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# Vibrational spectroscopy and DFT calculations 1, 3-dibromo-2,4,6-trimethylbenzene. (Anharmonicity, Coupling and Methyl group tunneling)

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**Abstract-** The Raman, IR and INS spectra of 1,3-dibromo-2,4,6-trimethylbenzene (DBMH) were recorded in the 80-3200 cm<sup>-1</sup> range. The molecular conformation and vibrational spectra of DBMH were computed at the MPW1PW91/LANL2DZ level. Except for the methyl 2 environment, the agreement between the DFT calculations and the neutron diffraction structure is almost perfect (deviations < 0.01 Å for bond lengths, < 0.2° for angles). The frequencies of the internal modes of vibration were calculated with the harmonic and anharmonic approximations; the later method yields results that are in remarkable agreement with the spectroscopic data, resulting in a confident assignment of the vibrational bands. Thus, no scaling is necessary. The coupling, in phase or anti-phase, of the motions of symmetrical C-Br and C-Me bonds is highlighted. Our DFT calculations suggest that the torsion of methyl groups 4 and 6 is hindered in deep wells, whereas methyl group 2 is a quasi-free rotor. The failure of the calculations to determine the frequencies of the methyl torsional modes is explained as follows: DFT does not consider the methyl spins and assumes localization of the protons, whereas the methyl groups must be treated as quantum rotors.

**Keywords:** Raman spectroscopy; IR spectroscopy; Inelastic neutron scattering; DFT calculations; 1,3-dibromo-2,4,6-trimethylbenzene

## 1-Introduction

A tremendous amount of work has been devoted to the spectroscopy of polysubstituted benzenes. Among those relative to penta- and hexa-substituted compounds, two extensive papers were published in 1985: one paper compared eleven polymethyl-benzenes (pMeB) [1] and the second described six chloromethyl-benzenes  $C_6(CH_3,Cl)_6$  (ClMeB) [2]. In these studies, the experimental spectra were assigned using normal coordinate calculations to determine a common multi-parameter force field. Because the molecular conformations were not accurately known the authors assumed that the aromatic ring in all molecules was a regular hexagon, that the bond angles were  $120^\circ$ , and the bond lengths were 1.396 Å for the carbon-carbon aromatic bond, 1.509 Å for the methyl carbon-aromatic carbon bond, 1.083 Å for the aromatic carbon-hydrogen bond and 1.096 Å for the methyl-carbon hydrogen bond. For the pMeB species, the Raman and infra-red measurements were performed with samples in the gaseous phase at approximately 550 K. A 46-parameter valence force field had been proposed, and the average error between the 736 observed and calculated wave-numbers was  $6.4\text{ cm}^{-1}$ . This agreement suggests that the use of an excessive number of off-diagonal constants (interactions terms) allows us to compensate for the large errors done in the geometrical parameters. According to the authors in these previous studies: “there is a great stability of the internal methyl modes, in a limit of  $\pm 5\text{ cm}^{-1}$  they are located at:  $2930\text{ cm}^{-1}$  for the C-H stretching,  $1445\text{ cm}^{-1}$  for the asymmetric C-H bending,  $1380\text{ cm}^{-1}$  for the C-H symmetric bending, 1000 and  $1060\text{ cm}^{-1}$  for the C-H rocking”. For the “torsion mode” of the methyl groups, no experimental data were available. The authors assumed without proof that these groups are “free” rotors”, which is rather rare for a methyl group in solid state molecules. On the contrary this finding is well established for one of the three methyl groups of dibromomesitylene that was studied in this paper.

The parent ClMB molecules were studied as crystals from room temperature to liquid nitrogen temperature. It was also assumed that in these molecules the benzene rings are perfect hexagons. Because the  $CH_3$  group vibrations have been observed at similar frequencies to those in the pMeB molecules these groups were assumed to be not affected by the skeleton vibrations. Thus by considering the methyl group to be a point mass the authors were able to independently study the remaining 30 vibrations of the carbon-chlorine skeleton as those of a cyclic  $C_6(X,Cl)_6$  system. Then, the authors used an optimized multi-parameter valence force field to reproduce the 121 assigned frequencies in seven different molecules. The potential energy was determined using nine principal force constants corresponding to nine different types of deformations (bond lengths and angles). To incorporate the interactions between the substituent vibrations and the deformations of the benzene ring, 37 near neighbor interaction coefficients were considered. In fact, using these

new interaction parameters partially corrects the errors introduced when using a regular hexagon for the benzene ring in these asymmetrically substituted molecules. Finally the calculated frequencies reproduce the observed frequencies with an average error of only  $10.9\text{ cm}^{-1}$ .

In parallel to these empirical molecular mechanic calculations, quantum chemistry programs were established that were able to provide a molecular conformation in good agreement with the X-ray determination. In the late 1970s, the analytical second-derivative techniques were introduced at the Hartree-Fock (HF) level resulting in the sequence of the fundamental frequencies. However, when compared to the experimental frequencies, the computed HF frequencies were consistently 7-15 % too high. This overestimation of the frequencies may largely be attributed to the HF method itself because this method neglects electron correlation, basis set truncation and anharmonicity. Because the overestimation is fairly uniform, a proposed solution involves the application of a single scaling factor to the theoretical harmonic frequencies to provide better agreement with the experimental data. For example in 1981, Pople et al. [3] proposed a scaling factor of 0.89 for a set of 38 molecules (477 HF/3-21G harmonic frequencies). The need for this type of correction diminishes when using density functional methods because these methods include –more or less empirically– a significant part of the total electron correlation. In 1996, Scott and Radom [4] improved the global scaling technique in their scaled quantum mechanical force field (SQMFF) procedure. They have shown that all DFT calculations, in particular those using B3-PW91/6-31G(d) or B-LYP/6-311G(df,p), were much more successful than those based on HF or MP2. At the same time, Pulay et al. [5, 6] proposed an alternative approach to the SQMFF method that involves a direct scaling of the primitive valence force constants. In a study of 60 molecules (1500 vibrational frequencies), these authors obtained an average error of less than  $8.5\text{ cm}^{-1}$  in the scaled frequencies. This result corresponds to an average percentage error smaller than 1 %. Notably when one hydrogen atom is implied, the scaling factors proposed for the stretching (C-H, N-H, O-H) and bending (X-C-H, X-O-H, X-X-H, H-N-H) vibrations are in the range of 0.87 - 0.95, whereas the scaling factor for C-Cl stretching is 1.04. The agreement between theory and experiment was further improved using larger basis sets in the DFT computations. Thus, the majority of Raman and IR experimental spectra are now assigned using DFT calculations. In addition to the data provided by these two spectroscopic techniques, inelastic neutron scattering (INS) furnishes new information because of the absence of selection rules and the predominance of scattering due to hydrogen atoms. An example of a study of symmetrical molecules is a paper describing sym- $\text{C}_6\text{F}_3\text{H}_3$  and  $\text{C}_6\text{F}_6$  [7], in which the results of DFT calculations (using five different basis sets) were used to

assign the lines observed in Raman, IR and INS spectra. The INS spectrum of  $C_6F_6$  directly exhibits seven bands of the nine vibrations that are inactive in the Raman and IR spectra. The overall RMS for the agreement between calculation and experiment is better than  $20\text{ cm}^{-1}$ , demonstrating the high quality of the theoretical predictions.

Currently, each year, several new publications are presented describing studies of substituted benzenes using a combination of Raman and IR spectroscopies and DFT or HF theoretical calculations. Curiously, no systematic spectroscopic studies of mesitylene (or 1,3,5-trimethylbenzene) derivatives have been undertaken whereas extensive studies of the protons tunneling in methyl groups have been published [8]. For the trichloro- tribromo- and triiodo-mesitylene (TCM, TBM, TIM) species, INS on powder samples [9] and on single crystals [10, 11] has been used to establish that the three methyl groups in each molecule have a different tunneling gap, which is contrary to what is expected for a molecule with true  $C_{3h}$  symmetry. For TIM, for example, three excitations were observed at 15, 24 and 89  $\mu\text{eV}$  [8, 9]. This difference in tunneling gaps is due to the difference in the intermolecular interactions exerted on the methyl groups located in non equivalent sites in the triclinic crystal cell. Neutron diffraction studies of the molecular conformations have shown additional consequences: not only the proton probability density but also the orientation differs for each methyl group [12, 13,14]. These deviations from the ideal symmetry have visible consequences in the Raman and IR spectra at low temperatures ( $< 50\text{ K}$ ): the inactive modes for  $C_{3h}$  symmetry become visible whereas the previously degenerate modes  $E'$  and  $E''$  are split [15]. The influence of intra and intermolecular symmetry on the tunneling properties of the methyl groups is still more spectacular in the case of the fully hydrogenated dibromomesitylene ( $C_9Br_2H_{10}$  or DBMH) [16]. In this compound, the methyl group surrounded by two bromine atoms behaves as a quasi-free rotor, and its tunneling gap as using INS is large (387  $\mu\text{eV}$  or  $3.1\text{ cm}^{-1}$ ). Furthermore, several torsional modes have been observed in the range 1.7 - 4 meV ( $14\text{-}32\text{ cm}^{-1}$ ) using INS [16, 17] and also Raman diffusion [18]. In contrast, the rotation of the two other methyl groups located between one bromine atom and one hydrogen atom is highly hindered, and the tunneling gap is less than 1  $\mu\text{eV}$  [17]. Therefore, we have undertaken a detailed study of the vibrational modes of DBMH with the intention to detect other spectroscopic consequences of the presence of two types of methyl groups.

## 2. Molecular conformation of dibromomesitylene DBMH.

The structure of DBMH has already been determined using neutron diffraction on a single crystal. Table 1 lists the values of the bond lengths and angles obtained using neutron diffraction at 120 and 14 K that have been published [19]. These values must be compared with the results of DFT calculations obtained under different symmetry conditions.

### 2.1. Molecular conformation of DBMH determined using neutron diffraction at 14 K and 120 K.

The thermal differential analysis indicated that DBMH exists in two different crystal phases with an order-disorder transition at  $292 \pm 0.3$  K.. The structure of the low temperature ordered phase was solved by Hernandez et al. using neutron diffraction with a single crystal [19]. The crystal is monoclinic and assigned to the space group  $P2_1/n$  with 4 molecules in the unit cell. The molecules form molecular columns parallel to the *a* axis. In these columns, the arrangement is ‘anti-dipolar’. The structure determination was performed with a great accuracy, the goodness factor R being 0.032 at 120 K and 0.026 at 14 K. The experimental bond lengths and angles are listed in Table 1 and depicted in Fig. 1b.

Table 1

Figures 1a, 1b, 1c

The hexagonal ring in DBMH is distorted. There are large differences between the six intracyclic angles, these angles vary from  $115.8^\circ$  for  $C_1-C_2-C_3$  to  $123.6^\circ$  for  $C_2-C_3-C_4$ . Thus, using a mean value of  $120^\circ$  as was used in [1, 2] should be a very crude hypothesis. Fig. 1b relative to the molecular conformation determined at 14 K suggests that the  $C_{2v}$  symmetry prediction is fulfilled to a good degree of accuracy. There is a plane of symmetry  $P_{\text{perp}}$  perpendicular to the benzene ring bisecting carbon atoms  $C_{m2}$ ,  $C_2$ ,  $C_5$  and hydrogen atom  $H_5$ . The proton probability densities (PPD) relative to  $Me_4$  and  $Me_6$  are well localized with three peaks (Fig. 2), corresponding to three well localized protons. In contrast, the PPD in  $Me_2$  is delocalized on a ring with six broad maxima from the refinement (Fig. 1b and 2) or from the map of a Fourier difference performed in the plane of the nuclei of the three hydrogen atoms of the methyl group (Fig. 2). The extra-cyclic angles  $C_1C_2C_{m2}$  and  $C_3C_2C_{m2}$  are equal to  $122.0^\circ$  in the limit of  $0.1^\circ$ , which is contrary to those observed for  $Me_4$  and  $Me_6$ , for which the extra-cyclic angles are  $120.2^\circ$  on one side and  $122.7^\circ$  on the other side (Table 1 and Fig. 1b). The small departure from symmetry observed at 120 K between the angles  $Br_1C_1C_2$  and  $Br_3C_3C_2$  was not observed at 14 K. One important result is that if the thermal librational scattering

(TLS) is considered during the neutron diffraction (ND) refinement, the benzene ring is determined to be planar with an accuracy greater than 0.005 Å. Therefore, it can be concluded from the ND results that in the solid state, the DBMH molecule retains two planes of symmetry: the ring plane ( $P_{\text{ring}}$ ) and one plane ( $P_{\text{perp}}$ ) perpendicular to  $P_{\text{ring}}$ . The intersection between these two planes is the axis of symmetry  $C_2$ .

Figure 2

## 2.2. DFT calculations of the molecular conformation

All of the computations were performed using the Gaussian 2009 program package [20]. The DFT calculation that yielded the equilibrium conformation is referred to as “Cal<sub>1</sub>  $P_{\text{ring}}$ ” in Table 1 and Fig. 1a. This conformation was obtained using the MPW1PW91 functional with the LANL2DZ basis set augmented with polarization functions on all atoms except for the hydrogen atoms. At equilibrium the benzene ring is a plane of symmetry. There is a remarkable agreement between the calculated values and the experimental data for the molecular region containing Me<sub>4</sub>, H<sub>5</sub> and Me<sub>6</sub> and good agreement for the angles around C<sub>4</sub>, C<sub>5</sub> and C<sub>6</sub> with a tolerance of 0.1°. The calculated bond distances are slightly larger than the experimental distances by only 2 to 4x10<sup>-3</sup> Å. Importantly, in the equilibrium conformation, the two methyl groups Me<sub>4</sub> and Me<sub>6</sub> have one CH bond located in the ring plane  $P_{\text{ring}}$  that points towards the atom H<sub>5</sub>. This finding is in perfect agreement with the  $C_{2v}$  symmetry observed for the molecule using neutron diffraction. These orientations of Me<sub>4</sub> and Me<sub>6</sub> groups correspond to very stable orientations: the barrier hindering the rotation of these groups has been calculated to be 490 cm<sup>-1</sup>. In contrast, the configuration calculated for the environment of the third methyl group, Me<sub>2</sub>, is not consistent with the  $C_{2v}$  symmetry. At equilibrium, the calculated structure corresponds to a methyl group with three punctual protons and one CH bond in the plane of the aromatic ring. It is clear from Fig.1a that the left and right sides of the molecule around Me<sub>2</sub> are not symmetric; this lack of symmetry occurs because the trigonal symmetry of the methyl group is not compatible with the  $C_{2v}$  symmetry of its environment which consists of two equivalent pairs of carbon-bromine atoms C<sub>1</sub>-Br<sub>1</sub> and C<sub>3</sub>-Br<sub>3</sub>. The DFT optimized geometry has only  $C_s$  symmetry, with the aromatic ring as the symmetry plane  $P_{\text{ring}}$ . It is possible to estimate the height of the barrier hindering the Me<sub>2</sub> “rotation” using this DFT computation. The calculated barrier of 53 cm<sup>-1</sup> is ten times smaller than that of the two other methyl groups. This small value partially explains why the DFT computations could lead to the actual conformation of this group. The true source of discrepancy is the strict application of the Born-Oppenheimer approximation to protons and not

considering the spin of the methyl groups: proton spins are central to the phenomenon of “proton rotational tunneling” and thus to spatial delocalization [10] and to the existence of NMR. Another molecular conformation was obtained using the DFT computation “Cal<sub>2</sub> P<sub>perp</sub>” involving the same functional and basis set, but constraining the plane P<sub>perp</sub> containing C<sub>m2</sub> and H<sub>5</sub> to be perpendicular to the initial ring plane P<sub>ring</sub> as a symmetry plane (Fig.1c). In a normal projection onto P<sub>ring</sub>, the “left” and “right” sides of the molecule are now symmetrical. The agreement between the experimental results and those obtained using Cal<sub>2</sub> is better than 0.005 Å for the bond lengths and better than 0.3° for all of the angles; therefore, Cal<sub>2</sub> provides a better representation of the molecular geometry than Cal<sub>1</sub> does for the environment around Me<sub>2</sub>. To obtain this new orientation of Me<sub>2</sub>, it is necessary to admit the loss of planarity of the aromatic ring; however, these deviations are very small: 8.10<sup>-3</sup> Å for atom C<sub>2</sub>, 15x10<sup>-3</sup> Å for atom C<sub>m2</sub> and -3x10<sup>-3</sup> Å for the Br atoms. The extra-cyclic angles around the C<sub>2</sub>-C<sub>m2</sub> bond are now equivalent and equal to 122.9°. But, a new problem has been introduced. There are now two possibilities for the displacement of the atom C<sub>m2</sub> relative to the plane P<sub>ring</sub>: i.e. above or below the plane? To summarize, the two calculations Cal<sub>1</sub> and Cal<sub>2</sub> provide molecular conformations that, with the exception of the Me<sub>2</sub> group, are not significantly different from the experimental geometry. The energy of formation found by Cal<sub>2</sub> is less favorable than that found by Cal<sub>1</sub> by only 200 cm<sup>-1</sup>, it concerns almost exclusively Me<sub>2</sub> environment. This explains that only the modes implying mainly Me<sub>2</sub> could be altered by the choice of the equilibrium structure.

### 3. Spectroscopic measurements, experimental details

The commercially purchased DBMH (Sigma-Aldrich, F) was purified by crystallization using a solution in chloroform and pumping the solvent. This step was followed by a vacuum sublimation. Small single crystals were obtained by slowly crystallizing the purified compound from a solution of methanol.

#### 3.1. Raman spectra: Table 2, Fig. 3a, 3b, 3c and 6b.

The Raman spectra of DBMH were measured using a Horiba Jobin-Yvon HR 800 Raman spectrometer. The excitation source was a He-Ne laser operating on the 632.817 nm line. The spectra were recorded at 288 K in the 80 - 4000 cm<sup>-1</sup> region with a resolution of 3 cm<sup>-1</sup>. It was possible to use single crystals with dimensions as small as 20x20x100 μm<sup>3</sup>. Thus, we obtained two types of Raman polarized spectra. In the first orientation (Fig. 3b), the polarization vector  $\mathcal{E}$  of the incident



laser beam was approximately in the *bc* crystallographic plane, i.e., it was oriented at an angle smaller than 20° relative to the DBMH ring plane. In the second orientation (Fig. 3c), the electric field  $\mathbf{\epsilon}$  was oriented approximately along the crystallographic axis *a*, i.e., it was oriented at an angle approximately 18° relative to a line perpendicular to the ring plane. Thus, in the case of similar Raman activities, the intensity of an in-plane vibration band is larger in the spectrum shown in Fig. 3b, while an out-of-plane vibration band is stronger in the spectrum shown in Fig. 3c. Table 2 lists the Raman band positions observed at 288 K (column 3) and the relative experimental intensities of these spectral features (column 16) in comparison with the calculated intensities (column 15). The spectrum recorded in the 2800 - 3200 cm<sup>-1</sup> range is shown in Fig. 6b.

### 3. 2. *Infra-red spectra at room temperature: Table 2, Fig. 4a, 4b & 6c.*

The infra-red absorption spectra were measured from 400 to 4000 cm<sup>-1</sup>, using a Fourier transform spectrograph (Bruker Equinox 55) and using the KBr pellet technique. The resolution was approximately 2.5 cm<sup>-1</sup>. The IR range of 100 - 400 cm<sup>-1</sup> was investigated using another Fourier transform spectrograph (IFS 125 MR) with synchrotron radiation on the AILES line at SOLEIL. The sample was a pellet of compressed pure DBMH that was 1mm thick. The resolution was better than 1 cm<sup>-1</sup>. The frequencies of the IR absorption lines observed experimentally in the range (Fig. 4b) are given in Table 2, column 4. The calculated relative intensities of these features are listed in column 17, and those observed experimentally are in column 18. The spectrum recorded in the 2800 – 3200 cm<sup>-1</sup> region is depicted in Fig.6c.

### 3. 3. *Inelastic neutron scattering spectra*

The INS spectra were obtained using the TOSCA-II inelastic spectrometer of the ISIS pulsed neutron source at Rutherford Appleton Laboratory (Chilton, Didcot, U.K.). TOSCA-II is a crystal-analyzer inverse-geometry spectrometer. The incident neutron beam spans a broad energy range that covers the energy transfer region from 4 to 500 meV (30 - 4000 cm<sup>-1</sup>). The final energy of the scattered neutrons (approximately 3.35 meV) was selected by two sets of pyrolytic graphite crystals placed nearly at 43° and 138° relative to the incident beam. Because of this fixed geometry, the wave-vector transfer *Q* is roughly proportional to the square root of the incoming neutron energy. The energy resolution in the accessible energy range allows a resolution better than 3 cm<sup>-1</sup> for *Q* smaller than 700 cm<sup>-1</sup>. The samples were pellets approximately 1 mm thick, resulting in a surface area of 3x4 cm<sup>2</sup>,

packed in an aluminum pocket. Approximately 2 g of DBMH was used, and the samples were frozen below 30 K in an “orange” helium cryostat. The frequencies of the maxima of the scattered neutrons are given in Table 2, column 5, and illustrated by Fig.5.

#### 4. Calculations of the frequencies of the normal modes of vibration

##### 4.1. Comparison of “harmonic”, “harmonic scaled” and “anharmonic” frequencies calculated by DFT, the ring plane $P_{ring}$ being a plane of symmetry

Table 2 lists the results of three DFT calculations of the normal modes of vibration that were performed using the Gaussian 2010 [20] program package. These results were obtained using the MPW1PW91 functional with the LANL2DZ basis set augmented with polarization functions. The equilibrium molecular geometry is that presented in section 2.2, the ring plane  $P_{ring}$  being a plane of symmetry. The DBMH molecule containing 21 atoms has 57 optical modes of vibration. Relatively to the plane of symmetry, the 57 modes of vibration may be divided into two types: 36 are of symmetry  $A'$  (in-plane), and 21 are of symmetry  $A''$  (out-of-plane) (see column 2, Table 2). The calculated modes are numbered from the largest frequency to the smallest within each fundamental frequency. The frequencies may be compared with the experimental results from the Raman spectroscopy at 288 K (column 3) or, if the Raman intensities are too small, to the IR data (column 4) or to the INS data (column 5). The  $\omega_h$  wave numbers (column 10) correspond to the frequencies that were calculated by assuming the harmonic approximation without any scaling. These  $\omega_h$  frequencies are larger than the experimental frequencies by 14 to 72  $\text{cm}^{-1}$  in the range of 600 - 1700  $\text{cm}^{-1}$  (i.e. 2 to 4 %), by 160 to 230  $\text{cm}^{-1}$  in the range of 2800 – 3250  $\text{cm}^{-1}$  (i.e. 6 to 8 %), while between 150 and 600  $\text{cm}^{-1}$  there are fluctuations amounting to  $\pm 3$  %. The five smallest calculated frequencies should be treated separately because they must contain the three methyl groups “torsions” and are very sensitive to the type of calculation. These DFT computations also predict the quantitative specificities of the IR and Raman lines. In fact, the Gaussian program calculates the Raman activity  $S_i$  for each normal mode  $n^\circ i$  of an isolated molecule. We calculated the Raman activity in parallel to the  $\omega_h$  frequencies assuming the harmonic approximation (see column 14 Table 2). The result was then converted to Raman intensity using the following relationship derived from the intensity theory of Raman scattering [21, 22, 23]

$$I_i = C (\omega_0 - \omega_i)^4 S_i / \omega_i [1 - \exp(-h c \omega_{Ri} / kT)]$$

where  $\omega_0$  is the laser excitation frequency in  $\text{cm}^{-1}$  (in our experiments  $\omega_0 = 15802 \text{ cm}^{-1}$ )  $\omega_i$  is the vibrational frequency of the  $i$ -th normal mode in  $\text{cm}^{-1}$ ,  $h$  and  $k$  are the Planck and Boltzmann

constants,  $T$  is the temperature in Kelvin and  $C$  is a suitably constant chosen that serves as a common normalization factor for all peak intensities. For these spectra, we chose a relative intensity of 100 for excitation  $n^\circ 41$  calculated at  $\omega_h = 560 \text{ cm}^{-1}$  and observed at  $\omega_i = 563 \text{ cm}^{-1}$ , because it is for this band that the intensity calculated in the range  $200\text{-}1600 \text{ cm}^{-1}$  is the highest. The relative intensities calculated for the Raman spectra are found in Table 2, column 15, and are represented in Fig. 3a. These intensities must be compared to those observed in the spectra: (column 16). However, because the Raman spectra were obtained with oriented crystals, it is possible to independently estimate those intensities arising from “in-plane” vibrations (Fig. 3b) and those arising from vibrations perpendicular to the molecular plane (Fig. 3c). Nevertheless, this estimation must be performed with caution because the spectra were obtained near room temperature, which means that the lines are broadened and only peak maximum intensities may be estimated.

To obtain a better agreement between theory and experiment, it is customary to scale down the calculated harmonic frequencies [4, 5, 6]. Recently, empirical scaling factors have been proposed after DFT calculations on benzenic compounds: 0.957 for tolunitrile [24], 0.960 for p-chlorobenzoic acid [25] and 0.958 and 0.983 for halogenoanilines above and below  $1700 \text{ cm}^{-1}$ , respectively [26, 27, 28]. Column 9b lists the values of the ratio  $\omega_R / \omega_h$  between our Raman experimental and calculated values. These values vary from 0.935 to 1.013. As a result, we multiplied frequencies  $\omega_h$  by the “often used” scaling factors of 0.958 for frequencies larger than  $1800 \text{ cm}^{-1}$  and 0.983 for the other frequencies, thus obtaining the  $\omega_{hsc}$  frequencies listed in column 8. The new differences between theory and experiment ( $\omega_{hsc} - \omega_R$ ) (column 9) are smaller than the preceding differences ( $\omega_h - \omega_R$ ) by a factor 2 to 3 for frequencies larger than  $700 \text{ cm}^{-1}$ . Only the frequencies calculated for the C-H stretching vibrations remain too large, by 40 to  $95 \text{ cm}^{-1}$ , and could be better estimated using a scaling factor of approximately 0.94. It is essential to determine whether the discrepancies between theory and experiment are mainly due to the neglect of the anharmonicity of the potential energy in the calculations. Thus, we computed the “anharmonic frequencies” ( $\omega_{anh}$ ) using the automated code implemented in Gaussian by Barone [29]. Starting from the optimized geometry, the program computes the third and semi-diagonal fourth derivatives of the formation energy and then evaluates the  $\omega_{anh}$  frequencies of the fundamental bands using second order perturbation theory (keyword: *freq(anharmonic)*). This calculation also provides the frequencies of the overtones ( $\omega_{anh}^o$ ) and those of the combination bands ( $\omega_{anh}^c$ ). The  $\omega_{anh}$  frequencies of DBMH are listed in column 6 of Table 2, and the deviations from the experimental values,  $\omega_{anh} - \omega_R$  are given in column 7. It appears that the agreement is remarkable in the entire range of  $130\text{-}1500 \text{ cm}^{-1}$ , this agreement is even better than that

observed for the scaled frequencies; the deviations are smaller than  $15\text{ cm}^{-1}$  for 39 of the 41 modes. The agreement is also greatly improved for the C-H stretching modes ( $\omega_{\text{anh}}$ ) and is now overestimated by less than 3% (compared to the previous 8%) for the “harmonic” frequencies  $\omega_{\text{h}}$ . These results for DBMH, a poly-methylated molecule, confirm the ability of the analytical second-order perturbation treatment to consider the anharmonicity during the calculation of the internal modes of vibration. Similar results have also been observed in studies of azabenzenes [29, 30, 31], pyrrole and furan [32], uracil [33] or, more recently, monofluoroanilines [26]. The anharmonic frequencies  $\omega_{\text{anh}}$  obtained for DBMH are in very good agreement with the experimental values, and this agreement is even better than that observed for  $\omega_{\text{hsc}}$ , whose values were obtained semi-empirically using the scaled quantum mechanical method. Thus, for the assignment of each type of motion to a specific frequency, we will only compare  $\omega_{\text{R}}$  (or  $\omega_{\text{IR}}$ ) and  $\omega_{\text{anh}}$ .

#### *4.2. Calculations of the frequencies of the normal modes of vibration assuming that the plane of symmetry is the plane $P_{\text{perp}}$ perpendicular to the benzene ring: Table 2 columns 12-13*

In the case of DBMH, it must be remembered that the model used in the DFT calculations does not perfectly represent the environment of the methyl group  $\text{Me}_2$ . Section 2.2 discussed another calculation that was performed when the molecule is constrained to have a symmetry plane  $P_{\text{perp}}$  normal to the ring plane. This conformation is illustrated in Fig. 1c and takes into account the symmetry around the  $\text{Br}_1$  and  $\text{Br}_2$  atoms and the  $\text{Me}_2$  substituent. It corresponds approximately to the neutron diffraction structure although it partially breaks the planarity of the benzene ring. However, for DFT calculations, this conformation corresponds to an unstable equilibrium in energy and is located on a saddle point. It is of interest to determine whether this conformation still provides valuable vibrational frequencies. Consequently, a fourth set of data was calculated using the harmonic approximation which is presented as  $\omega_{\text{hp}}$  in column 12. The vibrations of the new symmetries A' and A'' correspond to symmetrical (Sym) and antisymmetrical (AntiSym) motions, respectively for the bonds that are equivalent by symmetry such as  $\text{C}_1\text{-Br}_1$  and  $\text{C}_3\text{-Br}_3$  for example. The calculated frequencies are listed in Table 2, column 12. All of the differences ( $\omega_{\text{h}} - \omega_{\text{hp}}$ ) are smaller than  $4\text{ cm}^{-1}$  (column 13), with the exception of the lowest frequencies that are in direct connection with the  $\text{Me}_2$  torsion. Thus, the majority of the overestimation of the unscaled harmonic frequencies is effectively caused by the anharmonicity of the atomic motions, particularly those of the hydrogen atoms. The two geometric models, that are slightly “imperfect” relatively to the neutron

scattering structure, are practically equivalent for the DFT calculations, except for mode n° 57, which is relative to the torsion of the Me<sub>2</sub> group (see below).

**Table 2**

#### 4.3. A distinction between the internal vibration modes and the lattice modes

For a DBMH crystal below 292 K, there are four molecules in the monoclinic P2<sub>1</sub>2<sub>1</sub>2<sub>1</sub> cell, thus, 6 Z "lattice modes" are predicted, where Z is the number of molecules in the unit cell. Consequently, 21 new optical modes are expected. Generally, the intensity of the lattice modes in the Raman spectra is smaller than that of the internal modes. In a previous study conducted at approximately 4 K [18], almost all of these lattice modes occurred at frequencies smaller than 100 cm<sup>-1</sup>. The widths of the bands grew quickly as the temperature increases, so that the peaks are very flat at 288 K and difficult to detect. In Fig. 5, five peaks with maximum scattering of neutrons are detected below 135 cm<sup>-1</sup> at a temperature of approximately 10 K. In these experiments, the scattering due to hydrogen atoms is the most intense phenomenon. This region has been studied in previous experiments examining the proton tunneling of the methyl groups [16, 17]. Thus, if peaks at 61, 79, 91 cm<sup>-1</sup> originate in transitions 1<sup>A</sup>-2<sup>A</sup>, O<sup>E</sup>-2<sup>E</sup> and O<sup>A</sup>-2<sup>A</sup> between torsional levels of Me<sub>2</sub>, only the peaks at 117 (124 in R), 135 and 156 (157 in R) may be assigned to internal modes.

### 5. Vibrational spectra and their assignments

#### 5.1. Labelling of the internal coordinates used to describe the vibrations.

The labeling used to describe the vibrations is that proposed by Decius and alt. [34]. When the atomic displacements occur in the symmetry plane (i.p.), the variations in the lengths or angles (Fig. 1b) are labeled as follows: t<sub>n</sub> = stretching of a C<sub>n</sub>-C<sub>n+1</sub> bond into the ring (n=1-6), s<sub>p</sub> = stretching of a C<sub>p</sub>-C<sub>mp</sub> bond for a methyl substituent n° p (p = 2, 4, 6), s<sub>i</sub> = stretching of a C<sub>i</sub>-Br<sub>i</sub> bond for a bromine substituent n° i (i = 1 or 3), s<sub>5</sub> = stretching of the C<sub>5</sub>-H<sub>5</sub> bond, α<sub>n</sub> = C<sub>n-1</sub>-C<sub>n</sub>-C<sub>n+1</sub> ring angle breathing, and β<sub>n</sub> = in-plane bending of a C<sub>n</sub>-X<sub>n</sub> bond, knowing that φ<sub>nd</sub> = X<sub>n</sub>-C<sub>n</sub>-C<sub>n+1</sub> and φ<sub>ng</sub> = C<sub>n-1</sub>-C<sub>n</sub>-X<sub>n</sub> are outer ring angles, and in consequence β<sub>n</sub> = φ<sub>nd</sub> - φ<sub>ng</sub> and α<sub>n</sub> = 360° - (φ<sub>nd</sub> + φ<sub>ng</sub>). A positive φ<sub>n</sub> = φ<sub>nd</sub> + φ<sub>ng</sub> angle corresponds to a clockwise motion. The following labels apply for a motion within a methyl group n° p: r<sub>pa</sub> = stretching of the C<sub>mp</sub>-H<sub>pa</sub> bond (in the ring plane), r<sub>pb</sub> and r<sub>pc</sub> correspond to the two staggered C-H bonds, and η<sub>p</sub> = 'i.p.' torsion of the axis of this methyl group Me<sub>p</sub>. For example, η<sub>4</sub> corresponds to symmetrical rotations of planes C<sub>4</sub>-C<sub>m4</sub>H<sub>4b</sub> and C<sub>4</sub>-C<sub>m4</sub>-H<sub>4c</sub> around the

$C_4-C_{m4}$  bond relative to the ring plane.  $\chi_{pa}$  denotes the  $C_p-C_{mp}-H_{pa}$  variation with  $\chi_p = \chi_{pa} + \chi_{pb} + \chi_{pc}$ , whereas  $\theta_{pa}$  denotes the  $H_{pb}-C_{mp}-H_{pc}$  angle variation with  $\theta_p = \theta_{pa} + \theta_{pb} + \theta_{pc}$ .

The out-of-plane (o.p.) motions are labeled as follows:  $\gamma_n$  = out-of-plane bending of a  $C_n-X_n$  bond, it is also named wagging [2],  $\tau_i$  = out of-plane benzene ring torsion of a  $C_i-C_{i+1}$  bond and  $\delta_p$  = sum of out-of-plane torsion of the methyl C-H bonds around the  $C_p-C_{mp}$  bond. For deformations inside the methyl groups,  $r''_{mp} = C_{mp}-H_{pi}$  for the o-p methyl group stretching ( $i = a, b$  or  $c$ ) when the atom  $C_{mp}$  and  $H_{pi}$  move out of the ring plane  $P_{ring}$ ,  $\chi''_p$  = sum of angles for the  $C_p-H_{pi}-H_{pi+1}$  variations,  $\theta''_p$  = sum of angles for the variations  $C_p-C_{mp}-H_i$ . Two remarks must be added. First, if during a large wagging  $\gamma_n$  the substituent  $X_n$  moves up then the carbon  $C_n$  is moving down relative to the atoms  $C_{n-1}$  and  $C_{n+1}$  in the ring, this generates small ring torsions  $\tau_{i-1}$  and  $\tau_i$ ; secondly, for a ring torsion  $\tau_i$  there is also a displacement of the bond  $C_i-X_i$  in the same direction than  $C_i$ , this will be noted  $\kappa_i$  and is distinct from a  $\gamma_i$  for the PED calculations.

## 5.2. Calculation of the potential energy distribution (PED) for each mode.

To understand the molecular dynamics, it is useful to visualize the individual deformations of the bonds and angles. A first approach was obtained based on the information presented in Fig. 7, 8 and 9. These figures were drawn using the Molekel 98 package to visualize the deformations that are characteristics of each frequency  $\omega_h$ . This qualitative approach was completed using a quantitative estimation of the potential energy distribution, the PED. The PED was obtained using the key-word *freq(internal)* in the Gaussian package. Table 3 provides the contribution to the PED for each mode due to the principal in-plane deformations. Table 4 lists the PED for the out-of-plane deformations. The two approaches are not strictly equivalent because a large displacement of a light hydrogen atom has a smaller relative contribution to the PED than a displacement of a Br atom or a Me group does. These tables list the main deformations, for equivalent atoms, the + sign indicates motions in-phase, and the – sign indicates opposite motions. The numbers indicate the percentage of each contribution to the PED.

**Table 3**

## 5.3. A criterion for validity: the calculation-experiment agreement for the most intense frequencies:

A first test to confirm the ability of the DFT calculations to accurately predict the experimental data is to examine the sequence of the intense lines in the IR and Raman spectra that are presented in Table 2 (columns 15-18) and in Figs. 3 and 4. The criterion for the comparison requires the relative

intensity to be larger than 50 for the experiment and greater than 30 for the calculations (reference 100 for band n° 41 observed at  $560\text{ cm}^{-1}$  in the Raman spectra). The experimental and calculated frequencies will be compared and indicated in the following order: in first the frequency observed,  $\omega_R$ , and then, inside parentheses, ( $\omega_{\text{anh}}$ ) the “anharmonic” frequency calculated using DFT. The intense bands observed in the Raman spectra are located at: 157 (153), 252 (248), 561 (560), 590 (586), 1306 (1313), 1387 (1392) and  $2923\text{ (2977)}\text{ cm}^{-1}$ . The intense bands in the I.R. spectra are observed at: 190 (166), 228 (222), 253 (248), 296 (288), 636 (637), 852 (854), 960 (962), 1043 (1041), 1375 (1378), 1456 (1460) and  $2920\text{ (2977)}\text{ cm}^{-1}$ . The inelastic neutron scattering spectrum (Fig. 5) confirms the presence of essentially all of the modes calculated below  $1000\text{ cm}^{-1}$ , particularly several weak peaks that were difficult to detect in the IR or Raman spectra. For example, column 5 of Table 2, one may note the high intensity for the scattered lines at 193, 365, 542, 853 and  $1031\text{ cm}^{-1}$ . The agreement between the calculated and observed frequencies is remarkable. The success of the first inspection encouraged us to pursue a detailed assignment of the bands observed experimentally.

#### *5.4 Regarding the labels: in-plane or out-of-plane mode, “methyl” or “skeleton” vibrations, symmetric or anti-symmetric modes*

Because there are 21 atoms in the DBMH molecule, 57 normal modes of vibration are expected. In our calculations the plane  $P_{\text{ring}}$  is a plane of symmetry for the equilibrium conformation. Thus, there are 36 vibration frequencies with symmetry  $A'$  for which the motions of the Br, C, and  $H_a$  atoms occur within the plane  $P_{\text{ring}}$ . These vibrations are illustrated in Figs. 7 and 8. For the 21 other vibrations which have symmetry  $A''$ , these atoms are moving normally to  $P_{\text{ring}}$  as shown in Fig. 9. Of the 57 modes of vibration, 27 modes must correspond to the stretching of the  $C_m\text{-H}$  bonds ( $r_{\text{mp}}$ ,  $r''_{\text{mp}}$ ) or to deformations of the  $C_{\text{ring}}\text{-C}_{\text{methyl}}\text{-H}_i$  entities ( $\delta_{\text{pp}}$ ,  $\chi_p$ ,  $\theta_p$ ,  $\chi''_p$ ,  $\theta''_p$ ). These modes will be entitled “methyl modes of vibrations”. The criterion to validate this label is that the contribution of the methyl deformations to the PED is larger than 60 %. As mentioned in the introduction section, the methyl groups linked to an aromatic ring are characterized by specific modes concentrated in narrow bands, appearing around  $3000\text{ cm}^{-1}$ ,  $1445\text{ cm}^{-1}$ ,  $1380\text{ cm}^{-1}$ ,  $1060\text{ cm}^{-1}$  and  $1000\text{ cm}^{-1}$  and below  $150\text{ cm}^{-1}$ . The 30 other modes must therefore mainly involve the aromatic ring and its substituents and may be labeled as “skeleton vibrations”. These vibrations are located throughout the region below  $1700\text{ cm}^{-1}$ . A rapid examination of Table 2 shows that 24 frequencies are smaller than  $1000\text{ cm}^{-1}$  and may correspond to 21 skeleton modes and the three methyl torsion modes. Thus, only 9 “skeleton modes” are expected in the range of  $1000 - 1700\text{ cm}^{-1}$ . The pictures in Fig. 7 reveal the importance

of the pseudo-symmetry relative to  $P_{\text{perp}}$  perpendicular to  $P_{\text{ring}}$ . DFT calculations done with the conformation of Fig. 1c are given Table 2, column 12b. If the atoms that are located symmetrically relative to this plane have in-phase displacements, the vibrations are of symmetry species A', if they are moving in opposite directions they have the symmetry species A". In the text, we will refer to vibrations as Sym or Antisym, respectively.

### 5.5. Assignment of specific in-plane vibrations to normal modes of species A'.

To describe the spectral assignments, the characteristic number of the vibration in order of decreasing observed frequency is given in first, it is followed by the experimental frequency  $\omega_R$  (or  $\omega_{IR}$ ) and by the calculated frequency  $\omega_{\text{anh}}$  written within parentheses. The Sym "skeleton modes" of species A' correspond to the eleven pictures in the upper part of Fig. 7, whereas the ten Antisym modes are shown in the lower part of the figure. The "methyl modes" are represented in Fig. 8.

#### 5.5.1. In plane skeleton modes implying large radial displacements (Tables 2 and 3, Figs. 7)

Mode n° 49,  $\omega_R = 253$  (248)  $\text{cm}^{-1}$ , involves Sym stretches of the  $C_1\text{-Br}_1$  and  $C_3\text{-Br}_3$  bonds (22%) with a Sym variation of the opposite  $\alpha$  angles (9%) it is associated with Sym bendings of  $C_4\text{-C}_{m4}$  and  $C_6\text{-C}_{m6}$  ( $\beta_4\text{-}\beta_6$ ) (20%). All of these motions are in perfect agreement with the presence of a mirror of symmetry, which is the plane  $P_{\text{perp}}$ . Mode n° 47,  $\omega_{IR} = 295$  (288)  $\text{cm}^{-1}$ , consists in AntiSym motions of the same atoms: stretches (18%) and bending (15%), of C-Br bonds which are associated with AntiSym breathing vibrations involving angles  $-\alpha_4 + \alpha_6$  opposite to methyl groups 4 and 6) (24%). These two modes, n° 49 and 47, illustrate the splitting of the "C-Br stretching vibration" as a result of a coupling between the motions of the  $C_1\text{-Br}_1$  and  $C_3\text{-Br}_3$  bonds. The Sym mode n° 41,  $\omega_R = 561$  (560)  $\text{cm}^{-1}$ , is mainly due to ring breathing: two Sym angles enlargements ( $\alpha_1 + \alpha_3$ ) (20%), two angle reductions ( $-\alpha_2 - \alpha_5$ ) (28%). Mode n° 40,  $\omega_R = 590$  (586)  $\text{cm}^{-1}$ , is also a Sym mode, it is mainly due to Sym C-C<sub>m</sub> stretching vibrations (14%) plus reduction of the corresponding angles ( $\alpha_4 + \alpha_5 + \alpha_6$ ) (44%), Mode n° 38,  $\omega_{IR} = 636$  (637)  $\text{cm}^{-1}$ , has contributions similar to those of n° 40 but arranged in an AntiSym manner. Sym mode n° 35,  $\omega_R = 940$  (947)  $\text{cm}^{-1}$ , involves: Sym stretching vibrations of the C-C<sub>m</sub> bonds (11%), cycle angles variations (15%) and Sym. tilting vibrations (52%) of the Me<sub>4</sub> and Me<sub>6</sub> groups. These deformations are similar in nature and intensity to the "methyl" rocking vibrations indicating strong interactions between the ring and methyl group motions in this frequency range. For AntiSym mode n° 34,  $\omega_{IR} = 962$  (960)  $\text{cm}^{-1}$ ,



almost one half of the contributions are attributed to AntiSym pairs of C-Br and C-C<sub>m</sub> stretching vibrations: (16%) and the corresponding angles variations (21%), and the other half is a result of the Me<sub>2</sub> tilting vibration (34%). Note that for these two modes, n<sup>os</sup> 34 and 35, the methyl groups have a contribution to the PED as large as that of the “skeleton” motions. The Sym mode n° 27,  $\omega_{\text{IR}} = 1047$  (1043) cm<sup>-1</sup>, can be described as the trigonal ring-breathing mode or the ‘star of David’ vibration, it implies dilatations of three ring angles (25%) accompanied by contractions of the three others (23%). Sym mode n° 24,  $\omega_{\text{R}} = 1305$  (1314+1322) cm<sup>-1</sup>, is also a trigonal breathing mode consisting of stretching motions of the three C<sub>p</sub>-C<sub>mp</sub> bonds (15%) associated with alternate ring angles breathing (12%) while methyl umbrella deformations contribute 39%. Sym mode n° 20,  $\omega_{\text{R}} = 1387$  (1392) cm<sup>-1</sup>, is due to Sym stretching motion C<sub>p</sub>-C<sub>mp</sub> bonds (13%) but again contributions from the umbrella bending motions of Me<sub>2</sub> and Me<sub>6</sub> (65%) are larger than those of the skeleton indicating strong interactions in this frequency region. For Sym mode n° 1,  $\omega_{\text{IR}} 3014$  cm<sup>-1</sup> (3085), almost all the PED is due to stretching of the C<sub>5</sub>H<sub>5</sub> bond.

#### 5.5.2. In plane skeleton modes implying large tangential displacements (Tables 2 and 3, Figs. 7)

Among the remaining i.p. skeleton modes only four have relative intensities in the calculated Raman spectra that are larger than 20%: n<sup>os</sup> 53, 46, 11 and 12, which is in agreement with the experimental results. N° 53,  $\omega_{\text{R}} = 157$  cm<sup>-1</sup> (153), primarily consists of the Sym scissoring motion of the two C-Br bonds (52%) with a slight stretching contribution (14%). N° 50,  $\omega_{\text{R}} = (229$  cm<sup>-1</sup> (222), also involves the two C<sub>i</sub>-Br<sub>i</sub> bonds, but the bending motion in this mode is AntiSym (33%) with a small stretching character ( $s_1 - s_3$ ) (17%). N° 46 is observed at 323 cm<sup>-1</sup> (315), it involves the Sym bending (36%) of the pair C<sub>4</sub>-Me<sub>4</sub> and C<sub>6</sub>-Me<sub>6</sub> plus a Me<sub>4</sub> and Me<sub>6</sub> tilting motion (23%) N° 44 is observed at 374 cm<sup>-1</sup> (360), this mode consists of the AntiSym bending motion of the C-C<sub>m</sub> bonds (32%) and in-plane tilting motions of Me groups (24%). Mode n° 42 is observed 541 cm<sup>-1</sup> (542), it is an AntiSym bending motion involving an in-phase oscillation of the C-Br and C-C<sub>m</sub> bonds ( $\sum^n \beta_n \sim 50\%$ ) accompanied by in-plane tilting of the methyl group (34%). The AntiSym bending mode n° 26 (1217 cm<sup>-1</sup> (1235) includes contributions from the bending vibration of the C<sub>5</sub>-H<sub>5</sub> bond (28%) and tilting of the three methyl groups (34%). The mode n° 25,  $\omega_{\text{R}} = 1306$  cm<sup>-1</sup> (1314), corresponds to a trigonal deformation of the benzene ring caused by C<sub>a</sub>-C<sub>a</sub> stretching motions (33%) and large tilting motions of the methyl groups (42%). This mode is also referred to as the Kekule mode. The Sym stretching mode n° 21 may be a part of the broad band observed at 1387 cm<sup>-1</sup>, this “skeleton” mode has very significant contributions from methyl groups. The AntiSym mode n° 19,  $\omega_{\text{R}} = 1392$  (1397) cm<sup>-1</sup>, may

correspond to the peak in the IR spectrum at  $1392\text{ cm}^{-1}$ . It consists of C-Br and C-C<sub>m</sub> and C<sub>5</sub>-H<sub>5</sub> motions (30%), and to a larger extent, umbrella deformations of the methyl groups (52%). Mode n° 12,  $\omega_R = 1562\text{ cm}^{-1}$  (1588), corresponds to deformations of the C-C bonds in the benzene ring (32%), the vibrations are roughly parallel to the direction of the  $\mathcal{C}_2$  axis. These vibrations are accompanied by a large narrowing of the  $\alpha_2$  and  $\alpha_5$  angles (15%) and enlargements of the  $(\alpha_1+\alpha_3+\alpha_4+\alpha_6)$  angles (15%) $\square$ . AntiSym mode n° 11 has been detected at  $1587\text{ cm}^{-1}$  (1629) with  $I_R = 15$  (28). The deformations in the ring are perpendicular to  $P_{\text{perp}}$  and the  $\alpha_2$  and  $\alpha_5$  angles remain constant. All of these excitations are also detected in the INS spectrum (Fig. 5).

### 5.5.3. Assignment of the in-plane “methyl groups” modes of vibration

We previously proposed that the criteria for this labeling requires small contributions from the s, t,  $\alpha, \phi$  and/or  $\beta$  deformations but large contributions from the H-C-H ( $\theta$ ) and C-C-H ( $\chi$ ) bonds motions ( $> 30\%$ ) and from in-plane torsions ( $\eta$ ) of the C-C<sub>m</sub> bonds, i.e., more than 60% contribution to the PED from the methyl groups. In our experiments at 288 K, we observed a broadening of the lines so that the resolution was generally not sufficient to assign individual frequencies for Me<sub>4</sub> and Me<sub>6</sub> and individual intensities into these broad lines.

*The calculated rocking modes of CH<sub>3</sub>* are: N° 33,  $\omega_R = 1013\text{ cm}^{-1}$ , n° 32,  $\omega_R = 1028\text{ cm}^{-1}$ , n° 29  $\omega_R = 1024\text{ cm}^{-1}$ . The contribution of  $\eta_{\text{pp}}$  torsions to the total PED is almost equal to that of internal methyl deformations ( $\chi, \theta$ ). Small bands are observed at  $1015$  and  $1032\text{ cm}^{-1}$  in the Raman spectra (Fig.3) but it would be necessary to obtain spectra at much lower temperatures to propose a secure detailed assignment.

*In-plane bending modes (or umbrella modes):* For mode n° 23,  $\omega_R = 1380\text{ cm}^{-1}$  (1378), the deformations of the Me<sub>4</sub> and Me<sub>6</sub> groups are in phase but out of phase relative to those of Me<sub>2</sub>, For n° 22,  $\omega_R = 1387\text{ cm}^{-1}$  (1384), the deformations of Me<sub>4</sub> and Me<sub>6</sub> are in phase while Me<sub>2</sub> is at rest. For mode n° 20,  $\omega_R = 1387\text{ cm}^{-1}$  (1393), the deformations of Me<sub>2</sub> and Me<sub>6</sub> are in phase and out of phase to the small deformations of Me<sub>4</sub>. The total intensity for these modes is distributed into a broad Raman band at  $1387\text{ cm}^{-1}$  and must be associated with modes n° 19 and 20.

*In-plane bending (or scissor) modes of CH<sub>3</sub>:* In the IR spectra the sum of the relative intensities for n° 16 and that of the neighboring out-of-plane modes n° 17-18 is  $I_{\text{IR}} = 45$  in agreement with 30 calculated. For mode n° 16,  $\omega_R = 1435\text{ cm}^{-1}$  (1438), the torsions  $\eta_2$ ,  $\eta_4$  and  $\eta_6$  are calculated in phase while only the clippings for Me<sub>2</sub> and Me<sub>6</sub> are in phase. For mode n° 14,  $\omega_R = 1456\text{ cm}^{-1}$

(1462), the main contributions to the PED is calculated for Me<sub>4</sub>. For n° 13,  $\omega_R = 1456 \text{ cm}^{-1}$  (1460), Me<sub>4</sub> is almost at rest, while Me<sub>2</sub> and Me<sub>6</sub> are moving in opposition.

*Symmetric in-plane stretching (or breathing) modes of CH<sub>3</sub>:* N°s 10, 9 and 8 correspond to the broad and intense Raman line at  $\omega_R = 2923 \text{ cm}^{-1}$  (2977). For n° 10, groups Me<sub>4</sub> (48%) and Me<sub>6</sub> (42%) exhibit breathing motions that are out-of-phase (o.p.) For n° 9, these vibrations are in-phase (i.p.), and for n° 8, only Me<sub>2</sub> vibrates (88%).

*Asymmetric in-plane stretching modes of CH<sub>3</sub>:* For these modes, in each methyl group the in-plane bonds C<sub>pm</sub>-H<sub>pa</sub> move in opposite direction to the staggered C<sub>pm</sub>-H<sub>pb</sub> and C<sub>pm</sub>-H<sub>pc</sub> bonds. Modes n°s 3 and 4 correspond to  $\omega_R \sim 2978 \text{ cm}^{-1}$  (3041) with I<sub>R</sub> = 12 (36). N° 2 corresponds to motions only within Me<sub>2</sub>. This mode was calculated to be at  $3066 \text{ cm}^{-1}$  with a weak intensity, experimentally it appears as a bump in the IR spectra at  $2985 \text{ cm}^{-1}$ .

## 5.6. Out of plane modes (Tables 2 and 3, Figs. 3, 4 and 5)

In a Raman spectrum from a powdered sample of DBMH, it is more difficult to assign the o.p. modes than it is to assign the i.p. modes because the former have weaker intensities. Thus, we collected the Raman spectra shown in Fig. 3c using a small single crystal excited by a laser beam that was polarized approximately along the *b* axis of the crystal. Consequently, in this spectrum, the intensities of the i.p. vibrations are reduced and those of the o.p. modes are relatively enhanced which helps with the assignment. Nine normal modes of vibration of species A'' were calculated and observed in the range of  $1000\text{-}3050 \text{ cm}^{-1}$ . For each of them the contribution of the CH<sub>3</sub> groups to the PED is larger than 60%, so they are labelled "methyl o-p modes". They are represented in the lower part of Fig. 9. Six A'' modes located in the range of  $200\text{-}1000 \text{ cm}^{-1}$  were assigned as skeleton modes. For the six remaining modes, the torsion (hindered rotation) of the methyl groups plays a major part.

### 5.6.1. "Methyl groups" o. p. modes with frequencies higher than $1000 \text{ cm}^{-1}$

Table 4. Fig. 6 and 9

At 288 K, a broad band is observed at  $\omega_R = 2956 \text{ cm}^{-1}$  in the Raman spectra and at  $2954 \text{ cm}^{-1}$  in the IR spectrum. This band is located between two packets of i.p. modes already assigned: n°s 3 and 4 observed at  $\omega_R = 2978 \text{ cm}^{-1}$  (3041) and n°s 8, 9, 10 observed at  $2923 \text{ cm}^{-1}$  (2977). Thus, this band may be assigned to modes n°s 5, 6 and 7 calculated at  $\omega_{anh} = 3015 \text{ cm}^{-1}$ . For each mode, the internal motions involve stretchings of the two staggered C<sub>m</sub>-H<sub>mb</sub> and C<sub>m</sub>-H<sub>mc</sub> bonds, into a specific Me

group. The modes n<sup>o</sup>s 15, 17 and 18 at  $\omega_R = 1435 \text{ cm}^{-1}$  (1440) are caused by the o.p. tilting motion  $\delta$  and a type of clipping ( $\chi''$ ,  $\theta''$ ) within the methyl groups in almost equal proportions. The mode n<sup>o</sup> 15 involves only Me<sub>2</sub> ( $\delta_2 = 47\%$ ). Mode n<sup>o</sup> 18 implies out-of-phase vibrations of Me<sub>4</sub> and Me<sub>6</sub> ( $\delta_4 + \delta_6 = 48\%$ ) whereas those vibrations are in-phase for mode n<sup>o</sup> 17 ( $\delta_4 + \delta_6 = 49\%$ ). Modes n<sup>o</sup>s 28 and 30 ( $\omega_R = 1045 \text{ cm}^{-1}$  (1035)) are characteristic of the o.p. rocking motions of the methyl groups. Mode n<sup>o</sup> 28 concerns in phase o.p. vibrations involving Me<sub>4</sub> and Me<sub>6</sub>. The vibrations are out-of-phase for mode n<sup>o</sup> 30 whereas for mode n<sup>o</sup> 31,  $\omega_R = 1032 \text{ cm}^{-1}$  (1028) only Me<sub>2</sub> is excited.

### 5.6.2. “Skeleton out-of-plane” modes

Table 4 gives a detailed panorama of the motions for each mode which are shown in the upper part of Fig. 8. These motions are complex and then difficult to describe in detail. Thus, we will indicate only the main contributions to the PED. For example, Sym mode n<sup>o</sup> 36,  $\omega_{IR} = 853 \text{ cm}^{-1}$  (854), is composed almost entirely of the o-p bending  $\gamma_5$  (55%) of the C<sub>5</sub>H<sub>5</sub> bond. Sym mode n<sup>o</sup> 37 seen by neutrons at  $\omega_n = 698 \text{ cm}^{-1}$  (708) is the trigonal puckering of the ring it consists mainly of alternate torsions  $\tau_i$  up and down of the C-C in the ring (31%). AntiSym mode n<sup>o</sup> 39,  $\omega_R = 588 \text{ cm}^{-1}$  (583), consists of 20% of ring torsions, 21% of o.p. bending motions and 38% of methyl groups torsions. No<sup>o</sup> 43,  $\omega_R = 517 \text{ cm}^{-1}$  (526), appears as a Sym butterfly flapping around the C<sub>2</sub> axis, it consists of 16 % of ring torsions, 36% of o.p. bending motions. No<sup>o</sup> 45,  $\omega_R = 339 \text{ cm}^{-1}$  (332), appears as an AntiSym motion. The PED contains 16% of ring torsion, 33% of o.p. bending motions and 36% of torsions of the methyl groups around their axes. No<sup>o</sup> 48,  $\omega_R = 276 \text{ cm}^{-1}$  (269), involves nearly the same motion as n<sup>o</sup> 45, but the ring oscillates now around an axis perpendicular to the C<sub>2</sub> axis.

### 5.6.3. Methyl torsion modes and/or o.p. bending modes?

Mode n<sup>o</sup>51 ( $\omega_{IR} = 190 \text{ cm}^{-1}$ ,  $\omega_{anh} = 166$ ,  $\omega_n = 188$ ) is a type of breathing motion in which all atomic substituents move in the opposite direction of that of the ring carbon atoms, ( $\Sigma\gamma_i = 19\%$ ). This mode results in an intense neutron scattering at  $193 \text{ cm}^{-1}$  and a medium I.R. absorption at  $190 \text{ cm}^{-1}$ , but it is inactive in Raman spectroscopy. Interestingly, there is also an important contribution to the PED (56%) due to torsions around the axis of the Me<sub>4</sub> and Me<sub>6</sub> groups. These torsion phenomena occur for all the modes with frequencies smaller than  $200 \text{ cm}^{-1}$ . For mode n<sup>o</sup>52,  $\omega_R = 124 \text{ cm}^{-1}$  (135), the calculations indicate in-phase librations of Me<sub>4</sub> and Me<sub>6</sub> groups, the largest contribution to the PED coming from Me<sub>6</sub> (57%). For mode n<sup>o</sup>54,  $\omega_R = 124 \text{ cm}^{-1}$  (136), the librations of the Me<sub>2</sub> and Me<sub>6</sub> groups are out of phase and the largest contribution to the PED is coming from Me<sub>4</sub> (63%). The

assignment of a Raman frequency  $\omega_R$  for n°55 is difficult because its frequency cannot be known using the anharmonic approach ( $\omega_{anh} = -145$ ). Using the harmonic approximation, this mode was assigned as a  $Me_2$  libration ( $\delta_2 = 82\%$ ). Mode n° 56,  $\omega_R = 88\text{ cm}^{-1}$  (66), corresponds to AntiSym C-Br bending ( $-\gamma_1 + \gamma_3$ ) (26%) associated with a libration of  $Me_2$  (23%). The anharmonic DFT calculations did not result in a real frequency for mode n° 57 ( $\omega_{anh} = -56$ ). Using the harmonic approximation ( $\omega_h = 71$ ), the contributions to the PED are attributed to Sym C-Br bending vibrations ( $\gamma_1 + \gamma_3$ ) (16 %) with a larger contribution from the torsional motion of the  $Me_2$  group (53 %). These results confirm the particular behavior of the “symmetrical” methyl group  $Me_2$ . So it is now necessary to discuss the validity of these conclusions knowing the inadequacy of the DFT models to represent the molecular geometry around  $Me_2$ .

#### 5.6.4. Remarks about the methyl groups torsion modes.

In section 2.1, we summarized the results of the structural studies involving neutron diffraction at 120 and 14 K. For the methyl groups  $Me_4$  and  $Me_6$ , a “normal” conformation with three well defined maxima was observed with the protons probability density (PPD) [17]. In contrast, for  $Me_2$ , the PPD is widely spread over a torus. For  $Me_4$  and  $Me_6$ , the DFT calculations generated an equilibrium conformation that is in agreement with the experimental results: for each methyl group, there is a preferred orientation with one C-H bond pointing towards the hydrogen  $H_5$  (Table 1, Fig. 1a and 1c). Starting from these configurations, DFT calculations then provided values equal to 158 and 167  $\text{cm}^{-1}$  using the harmonic approximation and equal to 135 and 136  $\text{cm}^{-1}$  using the anharmonic approximation for the torsion mode of these groups. The calculated frequency does not vary significantly with the plane of symmetry ( $P_{ring}$  or  $P_{perp}$ ) retained for the model molecule studied (Table 3). The intensity calculated for these excitations is very low in I.R., INS and Raman. In fact, based on inelastic neutron scattering experiment using a reduced fixed window [17], it was possible to estimate the hindering potentials  $V_3$  of these methyl groups to be approximately 590  $\text{cm}^{-1}$  plus a small out of phase component  $V_6 = 90\text{ cm}^{-1}$ . Using this potential, the “libration” transition  $0^A-1^A$  (almost equal to  $0^E-1^E$ ) is approximately 145  $\text{cm}^{-1}$ ; this value is compatible with the preceding predictions relative to the Raman line at 136  $\text{cm}^{-1}$ . The problem is different for the  $Me_2$  group. A high-resolution inelastic neutron scattering study of the DBMH crystal [16] revealed that this group is a quasi-free rotor, which is contrary to the crude idea that the “big” bromine atoms will hamper the space around  $Me_2$  and, thus, hinder its “rotation”. Below 40 K, the INS spectra exhibit excitations at 14.7 and 24.6  $\text{cm}^{-1}$  that were assigned to transitions between levels of same symmetry  $0^A-1^A$  and  $0^E-$

<sup>1</sup>E In the Raman spectra [18], scattered lines exist at 15.2 and 23.7 cm<sup>-1</sup>. The fact that the intensities of these transitions are a function of time after freezing is another consequence of the quantum behavior of the rotor Me<sub>2</sub>. Specifically this phenomenon is a consequence of the progressive departure of a part of the population with spin 1/2 into level 0<sup>E</sup> to more stable O<sup>A</sup> level (spin 3/2) via spin conversion. In section 2.2, it was indicated that the DFT calculations could not reproduce the Me<sub>2</sub> experimental conformation because, at equilibrium, the methyl group must always present only three PPD maxima as in Fig. 1b. These calculations were performed strictly using the Born-Oppenheimer approximation to the hydrogen nuclei without considering their spin states. Consequently, the energy states of the methyl groups in small hindering potentials must always be studied as solutions of the Schrodinger equation of a uni-axial quantum rotor.

## 6. Conclusion

**6.1.** An experimental study of the Raman and infra-red spectra of the fully hydrogenated 1,3-dibromo-2,4,6-trimethylbenzene (DBMH) at 288 K was presented. These spectra cover the range 80 - 3500 cm<sup>-1</sup> with a resolution of approximately 3 cm<sup>-1</sup>. Additionally, INS spectra were recorded in the same frequency range below 30 K. In parallel, quantum chemistry calculations of the molecular conformation and of the internal modes of vibration were performed. They consisted of DFT calculations using the MPW1PW91 functional with the LANL2DZ basis set augmented with polarization functions. The molecular equilibrium conformation that was calculated theoretically is very close to that obtained using neutron diffraction at 120 and 14 K with the exception of the Me<sub>2</sub> environment. The DFT calculations fail to produce the exact molecular conformation because they cannot consider the fact that Me<sub>2</sub> is a symmetrical “quantum rotor” with spin states. Nevertheless, these calculations provide an acceptable value for the rotation hindering barrier because they take into account the energy related to the spatial part of the Hamiltonian.

**6.2.** Using the geometry with  $P_{ring}$  as a symmetry plane as the equilibrium molecular geometry, a first calculation of the internal modes frequencies was performed assuming the harmonic approximation. The use of scaling factors of 0.983 for frequencies smaller than 1800 cm<sup>-1</sup> and another factor equal to 0.958 for higher frequencies results in an agreement better than 1.5%. Furthermore, the values of the calculated intensities are qualitatively in agreement with those of the experimental peaks.

**6.3.** Another DFT calculation was performed using the same starting equilibrium conformation, but taking into account the anharmonicity of the motions. The agreement between calculation and experiment is now excellent for all of the frequencies below 1800 cm<sup>-1</sup>, and the calculations

overestimate the C-H stretching modes in the range 2900-3200  $\text{cm}^{-1}$  by less than 1.5 %. It is not necessary to use arbitrary scaling factors.

**6.4.** The contribution of internal vibrations to normal vibration potential energy (PED) was obtained within the harmonic approximation. Rocking, bending and stretching modes that are specific to internal methyl groups deformations during the vibration were localized and the calculations resulted in small differences in the frequencies of the  $\text{Me}_2$  vibrations and those of  $\text{Me}_4$  and  $\text{Me}_6$ . The specific excitations of the methyl torsion modes cannot be seriously calculated using DFT because this type of calculation neglects the spins of the methyl groups. Our previous INS studies [16, 17] demonstrated the necessity to take into account the quantum behavior of the methyl groups, particularly  $\text{Me}_2$ , which is subjected to a small hindering potential because of the symmetry of its environment.

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# Figures captions

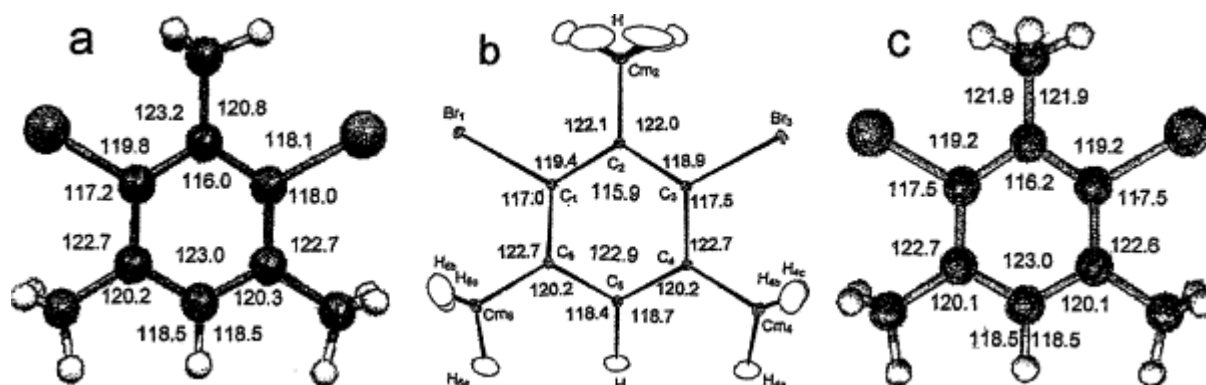


Fig. 1. Bond angles in the DBMH molecule: 1a- Optimized conformation using DFT computation; 1b- Obtained from neutron diffraction at 120 K; 1c- Conformation obtained from DFT calculations with the constraint that the plane perpendicular to the ring plane is a symmetry plane.

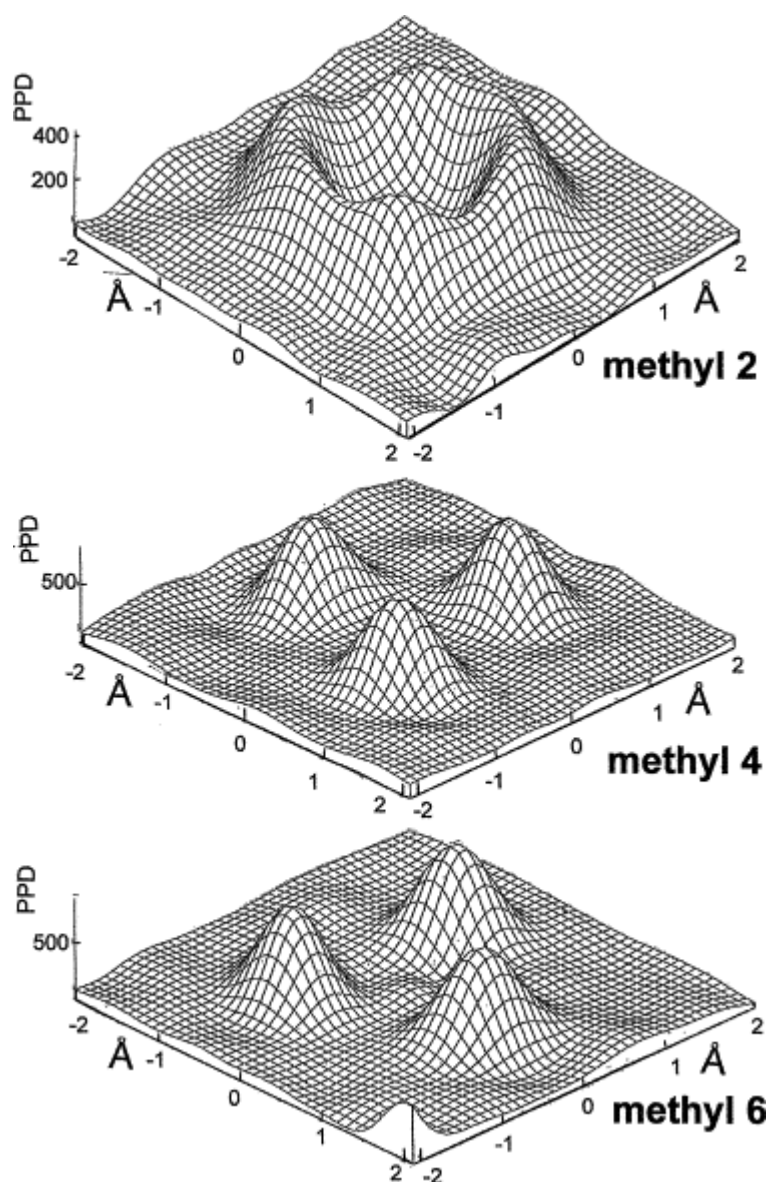


Fig. 2. Proton probability density for methyl groups Me<sub>2</sub>, Me<sub>4</sub> and Me<sub>6</sub> obtained using neutron diffraction on a single crystal of DBMH at 120 K [18]. Each probability density graph is a cut in a plane perpendicular to the C<sub>i</sub>-C<sub>mi</sub> bond and containing the maxima of the PPD (i= 2, 4 or 6).

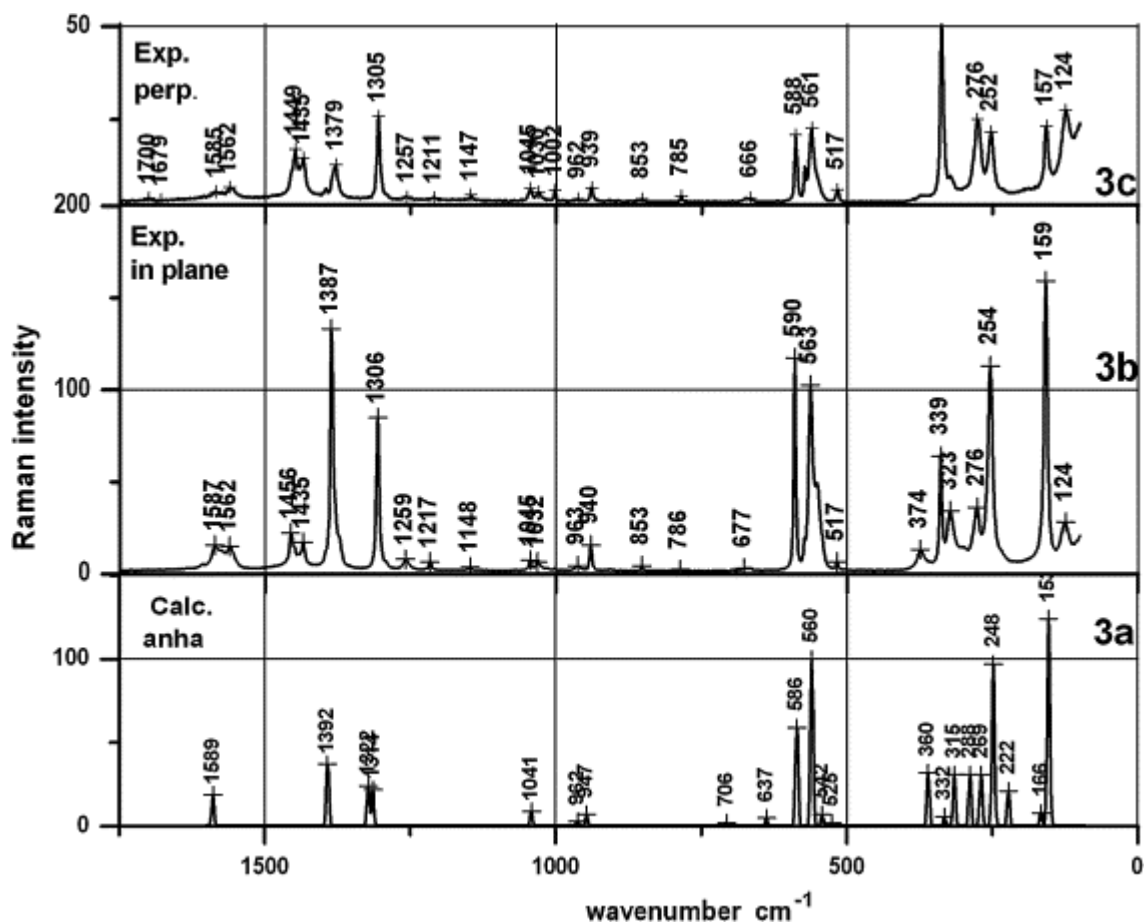


Fig. 3. DBMH Raman spectra: 3a- Calculated using DFT with the optimized conformation; 3b and 3c: Experimental spectra obtained at 288 K using excitation from an He-Ne laser line at 6328 Å, for 3b- the electric field  $E$  of the excitation beam was nearly in the molecular plane, for 3c-  $E$  was nearly normal to the molecular plane.

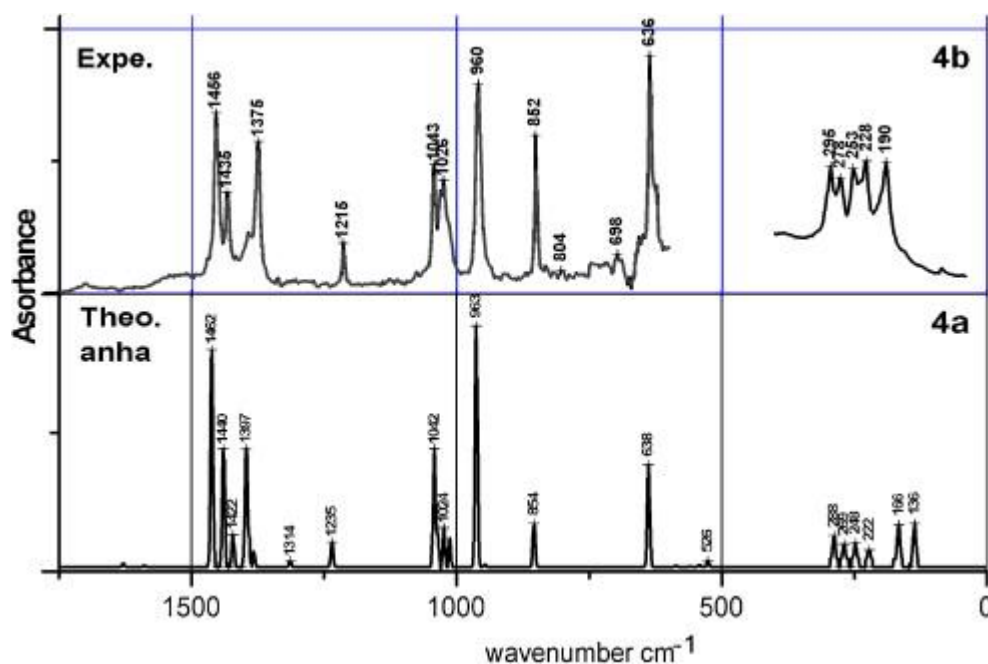


Fig. 4. DBMH infra-red absorption spectra: 4a- Calculated using DFT with the optimized conformation, 4b- In the range 500 - 1800  $\text{cm}^{-1}$ , spectrum obtained using a Varian Fourier transform spectrometer; In the range 100 – 400  $\text{cm}^{-1}$  spectrum obtained on the AILES line using synchrotron radiation at SOLEIL and a Bruker Fourier transform spectrometer.

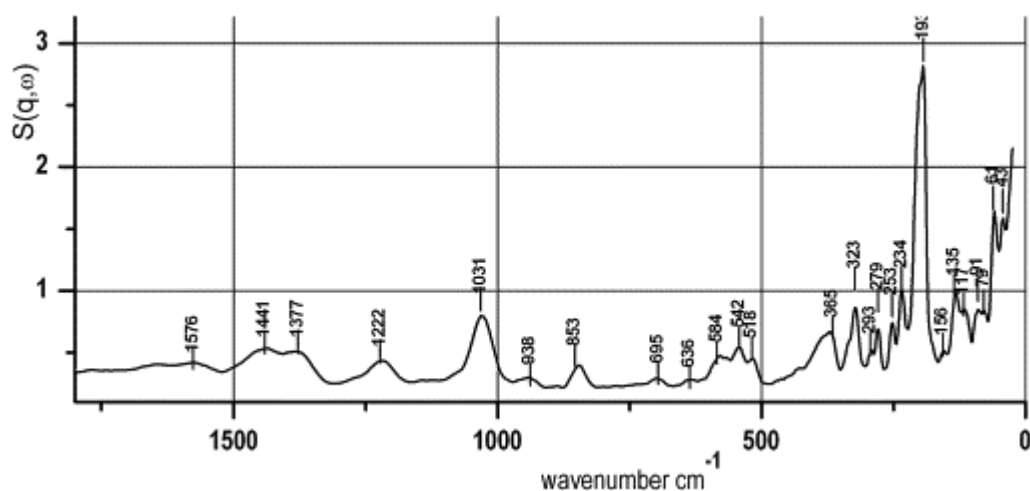


Fig. 5. Inelastic neutron scattering spectra which were obtained below 30 K using the TOSCA spectrometer at ISIS

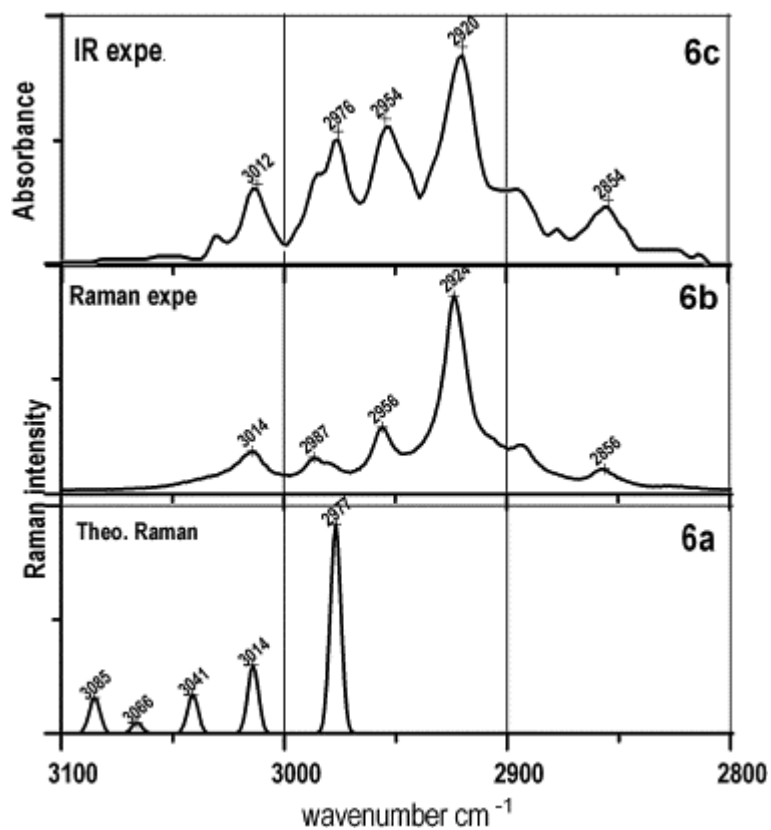


Fig. 6. Spectra in the 2800 - 3200  $\text{cm}^{-1}$  range specific to the methyl groups stretching vibrational modes; 6a- DBMH Raman spectrum calculated by DFT; 6b- Raman experimental spectrum at 288 K; 6c- Infra-red experimental absorption at 288 K.

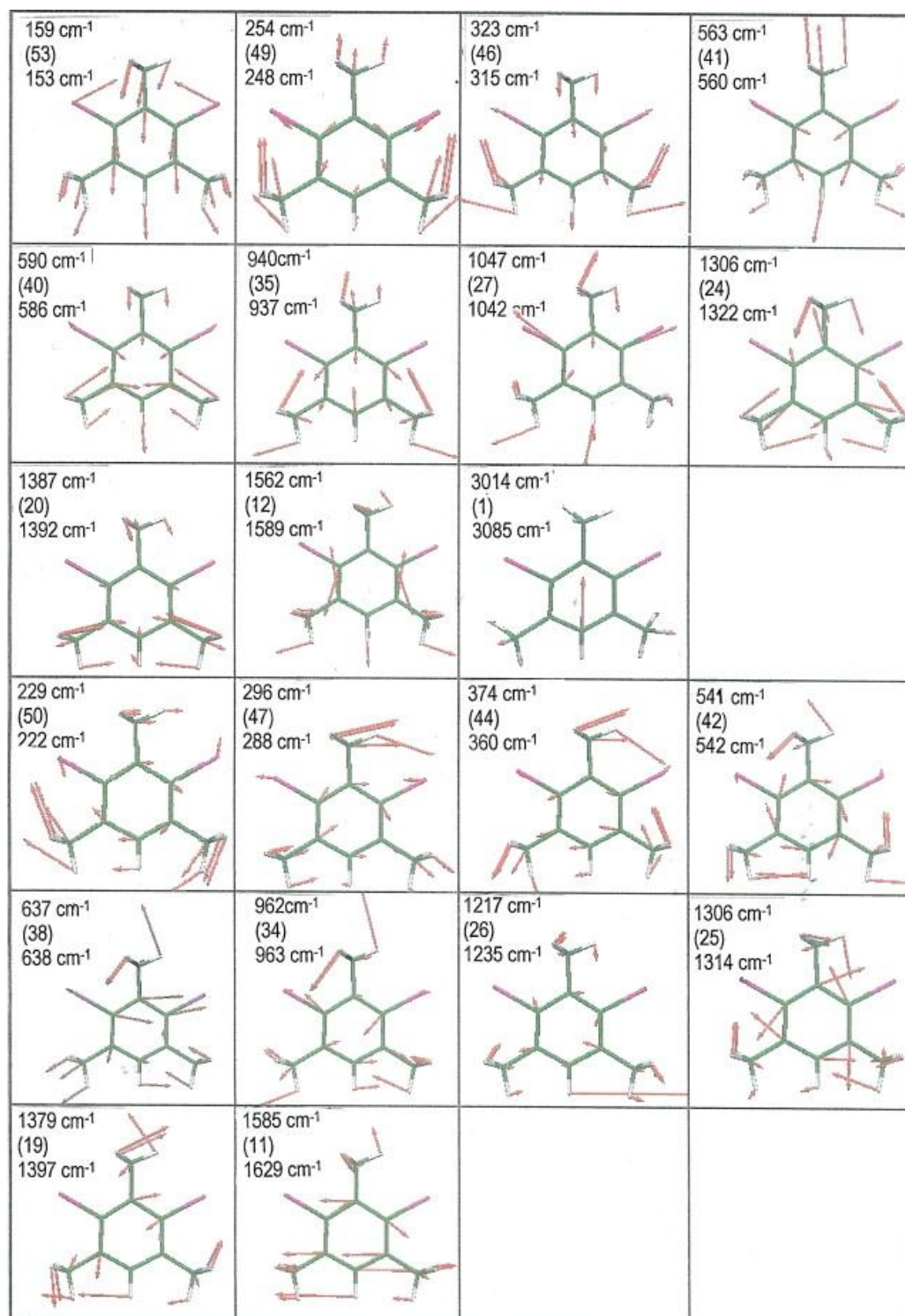


Fig. 7. Illustration of the atomic displacements for the twenty one “skeleton” in-plane vibrational modes (from DFT calculations): the first number at the top of each small figure corresponds to the frequency of an experimental observation (Raman or IR), the second (in parentheses) corresponds to the numbering from the highest frequency to the lowest, and the third number is the calculated anharmonic frequency.

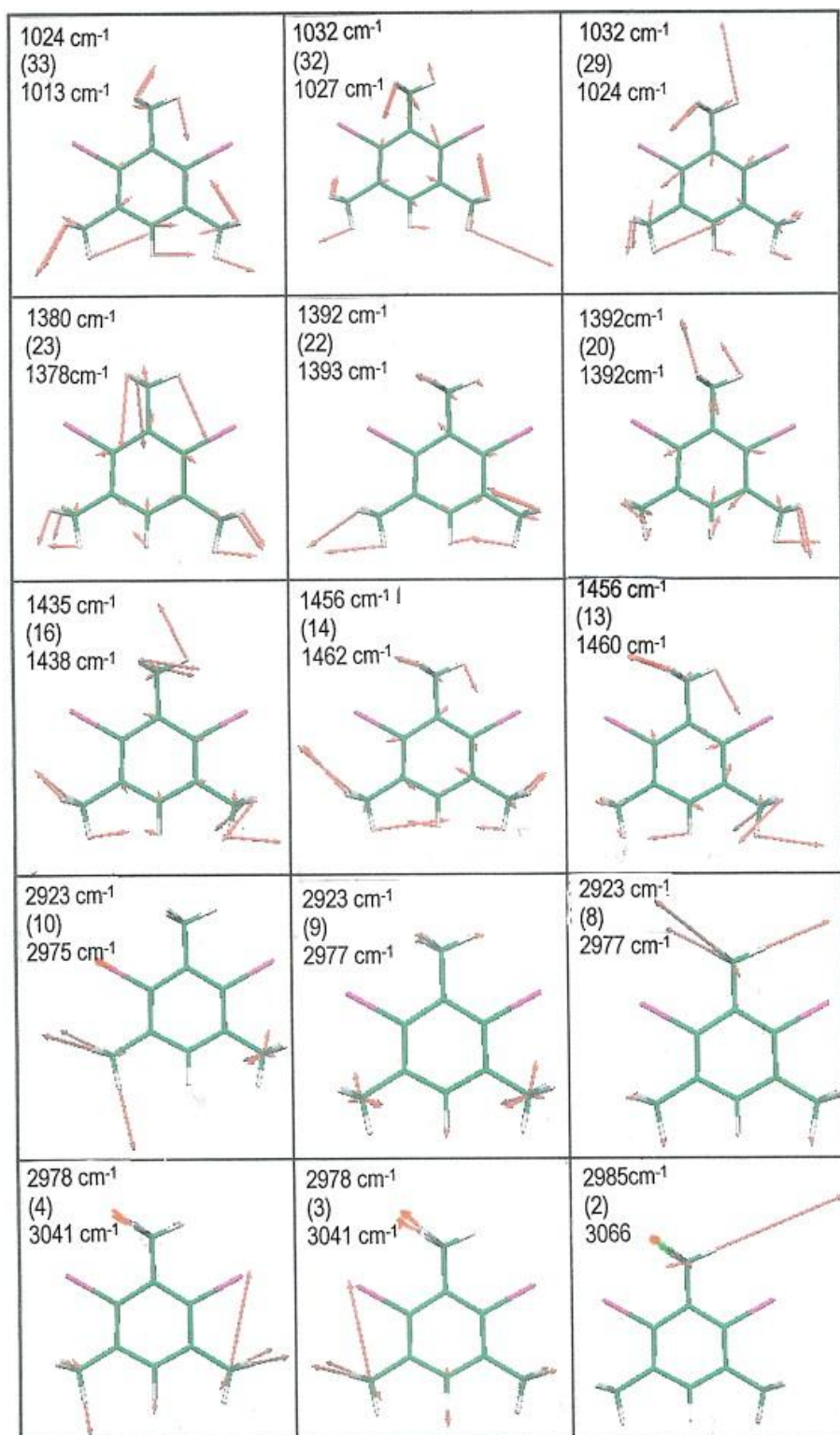


Fig. 8. Illustration of the atomic displacements for the fifteen in plane vibrational modes of “methyl groups” (from DFT calculations)



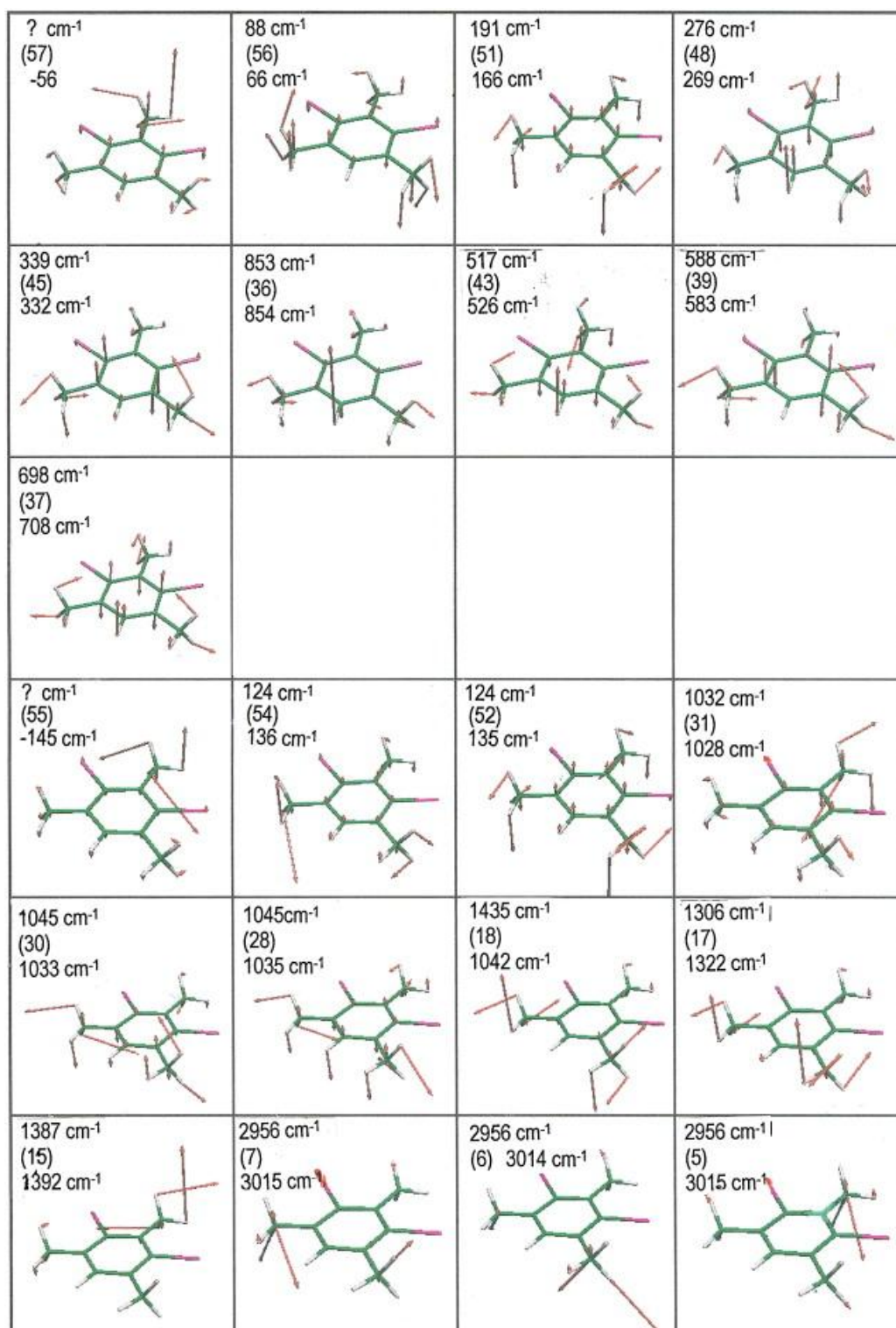


Fig. 9. Illustration of the atomic displacements for the nine “skeleton” and twelve “methyl groups” out of plane vibrational modes (from DFT calculations)

## Tables

**Table 1.** Bond lengths and angles in DBMH. Column n° 2: neutron diffraction at 14 K, col. n° 3: neutron diffraction at 120 K, col. n° 4 and 6: DFT calculations with the MPW1PW91 functional and the LANL2DZ basis set augmented with polarization functions (the plane of symmetry is the benzenic ring plane for n° 4 and the plane perpendicular to the ring for n° 6), col. n° 5: basis set without polarization functions.

1	2	3	4	5	6
Distances Angles	Expe 14 K	Expe 120 K	Cal <sub>1</sub> P <sub>ring</sub>	Cal <sub>2</sub> P <sub>perp</sub>	Cal <sub>3</sub> P <sub>perp</sub>
C <sub>1</sub> -C <sub>2</sub>	1.4025(19)	1.403	1.4061	1.4055	1.4058
C <sub>2</sub> -C <sub>3</sub>	1.4038(15)	1.399	1.4050	1.4108	1.4058
C <sub>3</sub> -C <sub>4</sub>	1.3962(14)	1.392	1.4030	1.4055	1.4012
C <sub>4</sub> -C <sub>5</sub>	1.3943(19)	1.391	1.3960	1.4030	1.3964
C <sub>5</sub> -C <sub>6</sub>	1.3939(15)	1.391	1.3968	1.4030	1.3964
C <sub>1</sub> -C <sub>6</sub>	1.3967(14)	1.392	1.3995	1.4055	1.4013
C <sub>1</sub> -Br <sub>1</sub>	1.9037(19)	1.898	1.9154	1.9645	1.9173
C <sub>2</sub> -C <sub>m2</sub>	1.4993(15)	1.497	1.5020	1.5068	1.5024
C <sub>3</sub> -Br <sub>3</sub>	1.9060(30)	1.904	1.9183	1.9645	1.9170
C <sub>4</sub> -C <sub>m4</sub>	1.5008(16)	1.498	1.5024	1.5067	1.5022
C <sub>5</sub> -H <sub>5</sub>	1.0911(16)	1.091	1.0864	1.0855	1.0864
C <sub>6</sub> -C <sub>m6</sub>	1.5000(20)	1.501	1.5019	1.5067	1.5022
C <sub>6</sub> -C <sub>1</sub> -C <sub>2</sub>	123.60(5)	123.5	123.49	123.46	123.23
C <sub>6</sub> -C <sub>1</sub> -Br <sub>1</sub>	117.06(9)	117.0	118.03	117.45	117.54
C <sub>2</sub> -C <sub>1</sub> -Br <sub>1</sub>	119.34(9)	119.5	118.42	119.08	119.23
C <sub>1</sub> -C <sub>2</sub> -C <sub>3</sub>	115.77(5)	115.9	116.27	116.00	116.24
C <sub>1</sub> -C <sub>2</sub> -C <sub>m2</sub>	121.65(5)	122.1	120.46	122.00	121.94
C <sub>3</sub> -C <sub>2</sub> -C <sub>m2</sub>	122.58(9)	122.0	123.27	122.00	121.82
C <sub>2</sub> -C <sub>3</sub> -C <sub>4</sub>	123.59(9)	123.5	122.99	123.46	123.26
C <sub>2</sub> -C <sub>3</sub> -Br <sub>3</sub>	119.02(4)	118.9	119.79	119.09	119.16
C <sub>4</sub> -C <sub>3</sub> -Br <sub>3</sub>	117.39(4)	117.5	117.23	117.45	117.58
C <sub>3</sub> -C <sub>4</sub> -C <sub>5</sub>	117.08(11)	117.2	117.31	117.09	117.14
C <sub>3</sub> -C <sub>4</sub> -C <sub>m4</sub>	122.70(10)	122.6	122.71	122.88	122.74
C <sub>5</sub> -C <sub>4</sub> -C <sub>m4</sub>	120.22(5)	120.2	119.98	120.03	120.12
C <sub>4</sub> -C <sub>5</sub> -C <sub>6</sub>	122.87(5)	122.9	123.02	122.89	122.99
C <sub>4</sub> -C <sub>5</sub> -H <sub>5</sub>	118.53(8)	118.7	118.50	118.55	118.50
C <sub>6</sub> -C <sub>5</sub> -H <sub>5</sub>	118.60(11)	118.4	118.49	118.55	118.50
C <sub>5</sub> -C <sub>6</sub> -C <sub>1</sub>	117.09(11)	117.1	116.93	117.09	117.14
C <sub>5</sub> -C <sub>6</sub> -C <sub>m6</sub>	120.33(9)	120.2	120.31	120.03	120.12
C <sub>1</sub> -C <sub>6</sub> -C <sub>m6</sub>	122.60(10)	122.7	122.76	122.88	122.74

**Table 2.** DBMH normal modes of vibration: Comparison of the experimental Raman (3), IR (4) and INS (5) bands with the theoretical anharmonic  $\omega_{anh}$  (6), harmonic scaled  $\omega_{hsc}$  (8) and harmonic  $\omega_h$  (10) frequencies, Raman scattering activities (14), Raman intensities: calculated (15) and observed (16), IR absorption: calculated (17) and observed (18)

1	2	Experimental frequencies $\omega$ cm <sup>-1</sup>			Calculated frequencies $\omega$ cm <sup>-1</sup>										Raman activity and intensities			IR absorption	
		3	4	5	6	7	8	9	9b	10	11	12	12b	13	14	15	16	17	18
$N^\circ$	Sym $P_{ring}$	$\omega_R$ Raman	$\omega_{IR}$ IR	$\omega_N$ INS	$\omega_{anh}$ cm <sup>-1</sup>	col.6 - col.3	$\omega_{hsc}$ cm <sup>-1</sup>	col.8 - col.3	$Ra/\omega_h$ cm <sup>-1</sup>	$\omega_h$ cm <sup>-1</sup>	col.10 - col.3	$\omega_{hp}$ cm <sup>-1</sup>	Sym $P_{perp}$	col.12 - col.10	Activity Å <sup>4</sup> /amu	$I_{cal}$ rel.	$I_{Rexp}$ rel.	$A_{cal.}$ km/mol	$I_{IR}$ exp.
1	A'	3014	3012	3000	3085	71	3088	74	0,935	3223	209	3223	A'	0	94	33	17	10	9
2	A'	2987	2985		3066	79	3080	95	0,928	3215	230	3201	A"	-14	27	11	14	5	3
3	A'	2978	2976	2970	3041	63	3051	71	0,935	3185	207	3184	A'	-1	70	36	12	19	15
4	A'	2978	2976		3041	63	3050	72	0,935	3184	206	3184	A"	0	29			9	
5	A"	2956	2954		3015	59	3030	74	0,935	3163	207	3172	A'	9	52	63	26	6	18
6	A"	2956	2954		3014	58	3026	70	0,936	3159	203	3159	A'	0	33			11	
7	A"	2956	2954		3015	59	3026	70	0,936	3159	203	3159	A"	0	90			4	
8	A'	2924	2920	2925	2977	53	2961	38	0,946	3091	168	3087	A'	-4	157	193	97	8	20
9	A'	2924	2920		2977	53	2963	40	0,948	3082	159	3082	A'	0	344			16	

		Experimental frequencies			Calculated frequencies										Raman activity and intensities			IR	
		$\omega$ cm <sup>-1</sup>			$\omega$ cm <sup>-1</sup>													absorption	
10	A'	2924	2920		2975	51	2962	39	0,948	3082	159	3082	A''	0	24			18	
11	A'	1587	-	1578	1629	42	1629	44	0,957	1657	72	1657	A''	0	33	28	15	1	2
12	A'	1562	-	1565	1588	26	1602	40	0,958	1630	68	1637	A'	7	21	19	15	0.5	2
13	A'	1449	1456		1460	11	1479	23	0,967	1505	49	1503	A''	-2	6	23	22	57	66
14	A'	1456	1456		1462	6	1475	19	0,970	1501	45	1502	A'	1	5			38	
15	A''	1435	1430	1441	1422	-13	1465	16	0,972	1490	41	1491	A'	1	8			10	
16	A'	1435	1435		1438	3	1464	29	0,964	1489	54	1488	A''	-1	4	24	17	4	45
17	A''	1435	1435		1441	6	1459	24	0,967	1484	49	1484	A'	0	14			25	
18	A''	1435	1435		1440	5	1458	23	0,935	1483	48	1483	A''	0	5			1	
19	A'	-	1397	1392	1397	0	1417	30	0,935	1442	55	1439	A''	-3	0.4	4	131	30	25
20	A'	1387	1392		1392	1	1407	30	0,936	1429	50	1429	A'	0	34	50		1	57
21	A'	1387	1392		1377	1393	6	1405	26	0,936	1424	37	1425	A'	1	5		20	
22	A'	1387	1375	1384		-3	1397	5	0,946	1420	33	1421	A''	1	7	3			
23	A'	1380	1375	1378		-2	1389	9	0,948	1413	33	1415	A''	2	4	4		4	
24	A'	1305	1313		1322	17	1330	25	0,948	1353	48	1351	A'	-2	20	30	84	0	5

		Experimental frequencies			Calculated frequencies										Raman activity and intensities			IR absorption	
		$\omega$ cm <sup>-1</sup>			$\omega$ cm <sup>-1</sup>														
25	A'	1305	1313		1314	8	1315	9	0,957	1338	32	1338	A"	0	9	22		2	
26	A'	1217	1215	1218	1235	18	1228	11	0,958	1249	32	1249	A"	0	1	1	6	6	20
27	A'	1047	1043	1045	1041	-6	1046	-1	0,967	1064	17	1064	A'	0	5.	9	7	30	45
28	A"	1045	1043		1035	-10	1043	-4	0,970	1061	14	1061	A'	0	2			7	
29	A'	1032	1028	1026	1024	-8	1041	9	0,972	1059	27	1058	A"	-1	0	2	6	6	32 20
30	A"	1045	1045		1033	-12	1040	-1	0,964	1058	13	1057	A"	-1	0.5			4	
31	A"	1032	1026		1028	-4	1033	1	0,967	1051	19	1054	A'	3	0			0	
32	A'	1032	1024		1024	-8	1032	0	0,968	1050	18	1049	A'	-1	1			1	
33	A'	1015	1015		1013	-2	1019	17	0,935	1037	35	1038	A"	1	0.3	1	2	8	
34	A'	962	960		962	0	974	12	0,935	991	29	992	A"	1	1.7	3	4	61	76
35	A'	940	-	938	947	7	947	8	0,936	963	24	964	A'	1	3.40	7	5	0	
36	A"	853	852	853	854	1	870	17	0,936	885	32	885	A'	0	0.2	1	4	10	40
37	A"	-	-	695	706	11	720	22	0,954	732	34	731	A'	-1	0.7	2	-	0	7
38	A'	-	636	637	637	1	641	4	0,977	652	16	652	A"	0	1.2	5	-	26	52
39	A"	588	-	585	583	-5	584	-4	0,990	594	6	594	A"	0	2.8	68	116	0	2

		Experimental frequencies			Calculated frequencies										Raman activity and intensities			IR absorption	
		$\omega$ cm <sup>-1</sup>			$\omega$ cm <sup>-1</sup>														
40	A'	590	-		586	-4	582	-6	0,997	592	2	592	A'	0	11			0	
41	A'	561	-		560	-1	558	-3	0,991	568	5	568	A'	0	19	100	100	0	
42	A'	-	540	542	542	2	537	-7	0,991	546	6	548	A"	2	1.2	7	-	1	5
43	A"	517	517	518	525	9	532	15	0,956	541	24	539	A'	-2	0.3	2	6	2	11
44	A'	374	-	369	360	-14	362	-12	1,016	368	-6	377	A"	9	2.7	32	13	0	-
45	A"	339	-		332	-7	336	-3	0,988	342	3	342	A"	0	0.4	6	63	0	-
46	A'	323	-	323	315	-8	314	-9	1,013	319	-4	319	A'	0	2	31	34	0	-
47	A'	-	295	293	288	-7	285	-11	1,021	290	-6	303	A"	13	1.7	31	-	3	4
48	A"	276	276	279	269	-7	277	1	0,979	282	6	281	A'	-1	1.5	31	36	3	3
49	A'	252	253	253	248	-4	248	-4	1,008	252	0	252	A'	0	4	97	111	3	3
50	A'	-	228	234	222	-7	224	-5	1,004	228	-1	230	A"	2	0.7	21	-	2	4
51	A"	-	190	193	166	-24	185	-6	1,016	188	-3	188	A'	0	0.2	8	-	4	4
52	A"	124	-	135	135	11	165	-1	0,743	167	43	167	A"	0	0.3			4	-
53	A'	157	-	156	153	-4	156	-1	1,004	158,4	1	160	A'	2	2	124	136		1
54	A"	124	-	135	136	12	155	21	0,785	158	34	159	A'	1	0.1	34	25	4	-

		Experimental frequencies			Calculated frequencies										Raman activity and intensities			IR absorption	
		$\omega$ cm <sup>-1</sup>			$\omega$ cm <sup>-1</sup>														
55	A"	-	-	-	-145	?	117	-		119	-	81	A"	-38	0.3				-
56	A"	88	-	91	66	?	80	-8	1,086	81	-7	80	A'	-1	0.6				-
57	A"	21	-	-	-56	?	70	49	?	71	-	-94	A"	?	0.0				-

**Table 3.** The predominant components of the PED matrix (relative phases and % contributions) for the in-plane modes of vibration of DBMH

1	2	3	4	5	6	7	8
n°	$\omega_R$ cm <sup>-1</sup>	% C <sub>a</sub> X stretching %C <sub>m</sub> H stretching	% -C <sub>a</sub> C <sub>a</sub> C <sub>a</sub> - bending	% -C <sub>a</sub> C <sub>a</sub> X bending	% ring C <sub>a</sub> C <sub>a</sub> stretching	% C <sub>a</sub> C <sub>a</sub> C <sub>m</sub> H torsion	% C <sub>a</sub> C <sub>m</sub> H + HC <sub>m</sub> H bending( $\chi_6, \theta_6$ )
1	3014	67 s <sub>5</sub>	14 $\alpha_5$	-	(t <sub>4</sub> +t <sub>5</sub> ) (6)		
2	2987	48 r <sub>2a</sub> - 30 (r <sub>2b</sub> + r <sub>2c</sub> )		-	-		
3	2978	-8 r <sub>4a</sub> + 12 (r <sub>4b</sub> + r <sub>4c</sub> ) - 35 r <sub>6a</sub> + 28 (r <sub>6b</sub> + r <sub>6c</sub> )		-	-		
4	2978	-35 r <sub>4a</sub> + 28 (r <sub>4b</sub> + r <sub>4c</sub> ) + 8 r <sub>6a</sub> + 12 (r <sub>6b</sub> + r <sub>6c</sub> )		-	-		
8	2924	20 r <sub>2a</sub> - 66 (r <sub>2b</sub> + r <sub>2c</sub> )					
9	2924	12 r <sub>4a</sub> + 30 (r <sub>4b</sub> + r <sub>4c</sub> ) + 14 r <sub>6a</sub> + 34(r <sub>6b</sub> + r <sub>6c</sub> )					
10	2924	14 r <sub>4a</sub> + 34 (r <sub>4b</sub> + r <sub>4c</sub> ) - 12 r <sub>6a</sub> - 30 (r <sub>6b</sub> + r <sub>6c</sub> )					
11	1587	4 s <sub>2</sub> + 7 (s <sub>4</sub> + s <sub>6</sub> )	10 ( $\alpha_1$ - $\alpha_3$ ) - 14 ( $\alpha_4$ - $\alpha_6$ )	15 $\beta_5$ + 3 $\beta_2$ - 5( $\beta_1$ + $\beta_3$ )	9 (t <sub>1</sub> -t <sub>2</sub> +t <sub>3</sub> )+16(t <sub>4</sub> - t <sub>5</sub> +t <sub>6</sub> )	-	18 ( $\chi_2, \theta_2$ + $\chi_4, \theta_4$ + $\chi_6, \theta_6$ )
12	1562	4 s <sub>2</sub>	15( $\alpha_2$ + $\alpha_5$ ) - 15( $\alpha_1$ + $\alpha_3$ + $\alpha_4$ + $\alpha_6$ )	12(- $\beta_1$ + $\beta_3$ - $\beta_4$ + $\beta_6$ )	16(t <sub>3</sub> +t <sub>6</sub> ) - 16(t <sub>1</sub> +t <sub>2</sub> +t <sub>4</sub> +t <sub>5</sub> )	-	14 ( $\chi_2, \theta_2$ + $\chi_4, \theta_4$ + $\chi_6, \theta_6$ )
13	1456	-	-	- 6 $\beta_5$	4(t <sub>2</sub> -t <sub>5</sub> )	11 $\eta_2$ + 15 $\eta_6$	- 18( $\chi_2, \theta_2$ ) + 23( $\chi_6, \theta_6$ )
14	1456	-	-	4 $\beta_5$	-	- 6 $\eta_2$ + 19 $\eta_4$ + 6 $\eta_6$	6 ( $\chi_2, \theta_2$ )-25 ( $\chi_4, \theta_4$ )+ 6( $\chi_6, \theta_6$ )
16	1435	-	-	-	-	16 $\eta_2$ +10 $\eta_4$ +14 $\eta_6$	-16( $\chi_2, \theta_2$ ) + 10( $\chi_4, \theta_4$ )- 14
19	1387	4 (s <sub>1</sub> -s <sub>3</sub> )	8( $\alpha_1$ - $\alpha_3$ )	4 $\beta_2$ - 6 $\beta_5$ - 4( $\beta_4$ - $\beta_6$ )	12 ( -t <sub>1</sub> + t <sub>2</sub> + t <sub>3</sub> - t <sub>4</sub> - t <sub>6</sub> )	-12 $\eta_2$ + 8 $\eta_4$ -7 $\eta_6$	17( $\chi_2, \theta_2$ )-11( $\chi_4, \theta_4$ )+7( $\chi_6, \theta_6$ )
20	1387	7 (s <sub>4</sub> +s <sub>6</sub> )	6( $\alpha_4\phi_4$ + $\alpha_6\phi_6$ )- 6( $\alpha_1\phi_1$ + $\alpha_3\phi_3$ + $\alpha_5\phi_5$ )	-	4(t <sub>4</sub> +t <sub>6</sub> )	- 18 ( $\eta_4$ + $\eta_6$ )	8(- $\chi_2, \theta_2$ )+21(- $\chi_4, \theta_4$ )+26 (- $\chi_6, \theta_6$ )
21	1387	- 5 s <sub>2</sub>	2 $\alpha_5\phi_5$ - 2 $\alpha_2\phi_2$	4(- $\beta_1$ + $\beta_3$ ) + 6( $\beta_4$ - $\beta_6$ )	6 (t <sub>1</sub> +t <sub>2</sub> +t <sub>3</sub> +t <sub>6</sub> ) - 6 (t <sub>4</sub> +t <sub>5</sub> )	25 ( $\eta_2$ - $\eta_4$ + $\eta_6$ )	18( $\chi_2, \theta_2$ )- 6( $\chi_4, \theta_4$ )+13( $\chi_6, \theta_6$ )
22	1387	4 (s <sub>4</sub> -s <sub>6</sub> )	-	4 $\beta_5$	-	-5 $\eta_2$ + 3 $\eta_4$ - 8 $\eta_6$	33( $\chi_4$ , - $\theta_4$ ) + 25(- $\chi_6$ , $\theta_6$ )
23	1380	-	-	10( $\beta_1$ - $\beta_3$ + $\beta_4$ - $\beta_6$ )	7 (t <sub>1</sub> +t <sub>2</sub> +t <sub>3</sub> -t <sub>6</sub> )	5 $\eta_2$ - 8 $\eta_4$ - 3 $\eta_6$	-32( $\chi_2, \theta_2$ )+11( $\chi_4$ , - $\theta_4$ )+7( $\chi_6$ , - $\theta_6$ )



24	1305	$15 (s_2 + s_4 + s_6)$	$-6(\alpha_2 + \alpha_4 + \alpha_6) + 6(\alpha_1 + \alpha_3 + \alpha_5)$	$-5(\beta_1 + \beta_3 + \beta_5)$	$-12 (t_2 + t_4 + t_6)$	$8 (\eta_4 + \eta_6)$	$31 (\chi_2, \theta_2 + \chi_4, \theta_4 + \chi_6, \theta_6)$
25	1305	$-5 (s_2 + s_4 + s_6)$		$15(\beta_1 - \beta_2 + \beta_3 - \beta_4 + \beta_5 - \beta_6)$	$33 (t_1 - t_2 + t_3 - t_4 + t_5 - t_6)$	$22 (\eta_2 + \eta_4 + \eta_6)$	$8(\chi_4, \theta_4) - 12 (\chi_2, \theta_2 + \chi_6, \theta_6)$
26	1217	$4 (s_1 - s_3) - 8 (s_4 - s_6)$	-	$28\beta_5 - 4\beta_2 - 6(\beta_1 + \beta_3)$	$10 (t_3 + t_4 - t_5 - t_6)$	$15 (\eta_2 + \eta_4 + \eta_6)$	$-7(\chi_2, \theta_2) + 12 (\chi_4, \theta_4 - \chi_6, \theta_6)$
27	1047	$12(s_1 + s_3) - 3(s_4 - s_6)$	$48(\alpha_1\phi_1 - \alpha_2\phi_2 + \alpha_3\phi_3 - \alpha_4\phi_4 + \alpha_5\phi_5 - \alpha_6\phi_6)$	-	$-6 (t_1 + t_2 + t_3 + t_4 + t_5 + t_6)$	$10 (\eta_2 + \eta_6)$	$9 (\chi_2, \theta_2)$
29	1032	$8 (s_2 - s_4)$	$6(-\alpha_1 + \alpha_2 - \alpha_3 + \alpha_4)$		$6 (t_1 + t_2 - t_3 - t_4)$	$-27 (\eta_2 + \eta_4)$	$19 (\chi_2, \theta_2) - 16 (\chi_4, \theta_4)$
32	1032	$8 (-s_2 + s_4)$		$-3\beta_5$		$13 (\eta_2 + \eta_4) + 18\eta_6$	$13 (\chi_2 + \chi_4) + 17\chi_6$
33	1015	$6 (s_4 - s_6)$		$+3\beta_5 + 2\beta_4$		$28(-\eta_2 - \eta_4) + 11\eta_6$	$-11(\chi_2, \theta_2) - 18(\chi_4, \theta_4) + 13(\chi_6, \theta_6)$
34	962	$16 (s_1 - s_3 + s_4 - s_6)$	$14 (\alpha_1 - \alpha_3) + 7(\alpha_4 - \alpha_6)$	-	-	$15\eta_2 + 4\eta_4$	$19 (\chi_2, \theta_2) + 6 (\chi_4, \theta_4)$
35	940	$5 s_2 - 6 (s_4 + s_6)$	$5(\alpha_2\phi_2 - \alpha_4 - \alpha_6) + 8\alpha_5\phi_5 - 2(\alpha_1 + \alpha_3)$	$5(-\beta_1 - \beta_3 + \beta_4 + \beta_6)$	$8 (-t_1 - t_2 + t_4 + t_5)$	$22 (\eta_4 + \eta_6)$	$22 (\chi_4 + \chi_6) - 8 (\theta_4 + \theta_6)$
38	637	$22(s_1 - s_3) - 8(s_4 - s_6)$	$8 (\alpha_1 - \alpha_3) + 10 (\alpha_4 - \alpha_6)$	$6\beta_2 - 5\beta_5$	-	$6\delta_2$	$7 (\chi_2, \theta_2)$
40	590	$8(s_1 + s_3) - 14(s_4 + s_6)$	$-8(\alpha_1 + \alpha_3) - 16\alpha_5\phi_5 + 28(\alpha_4 + \alpha_6)$	$-12(\beta_4 + \beta_6)$	-	$4 (\eta_4 + \eta_6)$	$-5 (\theta_4 + \theta_6)$
41	561	$11 (s_1 + s_3) + 8 s_2$	$10(\alpha_1\phi_1 + \alpha_3\phi_3) + 6(\alpha_4 + \alpha_6) - (11\alpha_2\phi_2 + 17\alpha_5\phi_5)$	$-10 (\beta_4 + \beta_6)$	$13 (t_2 + t_3 + t_4 + t_5)$	-	-
42	541	---	-	$16(\beta_1 + \beta_3) + 34(\beta_2 + \beta_4 + \beta_6)$	$4 (t_1 - t_3)$	$16 (\eta_2 + \eta_4 + \eta_6)$	$18 (\chi_2 + \chi_4 + \chi_6)$
44	374	$11(s_1 - s_3) + 5(s_4 - s_6)$	-	$-16\beta_2 + 16(\beta_4 + \beta_6)$	$6 (t_1 - t_2 + t_3 - t_6)$	$5\eta_2 - 7 (\eta_4 + \eta_6)$	$-12 (\chi_2 - \chi_4 + \chi_6)$
46	323	$14 (s_1 + s_3)$	-	$6(\beta_1 - \beta_3) + 36(\beta_4 - \beta_6)$	$6 (t_4 - t_6)$	$11 (\eta_4 + \eta_6)$	$12 (\chi_4 + \chi_6)$
47	296	$18 (s_1 - s_3)$	$12(-\alpha_1 + \alpha_3) + 23(-\alpha_4\phi_4 + \alpha_6\phi_6)$	$18(\beta_1 - \beta_3) - 12\beta_2$	-	-	-
49	252	$22 (s_1 + s_3)$	$-9(\alpha_1\phi_1 + \alpha_3\phi_3) + 12\alpha_2\phi_2 + 10\alpha_5\phi_5$	$-9(\beta_1 - \beta_3) - 20(\beta_4 - \beta_6)$	-	-	-
50	229	$11 (s_1 - s_3)$	$6 (-\alpha_1 + \alpha_3)$	$33(\beta_1 + \beta_3) - 21(\beta_4 + \beta_6)$	-	$5 (\eta_4 + \eta_6)$	$6 (\chi_4 + \chi_6)$
53	159	$14 (s_1 + s_3)$	-	$52(\beta_1 - \beta_3) - 6(\beta_4 - \beta_6)$	$8 (t_1 + t_2)$	-	-

Remark- the meaning of the symbols used is defined in the text (section 5.1)

**Table 4.** The predominant components of the PED matrix (relative phases and % contributions) for the out-of-plane modes of vibration of DBMH  
Remark- the meaning of the symbols used is defined in the text (section 5.1)

1	2	3	4	5	6	7	8	9
n°	$\omega_R$ cm <sup>-1</sup>	$\omega_{anha}$	Me	% C <sub>m</sub> H stretching	% C <sub>a</sub> C <sub>a</sub> C <sub>a</sub> bending o.p.	% C <sub>a</sub> C <sub>a</sub> X bending o.p.	% C <sub>a</sub> C <sub>a</sub> C <sub>m</sub> H torsion	% C <sub>a</sub> C <sub>m</sub> H + HC <sub>m</sub> H bending
5	2956	3014	2	72 (r <sub>2b</sub> <sup>2</sup> - r <sub>2c</sub> <sup>2</sup> )				
6	2956	3015	6	16(-r <sub>6b</sub> <sup>2</sup> +r <sub>6c</sub> <sup>2</sup> )+58(r <sub>4b</sub> <sup>2</sup> - r <sub>4c</sub> <sup>2</sup> )				
7	2956	3015	4	56(r <sub>6b</sub> <sup>2</sup> -r <sub>6c</sub> <sup>2</sup> ) +16(r <sub>4b</sub> <sup>2</sup> -r <sub>4c</sub> <sup>2</sup> )				
15	1435	1422	2				47 $\delta_2$	10 $\chi_2^2$ -33 $\theta_2^2$
17	1435	1441	4 + 6				28 $\delta_4$ + 20 $\delta_6$	5 $\chi_4^2$ + 4 $\chi_6^2$ -20 $\theta_4^2$ -14 $\theta_6^2$
18	1435	1440	4 - 6				20 $\delta_4$ - 29 $\delta_6$	4 $\chi_4^2$ -5 $\chi_6^2$ -14 $\theta_4^2$ +20 $\theta_6^2$
28	1045	1035	4 + 6		12 ( $\tau_3$ - $\tau_4$ + $\tau_5$ - $\tau_6$ )	-10( $\gamma_4$ + $\gamma_6$ ) +9 $\gamma_5$	11 $\delta_4$ +14 $\delta_6$	7 $\chi_4^2$ -10 $\chi_6^2$ - 3 $\theta_6^2$
30	1045	1033	- 4 + 6		15 ( $\tau_3$ - $\tau_4$ - $\tau_5$ + $\tau_6$ )	13 ( $\gamma_4$ - $\gamma_6$ )	-13 $\delta_4$ +18 $\delta_6$	13 $\chi_4^2$ +4 $\theta_4^2$ +10 $\chi_6^2$ +3 $\theta_6^2$
31	1032	1028	2		15 ( $\tau_1$ - $\tau_2$ )	9 $\gamma_2$ -6 ( $\gamma_1$ + $\gamma_3$ )	12 (- $\delta_4$ + $\delta_6$ ) + 24 $\delta_2$	16 $\chi_2^2$ +8 $\theta_2^2$ - 4 $\chi_6^2$
36	853	856			10 ( $\tau_4$ - $\tau_5$ )	55 $\gamma_5$	12 (- $\delta_4$ + $\delta_6$ )	9 ( $\chi_4^2$ - $\chi_6^2$ )
37	698	708			31( $\tau_1$ - $\tau_2$ + $\tau_3$ - $\tau_4$ + $\tau_5$ - $\tau_6$ )	11( $\kappa_1$ - $\kappa_2$ + $\kappa_3$ - $\kappa_4$ - $\kappa_6$ ) +8 $\kappa_5$	22 ( $\delta_4$ + $\delta_6$ )	6 ( $\chi_2^2$ - $\chi_4^2$ - $\chi_6^2$ )
39	588	583			20( $\tau_1$ - $\tau_2$ + $\tau_3$ - $\tau_4$ + $\tau_5$ - $\tau_6$ )	21( $\gamma_1$ - $\gamma_3$ + $\kappa_4$ - $\kappa_6$ )	30 ( $\delta_4$ - $\delta_6$ ) +8 $\delta_2$	4 ( - $\chi_4^2$ + $\chi_6^2$ )
43	517	526			16( $\tau_1$ - $\tau_2$ + $\tau_4$ - $\tau_5$ )	28( $\gamma_1$ + $\gamma_2$ + $\gamma_3$ + $\gamma_4$ + $\gamma_6$ ) - 8 $\kappa_5$	-19 $\delta_2$ +23(- $\delta_4$ + $\delta_6$ )	-
45	339	332			16( $\tau_1$ - $\tau_3$ - $\tau_4$ + $\tau_6$ )	33( $\gamma_1$ - $\gamma_3$ - $\gamma_4$ + $\gamma_6$ )	36 ( $\delta_2$ + $\delta_4$ - $\delta_6$ )	-
48	276	270			4( $\tau_1$ - $\tau_2$ ) - 6( $\tau_4$ - $\tau_5$ )	30 ( $\gamma_1$ - $\gamma_2$ + $\gamma_3$ - $\gamma_4$ - $\kappa_5$ - $\gamma_6$ )	-8 $\delta_2$ + 13 $\delta_4$ - 18 $\delta_6$	
51	191	166			-	7 ( $\gamma_1$ + $\gamma_2$ + $\gamma_3$ ) + 12( $\gamma_4$ + $\gamma_6$ )	8 $\delta_2$ + 35 $\delta_4$ + 21 $\delta_6$	
52	124	135	6	-			30 $\delta_4$ + 57 $\delta_6$	
54	124	136	4	-			- 16 $\delta_2$ + 63 $\delta_4$	
55	-	-145	2				82 $\delta_2$	
56	88	66				-14 $\gamma_1$ +12 $\gamma_3$ +11( $\kappa_4$ - $\kappa_6$ )	- 23 $\delta_2$ + 12 $\delta_6$	
57	21	-56				7 $\gamma_1$ + 9 $\gamma_3$	-53 $\delta_2$ + 9 $\delta_4$ - 7 $\delta_6$	

