

# Interaction of Lanthanum with Boron and Carbon: Phase Diagram and Structural Chemistry

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### ▶ To cite this version:

Volodymyr Babizhetskyy, Arndt Simon, Josef Bauer. Interaction of Lanthanum with Boron and Carbon: Phase Diagram and Structural Chemistry. Chemical Monthly = Monatshefte für Chemie, 2014, 145 (6), pp.869-876. 10.1007/s00706-014-1172-2. hal-01024172

HAL Id: hal-01024172

https://hal.science/hal-01024172

Submitted on 15 Jul 2014

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## 1 Interaction of Lanthanum with Boron and Carbon: Phase Diagram and

### **2 Structural Chemistry**

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4 Volodymyr Babizhetskyy ● Arndt Simon ● Josef Bauer

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#### 8 Abstract

9 The isothermal section of the La–B–C phase diagram at 1270 K has been investigated by 10 means of X-ray, neutron powder diffraction, microstructure and EPMA analyses. Eight ternary 11 compounds were found, and for six of them the crystal structures have been established. The 12 phase with the structure type of La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> has a broad homogeneity range described by the 13 formula  $La_5(BC)_x$  (5.6 $\le x \le 8.8$ ). The lanthanum sesquicarbide  $La_2C_3$  exhibits an extended solid 14 solution in the ternary domain  $La_2C_{3-x}B_x$  (x=0.4). The boron substitution of carbon leads to the 15 decrease of the superconducting temperature from 13.4 K for La<sub>2</sub>C<sub>3</sub> to 10.0 K for La<sub>2</sub>C<sub>2.8</sub>B<sub>0.2</sub> and 16 from 5.6K for La<sub>2</sub>C<sub>2.7</sub> to 4.1K for La<sub>2</sub>C<sub>2.6</sub>B<sub>0.4</sub>. The compositions of two new compounds,

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19 **Keywords:** Carbides ● Solid state ● Crystal structure ● Phase diagrams

~La<sub>4</sub>B<sub>3</sub>C<sub>12</sub> and ~La<sub>4</sub>B<sub>5</sub>C<sub>18</sub> were found via WDX analysis.

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#### Introduction

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The search for new ternary boride carbides predicted from several structural models recently led to a systematic study of ternary phase diagrams, especially of carbides, borides and boride carbides. The structures of the ternary rare earth (RE) boride carbides display a variety of different arrangements with boron carbon substructures. They extend from zero-dimensional units to chains and two-dimensional nets embedded in the metal atom sublattices as well as interconnected boron icosahedra [1, 2]. The substructures can be divided into three groups. In the first group the finite (0-D) quasi-molecular entities fill voids of the metal atom matrix and can have different lengths ranging from 2 to 13 non-metal atoms. Stretched units of different sizes as well as isolated C atoms can coexist. In the second category, the non-metal atoms form infinite one-dimensional planar or nearly planar ribbons (BC) $_{\infty}$  of zigzag chains of boron atoms to which carbon atoms are attached. In the third family, the boron and carbon atoms form infinite planar (2-D) nets which alternate with sheets of metal atoms.

The phase relations in the general Ln-B-C systems were not studied in sufficient detail so far. The solid-state phase equilibria in the ternary phase diagrams with Ln = Y, Eu, Gd, Ho were reported [3 - 6]. The crystal structures of the known La-containing compounds were determined from X-ray single crystal data. Two essential structural features of the model for LaB<sub>2</sub>C<sub>2</sub> solved in space group  $P\bar{4}$  2c presented in [7] remain uncertain. They address the questions, firstly, whether B and C atoms within the eight-membered rings are alternating or arranged in pairs and, secondly, whether B and C atoms are stacked identically or alternatingly along c. The reexamination of the LaB<sub>2</sub>C<sub>2</sub> structure using neutron powder diffraction has been performed [8]. The underlying coloring problem in extended networks and the correlation between symmetry and electronic stability has been thoroughly treated by Burdett et al. [9], and on the basis of Extended Hückel calculations they reached the conclusion that a coloring scheme with only heteroatomic B-C bonds in contrast to [7] should represent the stable pattern in the LaB<sub>2</sub>C<sub>2</sub> structure family. Refinements in P4/mbm are in agreement with the predicted structure

characterized by only heteroatomic B-C bonds, however, in distorted 2-D squares and octagons, in contrast to the structural model proposed in [8].

From the first category with 0-D quasi-molecular entities, the phase La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> containing CB<sub>3</sub> units and isolated carbon atoms has recently been described in detail [10, 11]. The atom displacements in the structure are indeed due to real structure effects and are not a consequence of experimental artefacts (e. g. absorption) or inappropriate structure models (e. g. wrong choice of space group, superstructures, or twinning). Tetragonal rare earth boride carbides with structure types derived from ~La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> and Sc<sub>3</sub>C<sub>4</sub>, respectively, can intergrow coherently. Slabs of both types with various thicknesses from one unit cell up to macroscopic domains may occur, sharing a common square net of metal atoms [12]. In the crystal structures additional La<sub>x</sub>B<sub>y</sub>C<sub>7</sub> compounds were found with B/C chains of different lengths. So, the structure of LaBC contains non-linear B<sub>5</sub>C<sub>5</sub> chains [13]. The structure of La<sub>10</sub>B<sub>9</sub>C<sub>12</sub> contains slightly corrugated twodimensional metal atom square nets forming two types of chains with 26 and 18 La atoms, respectively, wherein  $B_4C_4$  and  $B_5C_8$  units are located [14]. The structure of  $La_5B_4C_{5-x}$  (x = 0.15) is composed of slightly corrugated two-dimensional metal atom square nets hosting the finite boron-carbon units B<sub>4</sub>C<sub>4</sub>, B<sub>3</sub>C<sub>3</sub>, BC<sub>2</sub> and isolated carbon atoms. The positions of the carbon atoms in the unit BC<sub>2</sub> are not fully occupied [15]. Longer oligomeric anions are observed when the boron-carbon ratio increases. Eleven-membered chains of different B/C ratio, B<sub>4</sub>C<sub>7</sub> and B<sub>5</sub>C<sub>6</sub>, exist in  $La_{15}B_{14}C_{19}$  [16].

In order to pursue our systematic research on the  $La_xB_yC_z$  phases, we focused on the solid-state phase equilibria in the ternary La-B-C phase diagram. Hence, the presentation of the isothermal section of the ternary La-B-C phase diagram at 1270 K, crystallographic data of the ternary compounds and the structure and determination of the homogeneity range of  $\sim La_5B_2C_6$  are subjects of this work.

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Phase equilibria in the the La-B-C ternary system

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5 The isothermal section of the La-B-C phase diagram at 1270K is presented in Fig. 1. In 6 agreement with Schlesinger et al. [17], two binary compounds, LaB<sub>4</sub> and LaB<sub>6</sub> have been found 7 in the La–B binary system at 1270K. The existence of two binary compounds, La<sub>2</sub>C<sub>3</sub> and LaC<sub>2</sub>, 8 have been confirmed according to the data of Gschneidner et al. [18]. The superconducting 9 properties of the rare earth sesquicarbide and homogeneity region of La<sub>2</sub>C<sub>3-8</sub> (T<sub>c</sub>=13.4 K for 10 La<sub>2</sub>C<sub>3</sub> and 5.6K for La<sub>2</sub>C<sub>2.7</sub>) were recently studied [19, 20]. The standardized cell parameters at 293K (Germanium 99.9999%,  $a_{Ge}$ =5.657905 Å, served as internal standard) for La<sub>2</sub>C<sub>3- $\delta$ </sub> changes from a=8.8149(5) for La<sub>2</sub>C<sub>3</sub> to 8.8079(5)Å for La<sub>2</sub>C<sub>2.7</sub>. The results of EPMA analysis show that 12 this phase is extended in a ternary domain. The influence of the boron addition at the La<sub>2</sub>C<sub>3- $\delta$ </sub> was 13 14 also observed by magnetic measurements of crushed alloys buttons. So, the boron substitution in  $\text{La}_2\text{C}_3$  leads to the composition  $\text{La}_2\text{C}_{2.8(1)}\text{B}_{0.2(1)}$  (a=8.8136(4) Å) and decreasing  $T_c$  from 13.4 K to 15 16 10.0 K. For La<sub>2</sub>C<sub>2.7</sub> a limit of the homogeneity was observed at La<sub>2</sub>C<sub>2.6</sub>B<sub>0.4(1)</sub> (a=8.8218(6) Å), the phase becoming superconducting at 4.1 K in contrast to  $T_c = 5.6$  of La<sub>2</sub>C<sub>2.7</sub> [19]. The results 17 will be published separately. A very small range of homogeneity of LaC<sub>2</sub> was detected by 18 19 neutron and physical properties experiments [21]. No extension of a ternary domain has been 20 detected for the binary compound "B<sub>4</sub>C" [22, 23]. The phase boundary of the La-rich liquid was not defined within the ternary system. The micrograph analysis was unsuccessful because of 22 high sensitivity of samples against moisture. The form of the solid phases arrangements below 23 1190 K is sketched in Fig. 1 by dashed lines corresponding to the literature data of binary 24 systems presented in [17, 18].

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< Fig. 1 >26

1	Metallographic investigation, X-ray powder diffraction and EPMA analyses revealed the
2	new compound ~La <sub>7</sub> B <sub>9</sub> C <sub>34</sub> to be in equilibrium with LaB <sub>6</sub> (Fig.2 a, b). Fig. 2 c, d shows the
3	three-phase region of the La-B-C ternary system between LaB <sub>4</sub> , LaBC and La confirmed by X-
4	ray powder diffraction and EPMA analyses.
5	The crystallographic characteristics of the ternaries in the La-B-C system are listed in
6	Table 1. It is worth noting that after annealing at 1270 K the phase La <sub>15</sub> B <sub>14</sub> C <sub>19</sub> [16] was not
7	found. It exists only in arc-cast samples. The compositions of two new compounds $\sim La_4B_3C_{12}$
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10	< Fig. 2 >
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12	and ~La <sub>4</sub> B <sub>5</sub> C <sub>18</sub> were found using WDX analysis. They show up in the phase diagram according
13	to Fig. 1.
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15	Solid solution $La_5(BC)_x$ (5.6 $\leq x \leq 8.8$ )
16 17	$RE_5B_2C_6$ was first discovered and labeled as "Gd <sub>2</sub> BC <sub>2</sub> " [5a]. The phase exhibits a large
18	homogeneity region. The X-ray powder diagrams of "RE2BC2" show a similarity with
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20	< Table 1 >
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22	corresponding rare earth dicarbides (CaC <sub>2</sub> type structure) with an observed a parameter slightly
23	smaller than in $REC_2$ and doubled $c$ parameter. Firstly, the crystal structure of $\sim La_5B_2C_6$ has
24	been determined in space group P4 from single crystal x-ray data [24] and the correct formula
25	and superconducting transition at 6.9 K was deduced. Later, the structure of ~La <sub>5</sub> B <sub>2</sub> C <sub>6</sub> [10] was

reinvestigated in space group P4/ncc (a=8.590, c=12.398 Å). Crystal structure and real

structure investigations on the intermittent superconductivity of  $\sim La_5B_2C_6$  indicated in [11]

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1	clearly show that superconductivity is not an intrinsic property of the phase, but must be ascribed
2	to precipitations of a B containing $La_2C_{3-x}$ phase in the C rich region and $\beta$ -La in the metal-rich
3	region. Obviously, the complete coverage of borocarbide crystallites by La can result in a total
4	magnetic shielding of the samples

6 < Fig. 3 >

The structural arrangement of  $\sim$ La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> compounds consists of a three-dimensional framework of rare-earth atoms resulting from the stacking of slightly corrugated two-dimensional squares, which lead to the formation of octahedral voids and distorted bicapped square antiprismatic cavities (Fig. 3 a-c). They are filled with isolated carbon atoms and twofold disordered CBCC units, respectively. The C1, C2/B2 atoms in  $\sim$ La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> form bent CBCC units with C/B-C distances of 1.32 Å indicative of double bond character and C/B - C/B distances of 1.65 Å. The C– C/B – C/B angle is 158.1°.

Unambiguous locations of C/B in Wyckoff position 16g in disordered CBCC units as well as occupation of light atoms are not always obvious from the X-ray measurements. Therefore supplementary powder neutron and synchrotron single crystal experiments were performed. The

20 < Table 2 >

23 < Table 3 >

results of the structure refinement are in good agreement with WDX analyses. Tables 2 and 3 present the results of Rietveld refinements of x-ray, neutron and synchrotron single crystal

investigations of  $\sim$ La<sub>5</sub>B<sub>2</sub>C<sub>6</sub>. The unit cell parameter c for  $\sim$ La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> varies significantly and ranges from 12.315(2) to 12.843(4) Å. However, the occupation of the large voids by only disordered CBCC units is improbable. The observation of large displacement parameters for the central C/B atoms of the CBCC units is in line with the disorder of occupation by either CBC or CBCC units for the large voids and single carbon atoms in octahedral voids. That is the reason for the large homogeneity region of this compound. The structure refinement clearly shows that for the unit cells with c/a < 1.474 the BC<sub>3</sub> unit does not fully occupy the La<sub>10</sub> voids (Fig. 3c). The boron content in [C=B-C=C]<sup>7-</sup> varies from 15(1) to 21(1) at % which leads to a decrease in charge as well as disorder [5]. The overall charge of the anionic part seems to be compensated by introduction of single carbon atoms (C3) in La octahedra. The octahedron volume increases from  $30.524\text{Å}^3$  for empty octahedra to  $32.146\text{Å}^3$  for fully occupied ones by C3 atoms (Fig. 3b). So, the homogeneity region of  $\sim$ La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> is defined as La<sub>5</sub>(BC)<sub>x</sub> ( $5.6 \le x \le 8.8$ ) and presented in Fig. 1.

14 Conclusion

The isothermal section of the La–B–C phase diagram at 1270 K has been investigated by means of X-ray, neutron powder diffraction, microstructure and EPMA analyses. In the system eight ternary compounds were found. For six of them, namely  $\sim$ La<sub>5</sub>B<sub>2</sub>C<sub>6</sub>, La<sub>5</sub>B<sub>4</sub>C<sub>5</sub>, LaBC, La<sub>10</sub>B<sub>9</sub>C<sub>12</sub>, La<sub>15</sub>B<sub>4</sub>C<sub>14</sub>, LaB<sub>2</sub>C<sub>2</sub>, the crystal structures have been established. The phase La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> has a wide homogeneity range described by the formula La<sub>5</sub>(BC)<sub>x</sub> ( $5.6 \le x \le 8.8$ ). For the boundary compounds of the La-B system no extension in the ternary domain was found. In contrast to binary borides the lanthanum sesquicarbide La<sub>2</sub>C<sub>3</sub> exists as solid solution La<sub>2</sub>C<sub>3-x</sub>B<sub>x</sub> (x=0.4). The carbon substitution by boron leads to a decrease in  $T_c$  from 13.4 K for La<sub>2</sub>C<sub>3</sub> to 10.0 K for La<sub>2</sub>C<sub>2.8</sub>B<sub>0.2</sub> and from 5.6K for La<sub>2</sub>C<sub>2.7</sub> to 4.1K for La<sub>2</sub>C<sub>2.6</sub>B<sub>0.4</sub>. The new compounds  $\sim$ La<sub>4</sub>B<sub>3</sub>C<sub>12</sub> and  $\sim$ La<sub>4</sub>B<sub>5</sub>C<sub>18</sub> show broad diffraction peaks. In contrast to the rare earths with small atomic radius lanthanum boride carbides form no compounds with infinite one-dimensional planar or nearly planar ribbons (BC)<sub>∞</sub>. The investigated structures contain finite (0-D) quasi-molecular B/C entities which fill voids of the lanthanum matrix and can have different lengths ranging from

2 to 13 non-metal atoms. Linear or bent units of different sizes as well as isolated C atoms can 1

Polycrystalline samples were prepared from commercially available pure elements:

2 coexist.

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Tables 1, 2.

#### 3 **Experimental**

4 Synthesis and analysis

lanthanum with a claimed purity of 99.99 at%, Alfa – Aesar, Johnson Matthey Company, 6 7 sublimed bulk pieces; crystalline boron powder, purity 99.99 at%, H. C. Starck, Germany; 8 graphite powder, purity 99.98 at%, Aldrich. Before use, the graphite and boron powders were outgased overnight at 1220 K,  $p < 10^{-5}$  mbar. Lanthanum ingots were filed to coarse powders 9 10 with beryllium bronze files (Dönges GmbH, Germany). Stoichiometric mixtures of the constituents were compacted in stainless steel dies. The pellets were arc-melted under purified 12 argon atmosphere [25] on a water-cooled copper hearth. The alloy buttons of 1g were turned 13 over and remelted three times to improve homogeneity. The samples were then wrapped in 14 molybdenum foil and annealed in silica tubes under argon for 1 month at 1270 K. Subsequent 15 heating for some samples just above the melting point was carried out in a high-frequency 16 furnace (TIG-10/300, Hüttinger, FRG) under argon atmosphere for several hours at different 17 temperatures. Finally, the samples were wrapped in molybdenum foils, annealed in evacuated 18 silica tubes for one month at 1270 K and subsequently quenched in cold water. Sample handling 19 was carried out under argon atmosphere in a glove box ( $P_{H2O}$ <0.1 ppm) or through the Schlenk 20 technique. X-ray powder diffraction patterns were obtained on a powder diffractometer STOE 22 STADI P with  $MoK_{\alpha 1}$  radiation, using capillaries sealed under dried argon to avoid hydrolysis. 23 The unit cell parameters for the investigated compounds La<sub>x</sub>B<sub>y</sub>C<sub>z</sub>, as well as the Rietveld 24 refinements for some samples containing the La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> phase were refined with the help of the

CSD program package or FULLPROF programs packages [26, 27]. The results are given in

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For neutron powder experiments the <sup>11</sup>B containing sample ~La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> of 10 g was sealed into a vanadium can under 1 bar of He and measured at 300 K using the GEM diffractometer at the ISIS laboratory. The structural parameters as well as the composition were gained from Rietveld refinements using the FULLPROF program. Figure 4 shows the neutron powder diffraction patterns of composition  $La_5B_2C_{6-x}$ . The refined parameters of the  $La_5B_2C_{6-x}$  (x = 1.4) phase include the lattice constants, the fractional coordinates of the La, B/C and C atoms, isotropic displacement parameters, (anisotropic in case of La). After refinements with the aforementioned parameters, in a last step, the C occupancy was refined and the convergence by varying the occupancy was checked in those cases where the other refined parameters are fixed or free. For both cases, a C deficiency in 4c and 16g sites is derived consistently. The results of the refinements at T=300 K are listed in Table 1. The single crystal diffraction data of La<sub>5</sub>B<sub>2-x</sub>C<sub>6-</sub> were collected at room temperature on a synchrotron facility with  $\lambda = 0.41328$ Å. The starting atomic parameters taken from [10] were refined with the program SHELX-97 [28] within the WinGX program package [29] (full matrix least-squares on  $F^2$ ) with anisotropic atomic displacements for La atoms. The atomic coordinates and atoms displacement parameters are listed in Tables 3 and 4.

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- 21 Microprobe analysis.
  - For metallographic inspection and electron probe microanalysis (EPMA) some alloys were embedded in Woods metal (Fluka Chemie, Switzerland) with a melting point of ca. 75°C. The samples were polished on a nylon cloth using chromium oxide (Bühler Isomet) with grain sizes 1-5 µm. Quantitative and qualitative composition analyses of the samples were performed by energy dispersive X-ray spectroscopy (EDX) and wavelength dispersive X-ray spectroscopy

(WDX) on a scanning electron microscope TESCAN 5130 MM with an Oxford Si-detector and with an Oxford INCA WAVE 700 detector. LaB<sub>4</sub> and LaB<sub>2</sub>C<sub>2</sub> standards were used to deduce the compositions of compounds. For the chemical microprobe, the polishing procedure had to be performed or repeated just before the measurements. Metallographic and EPMA analyses of the La-B-C ternary system reveal the compound LaBC [13] to be in phase equilibrium with La<sub>5</sub>B<sub>4</sub>C<sub>5</sub> [15], La<sub>10</sub>B<sub>9</sub>C<sub>12</sub> [14] and La<sub>5</sub>(BC)<sub>x</sub> (5.6  $\leq$  x  $\leq$  8.8) (La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> structure type) [11] and LaB<sub>4</sub> at 1270K. Metallographic and EPMA analyses are in good agreement with results from x-ray

powder diffraction (e. g. see Fig. 2). Chemical analysis of the sample La<sub>5</sub>B<sub>2.4</sub>C<sub>5.3</sub> (La<sub>5</sub>B<sub>2</sub>C<sub>6</sub>

structure) was performed in the Mikroanalytisches Labor Pascher in Remagen-Bandorf,

10 Germany. The results of the EPMA and chemical analyses are listed in Tabl. 2.

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12 Magnetical properties.

13 The magnetic properties were studied in the temperature interval 1.8 - 330 K by use of a MPMS

14 XL-7 SQUID magnetometer (Quantum Design, Inc.) in external fields up to 7 T. The crushed

alloys buttons of irregular shapes were used for the measurements.

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#### Acknowledgements

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The authors gratefully thank M. Babizhetska for the sample preparation, Dr. C. Hoch for X-ray intensity data collection, E. Brücher for the magnetization measurements, V. Duppel for EPMA analysis.

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- 1 Figure Captions
- 2 Fig. 1 Isothermal section of the La–B–C phase diagram at 1270 K. Dashed tie-lines correspond
- 3 to equilibria in the solid state below 1190 K and  $La_{15}B_{14}C_{19}$  observed in arc-cast samples.

- 5 **Fig. 2** Secondary electron image (a) and x-ray powder pattern (b) of the annealed bulk sample
- 6 with nominal atomic composition La/B/C = 14/32/54. Indicated (\*) reflections are from LaB<sub>6</sub>
- and other of ~La<sub>7</sub>B<sub>9</sub>C<sub>34</sub>. Backscattered electron image (c) and x-ray powder pattern (d) of the
- 8 annealed bulk sample with nominal atomic composition La/B/C = 40/45/15. Indicated (\*)
- 9 reflections are from LaB<sub>4</sub>, (•) LaBC and (+) La.

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- Fig. 3 Crystal structure of  $\sim$ La<sub>5</sub>B<sub>2</sub>C<sub>6</sub> (a). The chains of edge-sharing La<sub>6</sub>C octahedra are
- emphasized. The rare earth metal atoms environments of discrete carbon atoms (b) and
- 13 C-B/C-B/C-C (c) units are shown. Lanthanum atoms in (b) and (c) are represented by their
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Fig. 4 Comparison of observed and calculated neutron powder profiles for La<sub>5</sub>B<sub>2</sub>C<sub>4.6</sub>

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Compound	Space	Structure	Unit cell parameters (Å)			References
	group	type	а	b	С	
$La_5(BC)_x$	P4/ncc	$La_5B_2C_6$	8.584(1)-		12.315(2) -	[10], <sup>a)</sup>
$(5.6 \le x \le 8.8)$			8.598(1)		12.730(1)	
$La_5B_4C_{5-x}(x=0.15)$	$Pna2_1$	$Ce_5B_4C_5$	24.682(3)	8.612(1)	8.647(1 b)	[15], <sup>a)</sup>
LaBC	$P2_12_12_1$	LaBC	8.666(1)	8.707(1)	12.485(2)	$[13],^{a)}$
$La_{10}B_{9}C_{12}$	$P4_12_12$	$Ce_{10}B_9C_{12}$	8.6678(5)		25.689(3)	[14], <sup>a)</sup>
$La_{15}B_{14}C_{19}^{\ \ b)}$	$P2_{1}/c$	$La_{15}B_{14}C_{19}$	8.640(3)	8.636(4)	19.823(6)	[16], <sup>a)</sup>
				β=94.28(3)°		
LaB <sub>2</sub> C <sub>2</sub>	P4/mbm	LaB <sub>2</sub> C <sub>2</sub>	5.4050(2)		3.9623(2)	[8], <sup>a)</sup>
$\sim$ La <sub>4</sub> B <sub>3</sub> C <sub>12</sub>						
$\sim$ La <sub>4</sub> B <sub>5</sub> C <sub>18</sub>						

a) Powder data of the present work, b) observed in arc-cast samples.

**Table 2** Crystallographic and chemical composition data for alloys  $La_5(BC)_x$  (5.6 $\le x \le 8.8$ )

Composition	Lattice parameters	c/a	$V(\text{Å}^3)$	B,	(C1-C2/B2)	C3
Composition	(a, c  in  Å)	C/U	V(A)	at % <sup>1)</sup>	occupation	occupation
La <sub>5</sub> B <sub>2.8</sub> C <sub>6</sub>	8.622(1)	1.489	955.0(4)	21(1)	1.00	1.00
	12.843(4)					
La <sub>5</sub> B <sub>2.8</sub> C <sub>5.7</sub>	8.5837(1)	1.487	940.39(4)	20(1)	1.00	0.75(4)
	12.7632(2)					
La <sub>5</sub> B <sub>2.4</sub> C <sub>5.3</sub> <sup>2)</sup>	8.5807(1)	1.474	931.70(4)	19(1)	0.8(1)	0.50(4)
	12.6539(3)					
$La_5B_2C_{4.6}^{3)}$	8.5810(1)	1.458	921.44(4)	17(1)	0.8(1)	0.20(3)
	12.5140(2)					
La <sub>5</sub> B <sub>1.8</sub> C <sub>3.9</sub>	8.5746(4)	1.450	914.4(2)	16(1)	0.7(2)	0.10(5)
	12.437(1)					
$La_5B_{1.75}C_{3.8}^{4)}$	8.578(1)	1.438	908.3(3)	16(1)	0.7(2)	0.0
	12.342(2)					
La <sub>5</sub> B <sub>1.5</sub> C <sub>4.1</sub>	8.584(1)	1.434	907.5(4)	15(1)	0.7(2)	0.0
	12.315(2)					

WDX microprobe analyses
Chemical analysis in at %: La-38.1, B-19.0, C-42.9
Neutron powder diffraction data

<sup>4)</sup> Synchrotron single crystal diffraction data

Methode	Noutron	Cymahaathaan		
Methode	Neutron	Synchrothron		
	powder data <sup>b</sup> )	single-crystal data c)		
Composition	$La_5B_2C_{4.6}$	$La_5B_{1.75}C_{3.8}$		
La1 x	0.6502(3)	0.64998(3)		
y	0.0508(3)	0.05149(2)		
z	0.1083(3)	0.10687(2)		
$B_{ m eq}$	0.38(2)	0.86(1)		
1				
La2 z	0.1363(6)	0.13729(4)		
$B_{ m eq}$	0.76(3)	1.10(1)		
cq				
C1   x	0.6545(5)	0.6550(6)		
y	0.0457(6)	0.0448(6)		
z.	-0.09351(4)	-0.0890(4)		
$\ddot{B_{ m iso}}$	0.98(3)	1.03(1)		
Diso	0.70(3)	1.03(1)		
C2/B2 x	0.6269(5)	0.6292(6)		
y	0.0632(6)	0.0609(6)		
y Z.	-0.2010(3)	-0.2000(4)		
=	` '	` '		
$B_{ m iso}$	1.16(3)	1.46(1)		
C3 z	0.257(1)			
	0.357(1)			
$B_{ m iso}$	1.51(3)			
Occupancy:		0.7(0)		
C1	0.80(1)	0.7(2)		
C2/B2	0.30(1)/0.50(1)	0.3(2)/0.4(2)		
C3	0.20(3)	-		

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Table 4 Anisotropic displacement parameters for lanthanum atoms in ~La<sub>5</sub>B<sub>2</sub>C<sub>6</sub>

Atom	$B_{11}$	$B_{22}$	$B_{33}$	$B_{23}$	$B_{13}$	$B_{12}$
La1 a)	0.32(2)	0.40(3)	1.25(6)	0.002(6)	-0.0022(9)	-0.010(7)
La2 <sup>a)</sup>	0.30(2)	$B_{22}$	1.88(1)	0.000	0.000	0.000
La1 b)	0.501(6)	0.656(7)	1.68(1)	0.004(4)	-0.0059(6)	-0.021(5)
La2 <sup>b)</sup>	0.466(9)	$B_{22}$	2.28(1)	0.000	0.000	0.000

<sup>&</sup>lt;sup>a)</sup> La2 and C3 in 4c (1/4, 1/4, z) <sup>b)</sup>  $R_{\rm I}$ =0.067,  $R_{\rm F}$ =0.056,  $R_{\rm p}$ =0.10 ,  $R_{\rm exp}$ =0.14,  $\chi^2$ =1.7 <sup>c)</sup>  $R_{\rm I}$ =0.035,  $wR_{\rm I}$ =0.087 for 2997 reflections with I<2 $\sigma(I_{\rm o})$ .

a) neutron powder diffraction data b) synchrotron single crystal data

Fig. 1 Isothermal section of the La–B–C phase diagram at 1270 K. Dashed tie-lines correspond to equilibria in the solid state below 1190 K and  $La_{15}B_{14}C_{19}$  observed in arc-cast samples.

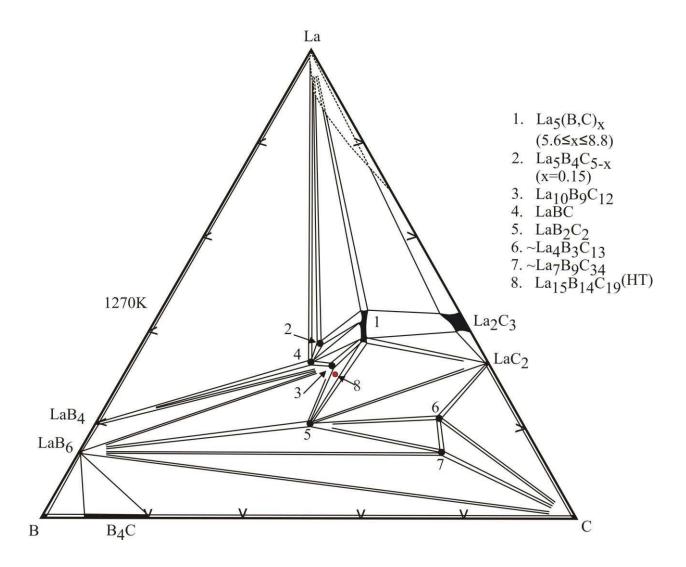


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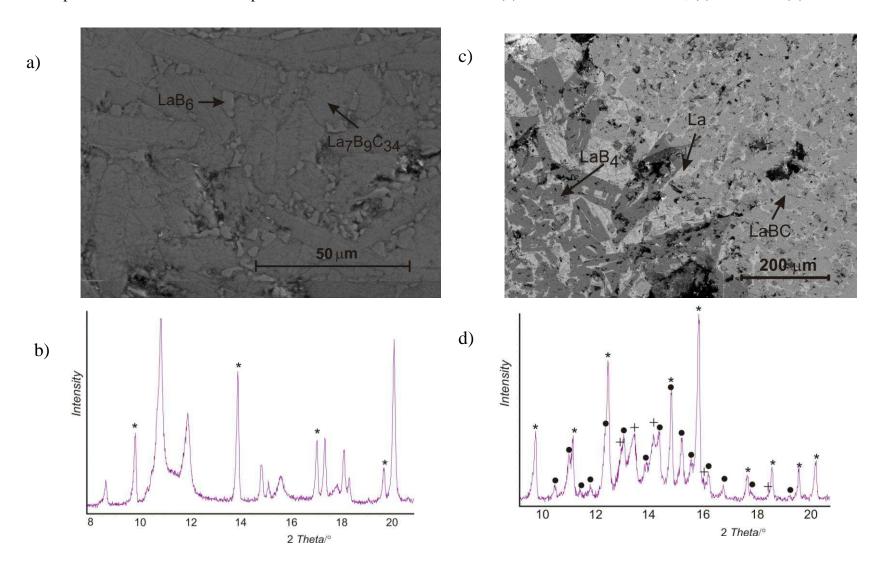


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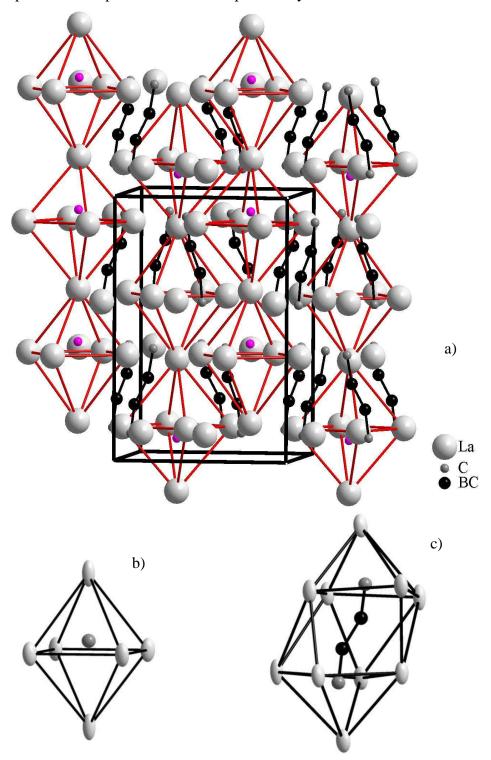
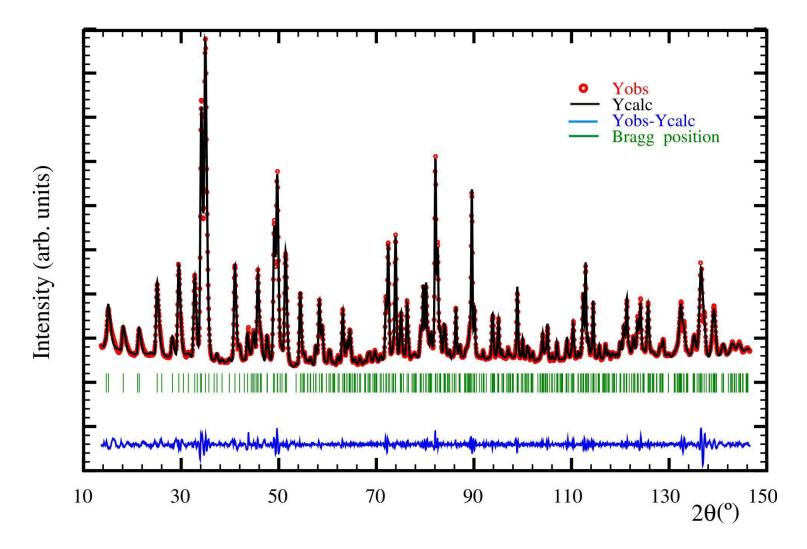


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# **Graphics for use in the Table of Contents**

