As_2S_3 , is given. We highlight the potential of this writing process to inscribe a buried waveguide with large ranges of both Δn amplitude and diameter considered independently.

II. Experimental section

II.1. Sample preparation

As_2S_3 glass sample was synthesized from high purity As (5N) and S (4N). These elements were separately distilled before weighting in stoichiometric proportions in a glove box under Ar atmosphere and introduced together in a silica set-up before evacuation (10^{-5} Pa). A final distillation was performed down to a melting ampoule which was isolated by sealing and placed in a rocking furnace for 24 hours at 750°C . The quench was operated in water at room temperature. The sample was finally annealed at the glass transition temperature T_g for 4 hours before being slowly cooled down to room temperature in order to relax internal stress. The rod of 10 mm diameter was then cut and polished down to optical quality into a parallelepiped of $7 \times 7 \times 25$ mm³.

II.2. Principle of the helical inscription

A Ti:sapphire ultrafast laser (Coherent Mira) was employed at a high repetition rate of 76 MHz and a nominal output pulse duration of 200 fs with a central wavelength of 800 nm. The energy per pulse was adjusted by using two crossed polarizers, the latter one being in a fixed position while the first one was rotated for adjustment. The beam was focused into the sample with a $f = 50$ mm lens ($\text{NA} = 0.03$) and the $1/e^2$ intensity radius at the focal point was estimated to be 20 μm in air. The sample was fixed on three translation stages: (i) two Newport XMS-160, for circular movement in the (x,y) plane, and (ii) one Newport GTS-150, for linear translation along the beam propagation axis (z). The rotation velocity was fixed to $\omega = 75$ rad/s meaning that a complete circle in the (x,y) plane was done in $2\pi/\omega \approx 0.084$ s. A scheme of the helical writing is depicted in Fig. 1(a), and an outline the corresponding Δn modification is given in Fig. 1(b) considering that the experimental conditions lead to an inscription with a negative Δn .

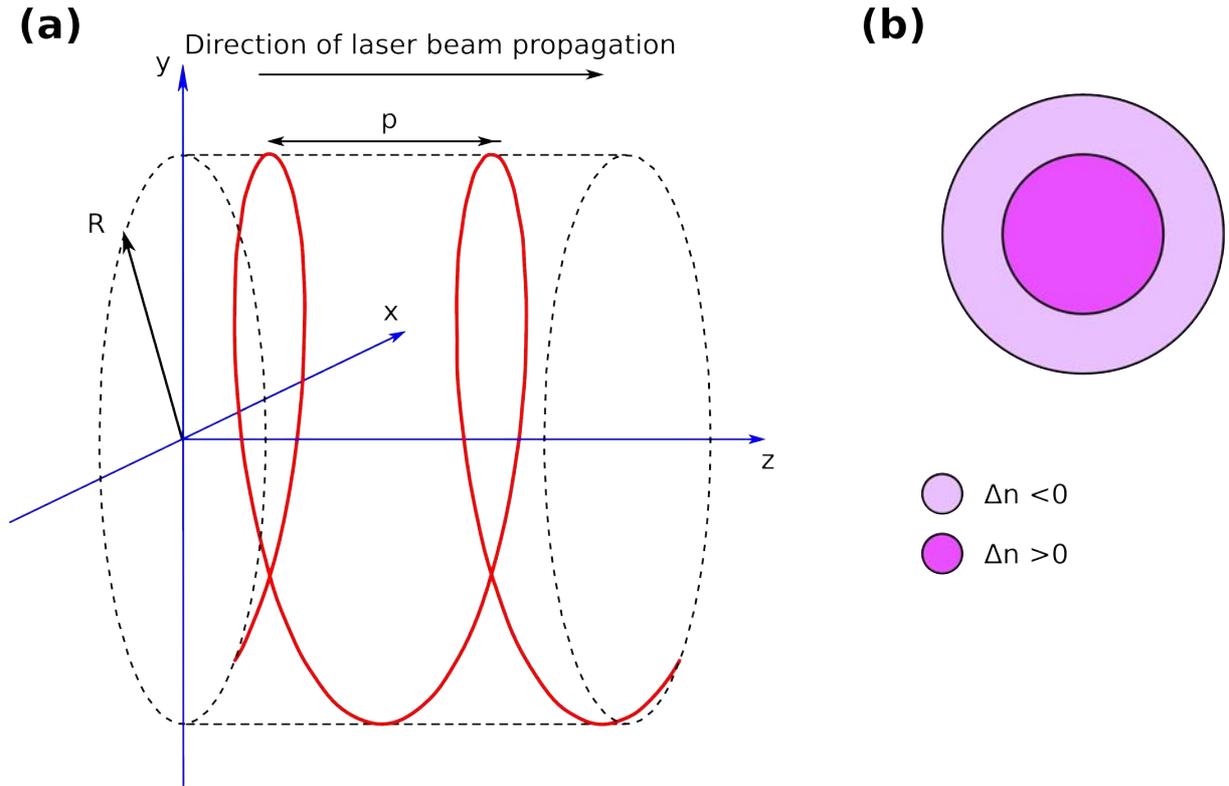


Figure 1: (a) Scheme of the helical writing in the bulk of the glass As_2S_3 . R is the radius of the helix and p is the helix pitch. (b) (x,y) plane image of the Δn modification after the helical writing.

II.3. Δn profile reconstruction

The first step in the reconstruction of the Δn profile is done by quantitative phase imaging [21] according to the procedure described in Ref [22]. All the images are recorded with a standard microscope equipped with a Hamamatsu OrcaR² camera. As the phase profile is axially symmetric, an Abel inversion, following a Nestor-Olsen algorithm can be applied to the phase image to obtain spatially resolved $\Delta n(r)$ [23,24] (r is the transverse direction). All measurements were carried out by using a wavelength around 630 nm.

The onset of the Fig. 2(a) is a three-dimensional transverse reconstruction of the Δn profile. The helical writing technique allows obtaining a perfect circular structure in term of Δn modification. Fig. 2(a) exhibits a typical longitudinal Δn reconstruction. In this example, the power P , the helix radius R , the helix pitch p , and the z -translation velocity were 50 mW, 9 μm , 25 μm , and 1000 $\mu\text{m/s}$, respectively. The longitudinal structure presents a positive Δn at the center of the modified Δn area surrounded by a negative Δn ring corresponding to the irradiated zone part. The Fig. 2(b) shows from top to bottom, a photograph of the obtained waveguide, the phase image recorded according to the procedure describes above, and the corresponding Δn . This last image emphasizes the homogeneity of the inscription.

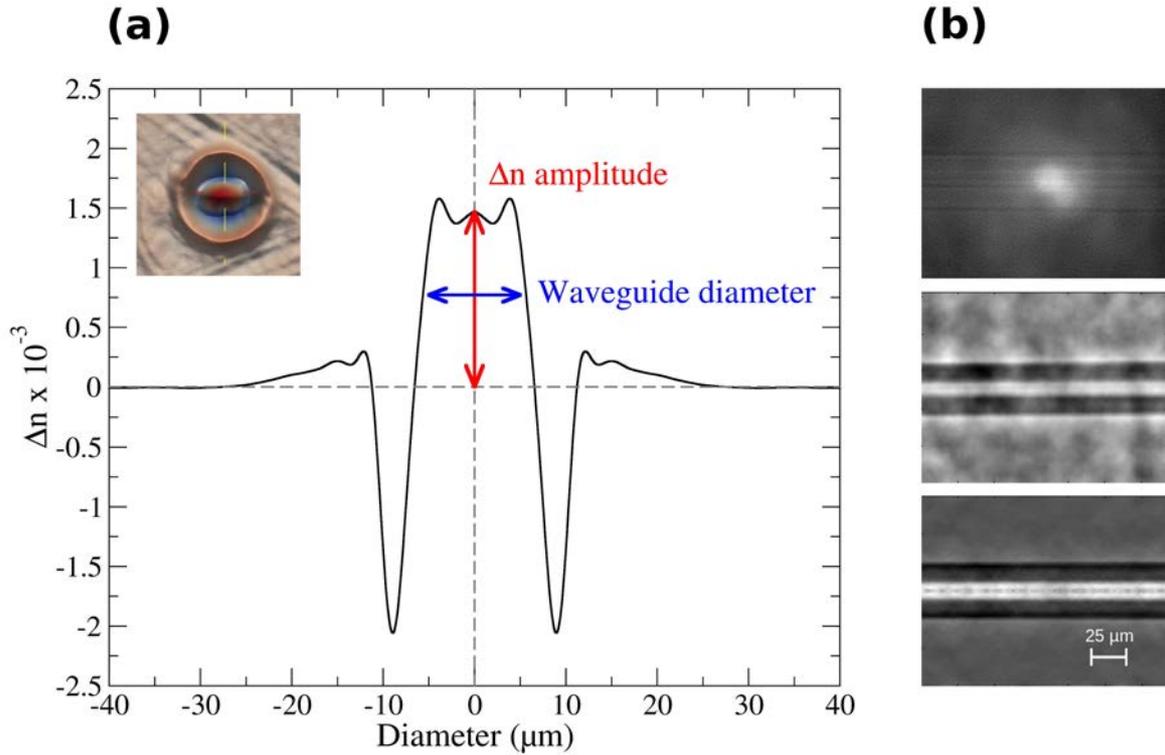


Figure 2: (a) Δn type profile obtained by helical writing by use of femtosecond laser beam. The onset represents the 3D transverse Δn profile of the helical writing. The following specific parameters were applied in the present case: $v_z = 1000 \mu\text{m/s}$, $R = 9 \mu\text{m}$, $p = 25 \mu\text{m}$, and $P = 50 \text{ mW}$. (b) From the top: microscope photo of an inscription, image of the phase matrix, and image of the index matrix ($v_z = 250 \mu\text{m/s}$, $R = 9 \mu\text{m}$, $p = 25 \mu\text{m}$, and $P = 50 \text{ mW}$)

III. Results

The profile of the Δn modification generated from a helical inscription is a function of different parameters that include the helix pitch, the z -translation velocity, the helix radius or the delivered power.

III.1. Influence of the helix pitch

The first investigated parameter is the length of the helix pitch (p). The other parameters were set as follows: the z -translation velocity equal to $250 \mu\text{m/s}$, the helix radius of $9 \mu\text{m}$, and a average power of 50 mW . The Fig. 3 depicts the Δn profiles obtained for four different values of p : $5 \mu\text{m}$, $25 \mu\text{m}$, $100 \mu\text{m}$, and $200 \mu\text{m}$. For $p = 5 \mu\text{m}$, no significant Δn modification is observed. The presence of a positive Δn surrounded by a negative ring is visible for other values of p . Other experiments that we have carried out, have also shown that if $p < 25 \mu\text{m}$, the Δn amplitude is very low and inhomogeneous. From $p = 25 \mu\text{m}$ up to $p = 150 \mu\text{m}$, Δn is homogeneous and its amplitude continuously decreases. Moreover, the Fig. 2(b) highlights the absence of spatial period corresponding to the pitch ($p = 25 \mu\text{m}$ in the represented

example). If $p > 200 \mu\text{m}$, the center of the waveguide begins to exhibit a decrease of the Δn until appearances of unscribed zones and inhomogeneities. Another point has to be emphasized: for homogeneous inscriptions, the waveguide diameter, as defined in Fig. 2(a) remains almost constant whatever the helix pitch.

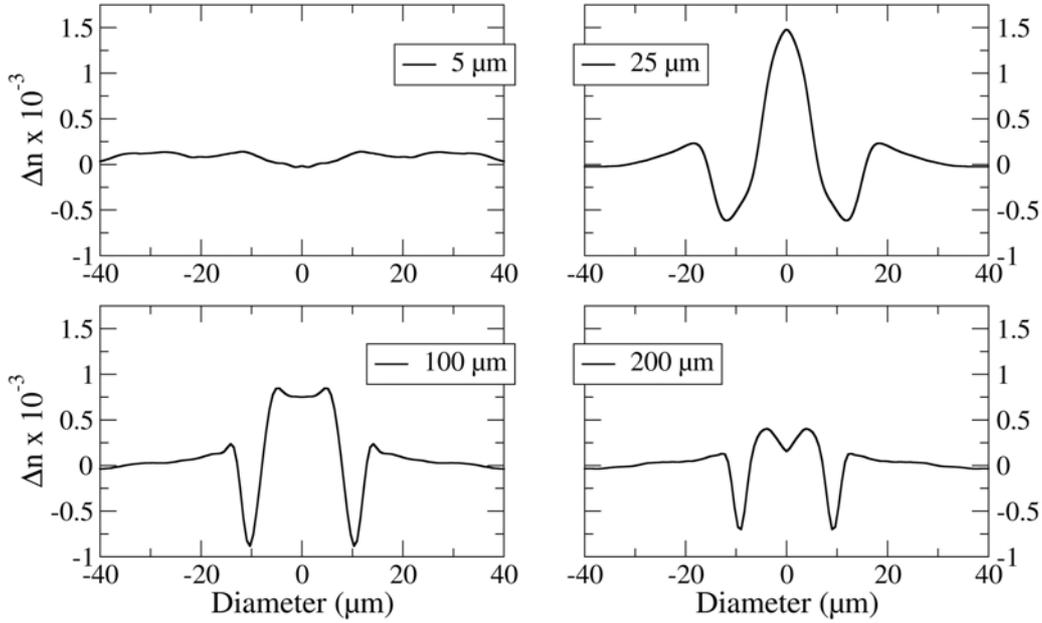


Figure 3: Examples of Δn profiles following the length of the helix pitch (p). Each inscription was performed using a z-translation velocity equal to $250 \mu\text{m/s}$, a helix radius of $9 \mu\text{m}$, and a power of 50 mW . Δn profile for 4 values of p are plotted: $p = 5; 25; 100; 200 \mu\text{m}$.

III.2. Influence of the translation speed

The second studied parameter is the z-translation velocity v_z , which is defined as follows:

$$\|s_{tz}\|^2 = \|v_z\|^2 + \|R \cdot \omega\|^2$$

where s_{tz} is the scalar velocity and $R \cdot \omega$ the rotation speed in the transverse plane.

In Fig. 4, the Δn profile is given for a few v_z values, between $v_z = 50 \mu\text{m/s}$ up to $v_z = 1500 \mu\text{m/s}$. The other parameters, p , R , and P were set at $25 \mu\text{m}$, $12 \mu\text{m}$, and 50 mW , respectively. At first, we can see that for $v_z = 50 \mu\text{m/s}$, the Δn modification is very low and no waveguide structure is discernable. For $v_z = 250 \mu\text{m/s}$, a positive Δn surrounded by a negative part is observed. Next, if the z-translation velocity is increased (typically up to $1000 \mu\text{m/s}$), the amplitude of Δn is reduced while the waveguide diameter is increased. At the same time we can also mention that the homogeneity of the inscription is also better. From $v_z = 1500 \mu\text{m/s}$, the trend is reversed since inhomogeneity appears. Consequently, v_z plays a key-role since low v_z allows high Δn amplitude to be obtained but the waveguide diameter and also its homogeneity require a higher v_z .

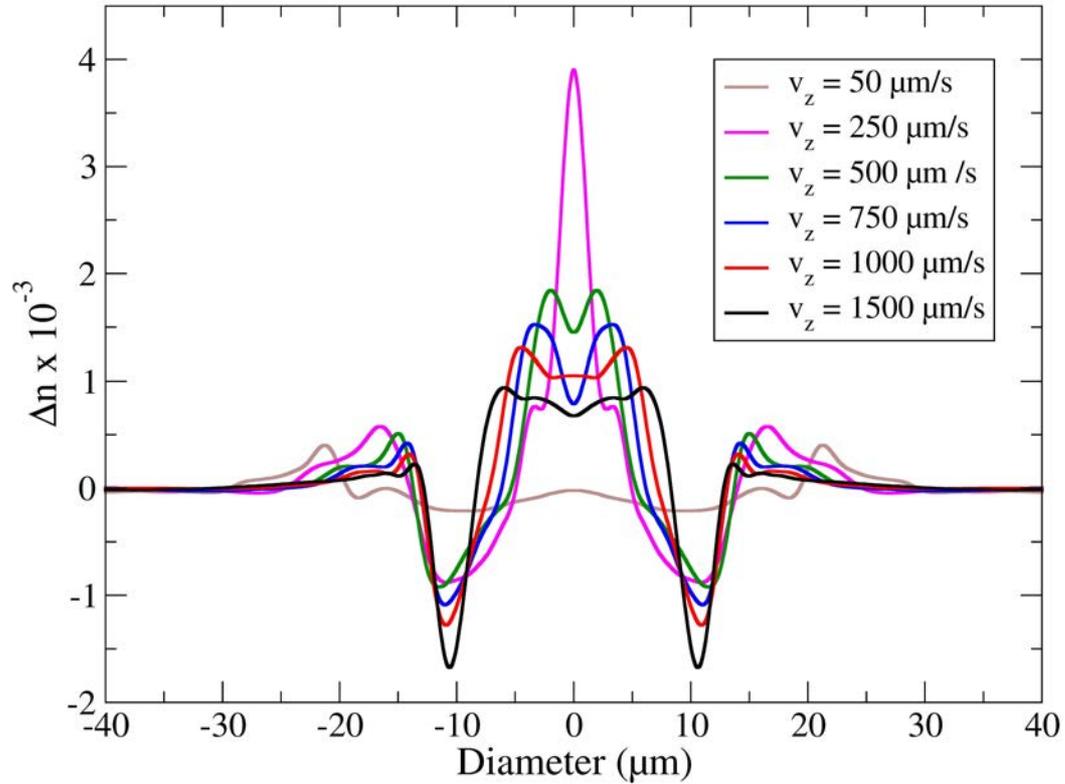


Figure 4: Examples of Δn profiles following the z -translation velocity (v_z) obtained with a helix pitch equal to $25 \mu\text{m}$ and a helix radius of $12 \mu\text{m}$, and a power of 50 mW . Δn profiles for 6 values of v_z are plotted.

III.3. Influence of the helix radius

The third investigated parameter is the helix radius R . The other parameters have been fixed as follows: $p = 25 \mu\text{m}$, $P = 50 \text{ mW}$; $v_z = 1000 \mu\text{m/s}$. The Δn profile evolution is displayed on Fig. 5 for $6 \mu\text{m} \leq R \leq 14 \mu\text{m}$. Actually, in these conditions we can observe that the higher is R , the larger is the diameter of the waveguide. For example, for $R = 6 \mu\text{m}$ the waveguide diameter is only $5 \mu\text{m}$, while for $R = 14 \mu\text{m}$ it is approximately $20 \mu\text{m}$. Nevertheless, at the same time a decrease of the maximum Δn amplitude is observed from 2.2×10^{-3} down to 1×10^{-3} . The amplitude of the negative Δn parts surrounding the positive Δn area remains constant. Consequently, it is noticed that this parameter has a relative strong influence on the waveguide diameter in regards to the amplitude of Δn .

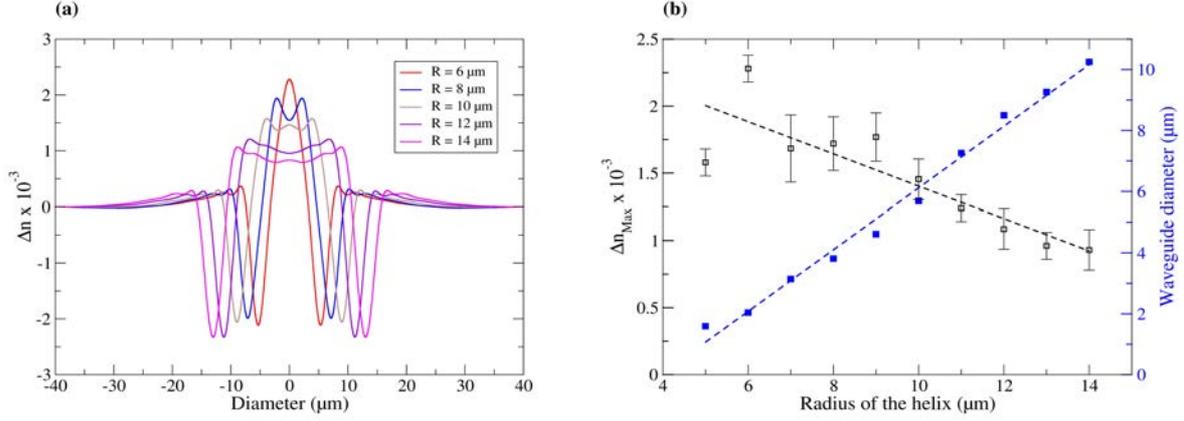


Figure 5: (a) Examples of Δn profiles following the radius of the helix (R) obtained with a helix pitch equal to 25 μm , a z translation velocity equal to 250 $\mu\text{m/s}$, and a power of 50 mW. Δn profiles for 5 values of R are plotted. (b) Evolution of both Δn and the waveguide diameter following the helix radius

III.4. Influence of the laser power

The last studied parameter is the average power P . In Fig. 6a, the Δn profile evolution following several P between 40 mW and 120 mW is plotted. Up to 100 mW, we clearly observe an increase of the Δn amplitude of the central peak following P , and from 100 mW up to 120 mW, it is almost constant just above 5×10^{-3} . Concerning the negative Δn parts, a slight increase of the negative amplitude is observed. Moreover, in Fig. 6b, the Δn amplitude is displayed more accurately by using additional values of P . As a result, three domains are well brought out: (i) below 40 mW, characterized by none observable Δn modification; (ii) between 40 mW and 90 mW characterized by a linear increase of the positive Δn amplitude; (iii) above 100 mW, where a limit in term of Δn amplitude seems to be achieved ($\approx 5 \times 10^{-3}$) and inhomogeneities also become noticeable.

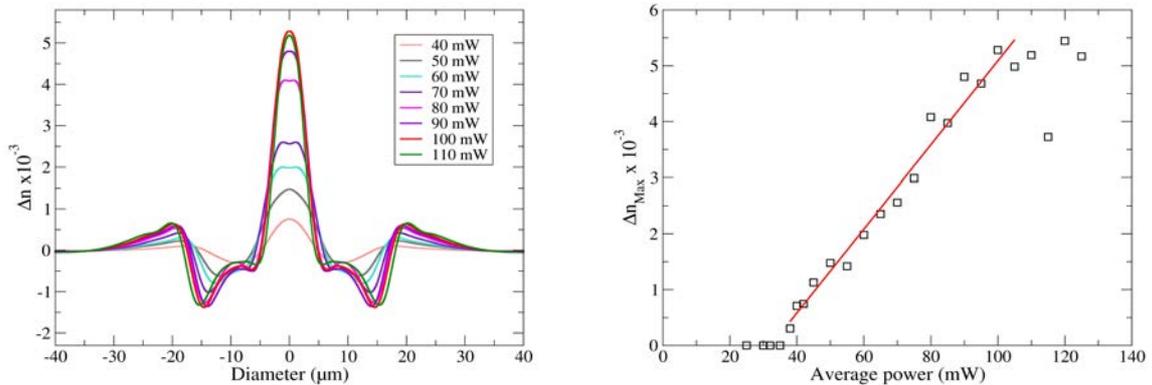


Figure 6: (a) Evolution of Δn profiles following the power (P) obtained with a helix pitch equal to 25 μm , a helix radius of 9 μm , and a z -translation velocity of 250 $\mu\text{m/s}$. Δn profiles for 8 values of P are plotted. (b) Evolution of the positive Δn amplitude as a function of the power

IV. Discussion

In order to understand all the results reported above, we have to consider the mechanism leading to Δn . We think that in all these experiments, the thermal effect is dominant and explains the observed behaviors.

The complex morphology of Δn can be understood according to this assumption. The consequence of the local heating is a pressure wave extending from the center of the focal volume toward the outside. When the glass temperature passes over T_g , the matter can move along the same direction. Consequently, the medium is less dense in the focal volume leading to a negative Δn . At the same time, we have to remember that the focal point follows a circular trajectory in the XY plan. Therefore, in all the point of this trajectory, the matter is pushed toward the center of the circle, resulting in a local densification providing that the temperature decreases below T_g before the matter has enough time to relax back into its original position. Consequently, the formation of a core of positive Δn results from the combination of the dynamic of matter displacement and temperature evolution, and the motion of the sample.

This later condition explains the observed behaviour. First, let's consider the dependence of Δn with the pitch of the helix. When the value of the pitch is low, the local temperature is higher because of the high spatial and temporal overlapping of the laser pulse [25]. Therefore, the time needed for the glass temperature to decrease below T_g is longer and the matter has enough time to relax back to its original location, resulting in a low Δn . Once the helix pitch is sufficiently important, the heat accumulation occurs again but the transition time to reach a temperature lower than T_g is shorter due to a less important overlap of the irradiated area and the matter is "quenched" in a densified state. Then, when the pitch increases further and the local temperature is lesser, the viscosity increases and the local displacement of the matter is less. Consequently, the Δn magnitude decreases when the pitch increases as observed. Also, if the pitch is too important, the inscription is not homogeneous anymore because of the non-uniformity of the heating and of the accumulation effect.

A similar explanation can be provided to interpret the dependence on the z-translation velocity. One can compare the Fig. 3(a) and Δn corresponding to $v_z = 50 \mu\text{m/s}$ in Fig. 4. These two curves correspond to the same situation when the local temperature is very high and the matter viscosity stays low enough so that the matter local displacement can relax back to its initial state. On the other hand, when the translation speed increases, the viscosity is higher and the pressure wave displaces the matter less efficiently. Consequently, the magnitude of the resulting Δn is less and its diameter is higher.

If we consider now the dependence with the laser power, a similar interpretation can also be given. In the domain of low power (less than 40 mW), any change of the index is detectable. This corresponds to a situation when the local temperature is not high enough to obtain a viscosity that allows efficient matter displacement. At high input power, Δn is not homogeneous anymore in reason of the too high temperature.

The decrease of the Δn magnitude with the radius of the helix can also be explained according to the thermal effect. In fact, as this radius is increased, the spatial overlapping of the heated zone by the laser pulse is less and once again the local temperature is less also.

Finally, we have to mention that a heating of the sample in an oven at a temperature close to T_g for a few hours erases all the Δn structures. Under such temperature condition, the viscosity is sufficiently low for the displaced matter to return to its original position and thus Δn to vanish.

It is worthy of note that a permanent Δn can also be obtained when the laser beam experiences strong self-focusing that enhances the multiphoton ionization probability [26]. However, this behavior is obtained only above a critical power that can be easily calculated [27]: under our experimental conditions, such calculation gives a value of 87 mW which is much higher than the threshold value observed in Fig. 4(b). Therefore, we can conclude that the influence of this mechanism does not play a significant role in our experiments.

V. Conclusion

In this paper, several parameters of the helical writing process have been investigated in order to evaluate their relative influence on the index profile generated in the As_2S_3 glass. The helix radius allows controlling the waveguide diameter from typically 3 to 20 μm under our experimental conditions. As for the mean power of the laser beam, it may be used to adjust the magnitude of the positive Δn of the waveguide between 0 and $5 \cdot 10^{-3}$. Thus, the ability of the helical writing process to inscribe infrared waveguide in the bulk of glass with well-defined characteristics has been clearly demonstrated. This technique offers a new possibility to fabricate monomode buried waveguide in glasses notably for the infrared range with an accurate control of both the Δn magnitude and the waveguide diameter.

Acknowledgements

Dr Ophélie Caulier is very grateful to the Direction Générale de l'Armement (DGA) and the Syndicat Mixte de la Côte d'Opale (SMCO) for financial support.

References

- [1] K Davis, K. Miura, N. Sugimoto, K. Hirao, *Optics Letters* 1996;21(21):1729-1731
- [2] S. Nolte, M. Will, J. Burghoff, and A. Tuennermann, *Appl. Phys. A* (2003);77:109-111
- [3] R. R. Thomson, T. A. Birks, S. G. Leon-Saval, A. K. Kar, J. Bland-Hawthorn, *Optics Express* 2011;19(6):5698-5705
- [4] J.C Ng, C. Li, P.R. Herman, L. Qian, *Optics Express* 2012;20(16):17894-17903
- [5] M. Pospiech, M. Emons, A. Steinmann, G. Palmer, R. Osellame, N. Bellini, G. Cerullo, U. Morgner, *Optics Express* 2009;17(5):3555-3563
- [6] C. Conseil, J.-C. Bastien, C. Boussard-Plédel, X.-H. Zhang, P. Lucas, S. Dai, J. Lucas, B. Bureau, *Opt. Mater. Express* 2012;2:470-1477
- [7] “Chalcogenide glasses for infrared optics”, Eds A. R. Hilton, S. Kemp, The McGraw Hill Companies Inc (2010)
- [8] E. Bychkov, Yu. Tveryanovich, Yu. Vlasov, in: R. Fairman and B. Ushkov (Eds.) *Applications of chalcogenide glasses, Semiconductors and Semimetals Series*, Vol. 80, Elsevier, New York – London, 2004
- [9] A. Hayashi, K. Noi, A. Sakuda, M. Tatsumisago, *Nature Communications* 2012;3:856
- [10] *Chalcogenide glasses: preparation, properties and applications*”, Eds. J.L. Adam, X.H. Zhang, Woodhead Publishing 2014
- [11] C. Lin, L. Calvez, B. Bureau, H. Tao, M. Allix, Z. Hao, V. Seznec, X. Zhang, X. Zhao, *Phys. Chem. Chemical Phys.* 2010;12(15):3780-3787
- [12] B. J. Eggleton, B. Luther-Davies, and K. Richardson, *Nat. Photonics* 2011;5:141-148
- [13] I Blonskyi, V. Kadan, O. Shpotyuk, M. Iovu, P. Korenyuk, I. Dmitruk, *Applied Phys B* 2011;104(4):951-956
- [14] O.M. Efimov, L.B. Glebov, K.A. Richardson, E. Van Stryland, T. Cardinal, S.H. Park, M. Couzi, J.L. Brunéel, *Opt. Mater.* 2001;17:379-386
- [15] J. F. Viens, C. Meneghini, A. Villeneuve, T.V. Galstian, E.J. Knystautas, M.A. Duguay, K.A. Richardson, T. Cardinal, *J. Lightwave Technol.* 1999;17:1184-1191
- [16] D. Le Coq, P. Masselin, C. Przygodski, E. Bychkov, *J. Non-Cryst. Solids* 2009;355:1832-1835
- [17] P. Masselin, D. Le Coq, A. Cuisset, E. Bychkov, *Opt. Mater. Express* **2012**;2(12):1768-1775
- [18] C. D’Amico, G. Cheng, C. Maclair, J. Troles, L. Calvez, V. Nazabal, C. Caillaud, G. Martin, B. Arezki, E. Le Coarer, P. Kern, R. Stoian, *Optics Express* 2014;22(11):13091-13101
- [19] D. Le Coq, P. Masselin, C. Przygodski, E. Bychkov, *J. Non-Cryst. Solids* 2009;355:1832-1835
- [20] O. Caulier, D. Le Coq, E. Bychkov, P. Masselin, *Optic Letters* 2013;38(20):4212-4215
- [21] A. Roberts, E. Ampem-Lassen, A. Barty, K. A. Nugent, G. W. Baxter, N. M. Dragomir, and S. T. Huntington, *Optic Letters* 2002;27:2061-2063
- [22] A. Barty, K. A. Nugent, D. Paganin, and A. Roberts, *Optic Letters* 1998;23:817–819
- [23] O. H. Nestor and H. N. Olsen, *SIAM Review* 1960;2:200-207

-
- [24] E. Ampem-Lassen, S. T. Huntington, N. M. Dragomir, K. A. Nugent A. Roberts, *Optics Express* 2005;13:3277-3282
- [25] S. Eaton, H. Zhang, P. Herman, F. Yoshino, L. Shah, J. Bovatsek, A. Arai, *Optics Express* 2005;13:4708-4716
- [26] P. Masselin, D. Le Coq, E. Bychkov, E. Lépine, C. Lin, L. Calvez, *Proc. SPIE* 2011;7993:79931B
- [27] J. Marburger, *Progress in Quantum Electronics* 1975;4:35-110